

DNS AND LES OF BUOYANCY-DRIVEN TURBULENCE AT HIGH RAYLEIGH NUMBERS: NUMERICAL METHODS AND SUBGRID-SCALE MODELS

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INTRODUCTION

Buoyancy-driven flows have always been an important subject of scientific studies with numerous applications in environment and technology. The most famous example thereof is the thermally driven flow developed in a fluid layer heated from below and cooled from above, *i.e.* the Rayleigh-Bénard convection (RBC). It constitutes a canonical flow configuration that resembles many natural and industrial processes, such as solar thermal power plants, indoor space heating and cooling, flows in nuclear reactors, electronic devices, and convection in the atmosphere, oceans and the deep mantle.

In the last decades significant efforts, both numerically and experimentally, have been directed at investigating the mechanisms and the detailed scaling behavior of the Nusselt number as a function of Rayleigh and Prandtl numbers in the general form $Nu \propto Ra^\gamma Pr^\beta$. In this regard, Figure 1 shows the predictions of the Nu -number based on the classical Grossmann-Lohse (GL) theory [1] and its subsequent corrections [2, 3] where different scaling regimes, characterized by their corresponding exponents γ and β , are identified. Assuming this power-law scalings and following the same reasonings as in Ref. [4] leads to the estimations for the number of grid points shown in Figure 2 (top). This corresponds to mesh resolution requirements for DNS and clearly explain why nowadays DNS of RBC is still limited to relatively low Ra -numbers. However, many of the above-mentioned applications are governed by much higher Ra numbers, located in the region of the $\{Ra, Pr\}$ phase space where the thermal boundary layer becomes turbulent (see the black dash-dotted line in Figure 2). This region corresponds to the so-called asymptotic Kraichnan or ultimate regime of turbulence, with $\gamma = 1/2$. On the other hand, reaching such Ra -numbers experimentally while keeping the basic assumptions (Boussinesq approximation, adiabaticity of the closing walls, isothermal horizontal walls, perfectly smooth surfaces...) is a very hard task; therefore, the observation of the Kraichnan regime also remains elusive [2, 3].

LES OF BUOYANCY-DRIVEN TURBULENCE

In this context, we may turn to LES to predict the large-scale behavior of incompressible turbulent flows driven by buoyancy at very high Ra -numbers. In LES, the large-scale motions are explicitly computed, whereas the effects of small-scale motions are modeled. Since the advent of CFD, many subgrid-scale (SGS) models have been proposed and success-

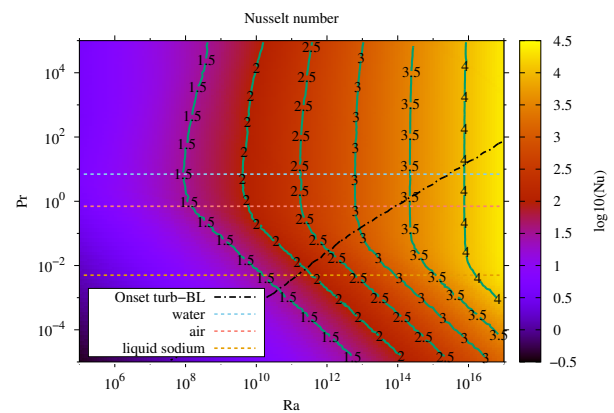


Figure 1: Estimation of the Nusselt number of a RBC in the $\{Ra, Pr\}$ phase space given by the classical GL theory [1] and its subsequent corrections [2]. Green solid isolines represent the \log_{10} of the Nusselt. Three dashed horizontal lines correspond to three different working fluids: water ($Pr = 7$), air ($Pr = 0.7$) and liquid sodium ($Pr = 0.005$). Black dash-dotted line is an estimation for the onset of turbulence in the thermal boundary layer.

fully applied to a wide range of flows. However, there still exists inherent difficulties in the proper modelization of the SGS heat flux. This was analyzed in detail in the PRACE project entitled "Exploring new frontiers in Rayleigh-Bénard convection" (33.1 millions of CPU hours on MareNostrum4 in 2018-2019), where DNS simulations of air-filled ($Pr = 0.7$) RBC up to $Ra = 10^{11}$ were carried out using meshes up to 5600M grid points (see dots displayed in Figure 2, top). These results shed light into the flow topology and the small-scale dynamics which are crucial in constructing the turbulent wind and energy budgets [5]. Moreover, it also provided new insights into the preferential alignments of the SGS and its dependence with the Ra -numbers [6], highlighting that the modelization of the SGS heat flux is the main difficulty that (still) precludes reliable LES of buoyancy-driven flows at (very) high Ra -numbers. This inherent difficulty can be by-passed by carrying out simulations at low-Prandtl numbers. In this case, the ratio between the Kolmogorov length scale and the Obukhov-Corrsin length scale (the smallest scale for the temperature field) is given by $Pr^{3/4}$; therefore, for instance, at $Pr = 0.005$ (liquid sodium) we have a separation of more than one decade. Hence, it is possible to combine an LES simulation for the velocity field (momentum equation) with the numerical resolution of all the

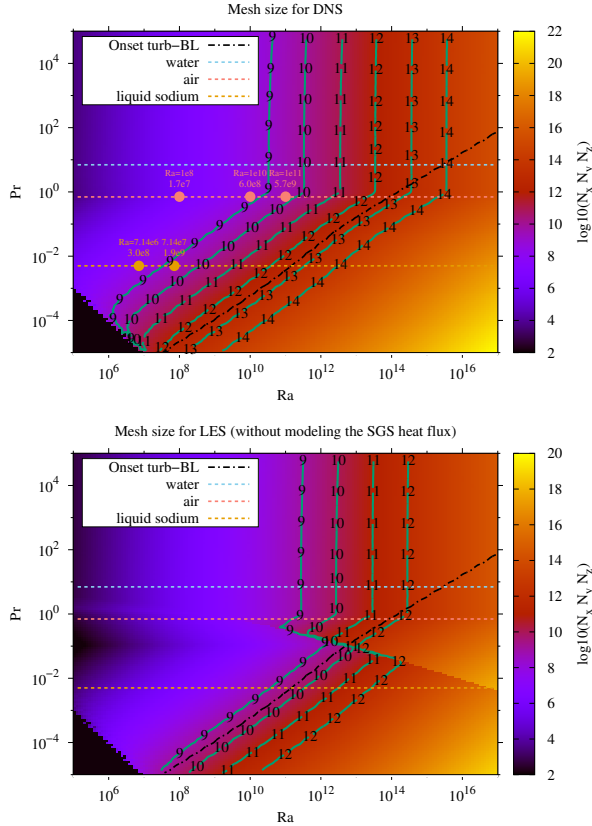


Figure 2: Estimation of the mesh sizes for DNS (top) and LES (bottom) simulations of RBC in the $\{Ra, Pr\}$ phase space. LES estimations assume that thermal scales are fully resolved, *i.e.* no SGS heat flux model is needed. Green solid isolines represent the \log_{10} of the total number of grid points. Dashed horizontal lines correspond to three different working fluids: water ($Pr = 7$), air ($Pr = 0.7$) and liquid sodium ($Pr = 0.005$). Dots displayed on top of these lines correspond to the DNS simulations carried out in previous studies [4, 5, 6]. Black dash-dotted line is an estimation for the onset of turbulence in the thermal boundary layer.

thermal scales. Results obtained in Ref. [6] suggest that accurate predictions of the overall Nu can be obtained with meshes significantly coarser than for DNS (*e.g.* in practice for $Pr = 0.005$ we can expect mesh reductions in the range 10^2 - 10^3 for the total number of grid points). This can be clearly observed in Figure 2 (bottom), where estimations of the mesh size for LES are given with the assumption that thermal scales are fully resolved. This opens the possibility to reach the ultimate regime carrying out LES at low- Pr using meshes of 10^{10} - 10^{11} grid points. Nevertheless, to do so, we firstly need to combine proper numerical techniques for LES (also DNS) with an efficient use of modern supercomputers.

NUMERICAL METHODS AND ALGORITHMS FOR LARGE-SCALE SIMULATIONS ON MODERN SUPERCOMPUTERS

The essence of turbulence are the smallest scales of motion. They result from a subtle balance between convective transport and diffusive dissipation. Mathematically, these terms are governed by two differential operators differing in symmetry: the convective operator is skew-symmetric, whereas the diffusive is symmetric and negative-definite. At discrete level, operator symmetries must be retained to preserve the

analogous (invariant) properties of the continuous equations: namely, the convective operator is represented by a skew-symmetric matrix, the diffusive operator by a symmetric, negative-definite matrix and the divergence is minus the transpose of the gradient operator. In our opinion, this is the first requirement for reliable DNS and LES simulations. Furthermore, these (large-scale) simulations should run efficiently on the variety of modern HPC systems (CPUs, GPUs, ARM,...) while keeping the code easy to port and maintain.

In this regard, a fully-conservative discretization for collocated unstructured grids was proposed [7]. It exactly preserves the symmetries of the underlying differential operators and is based on only five discrete operators (*i.e.* matrices): the cell-centered and staggered control volumes (diagonal matrices), Ω_c and Ω_s , the face normal vectors, N_s , the cell-to-face interpolation, $\Pi_{c \rightarrow s}$ and the cell-to-face divergence operator, M . Therefore, it constitutes a robust approach that can be easily implemented in existing codes such as OpenFOAM[®] [8]. Then, for the sake of cross-platform portability and optimization, CFD algorithms must rely on a very reduced set of (algebraic) kernels [9] (*e.g.* sparse-matrix vector product, SpMV; dot product; linear combination of vectors). Results showing the benefits of symmetry-preserving discretizations will be presented together with novel methods aiming to keep a good balance between code portability and performance. In particular, results of DNS and LES results of RBC at different Ra will be presented focusing of the feasibility of the developed LES technology to give accurate predictions of the above-explained Nu -vs- Ra scalings. Comparison with DNS results and with the classical GL theory will be conducted.

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