

# An energy-preserving unconditionally stable fractional step method for DNS/LES on collocated grids

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- 1 Motivation: Find an energy-preserving unconditionally stable fractional step method on collocated grids.
- 2 Symmetry-Preserving discretization of NS equations on collocated unstructured grids.
- 3 Conservation of global kinetic energy.
- 4 A stable pressure gradient interpolation for the velocity correction.
- 5 Test case: Turbulent channel flow  $Re_\tau = 395$
- 6 Conclusions.

# 1. Motivation of this work

**Motivation:** Is it possible to find an energy-preserving unconditionally stable fractional step method on collocated grids for any mesh? Instead of changing randomly the numerical schemes, the mesh or...



## 2. Definition of basic collocated operators

Let us suppose we have  $n$  control volumes and  $m$  faces.

Finite volume discretization of incompressible NS equations on an arbitrary collocated mesh

$$\Omega \frac{d\mathbf{u}_c}{dt} + C(\mathbf{u}_s)\mathbf{u}_c = -D\mathbf{u}_c - \Omega G_c \mathbf{p}_c, \quad (1)$$

$$M\mathbf{u}_s = \mathbf{0}_c. \quad (2)$$

- $\mathbf{p}_c = (p_1, \dots, p_n)^T \in \mathbb{R}^n$  is the cell-centered pressure.
- $\mathbf{u}_c = (\mathbf{u}_1, \mathbf{u}_2, \mathbf{u}_3)^T \in \mathbb{R}^{3n}$ , where  $\mathbf{u}_i = ((u_i)_1, \dots, (u_i)_n)^T$  are the vectors containing the velocity components corresponding to the  $x_i$ -spatial direction.
- $\mathbf{u}_s = ((u_s)_1, \dots, (u_s)_m)^T \in \mathbb{R}^m$  is the staggered velocity.
- The velocities are related via the interpolator from cells to faces  
 $\Gamma_{c \rightarrow s} \in \mathbb{R}^{m \times 3n} \implies \mathbf{u}_s = \Gamma_{c \rightarrow s} \mathbf{u}_c.$

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### 3. Conservation of global kinetic energy

#### Global kinetic energy equation

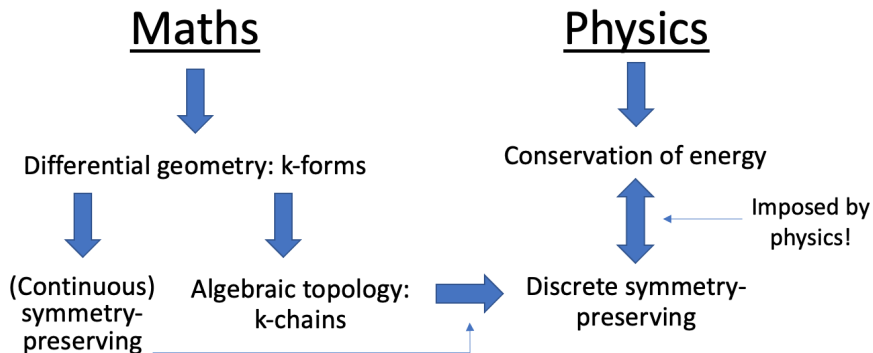
$$\begin{aligned} \frac{d\|\mathbf{u}_c\|^2}{dt} = & -\mathbf{u}_c^T (C(\mathbf{u}_s) + C^T(\mathbf{u}_s))\mathbf{u}_c - \mathbf{u}_c^T (D + D^T)\mathbf{u}_c \\ & - \mathbf{u}_c^T \Omega G_c \mathbf{p}_c - \mathbf{p}_c^T G_c^T \Omega^T \mathbf{u}_c. \end{aligned} \quad (3)$$

In absence of diffusion, that is  $D = 0$ , the global kinetic energy is conserved if:

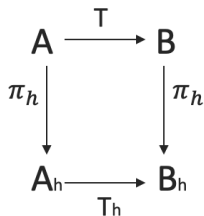
- $C(\mathbf{u}_s) = -C^T(\mathbf{u}_s)$ , i.e, the convective operator should be skew-symmetric.
- $(-\Omega G_c)^T = M \Gamma_{c \rightarrow s}$ , because  $M \mathbf{u}_s = \mathbf{0}_c$ .

**Question:** Can we find a mathematical reason to justify these relations, instead of a physical one?

# Mimicking continuous properties



# Mimicking continuous properties



- $A$  and  $B$  are two vectorial spaces.
- $\pi_h$  is a discretization operator.
- $T$  is a continuous operator.
- $A_h$ ,  $B_h$  and  $T_h$  are the discrete counterparts of  $A$ ,  $B$  and  $T$ , respectively.

**We will require this diagram to be commutative.**

## 4. A stable pressure gradient interpolation

### FSM iterative Poisson equation in collocated meshes

$$L\tilde{p}_c^{n+1} = M_c u_c^n \quad \longrightarrow \quad u_c^{n+1} = u_c^n - G_c \tilde{p}_c^{n+1}, \quad (4)$$

where  $M_c = M\Gamma_{c \rightarrow s}$  and  $G_c = \Gamma_{s \rightarrow c} G = -\Omega_c^{-1} \Gamma_{c \rightarrow s}^T M^T$  are the collocated divergence and the collocated gradient. Developing the correction in  $u_c^n$ :

$$u_c^n = u_c^{n-1} - G_c \tilde{p}_c^n = u_c^{n-2} - G_c \tilde{p}_c^n - G_c \tilde{p}_c^{n-1} = \dots = u_c^p - G_c \sum_{i=1}^n \tilde{p}_c^i \quad (5)$$

So, the accumulated pressure at  $n$  iteration is:

$$p_c^n = \sum_{i=1}^n \tilde{p}_c^i \quad (6)$$

Introducing all this in (7) we obtain:

$$Lp_c^{n+1} = M\Gamma_{c \rightarrow s} u_c^p + (L - L_c) p_c^n. \quad (7)$$

# A stable pressure gradient interpolation

- **Mid-point scheme:**  $\phi_f = \frac{1}{2}(\phi_{c1} + \phi_{c2})$ .
- **Volume weighted scheme:**  $\phi_f = \frac{V_{s1}}{V_{s1}+V_{s2}}\phi_{c1} + \frac{V_{s2}}{V_{s1}+V_{s2}}\phi_{c2}$ .

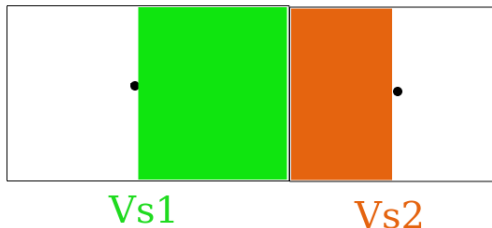


Figure 1: Volume weighted volumes

# A stable pressure gradient interpolation

- Stable solutions  $\rightarrow$  Eigenvalues of  $L - L_c$  negative.
- This can be achieved by using the volume weighted scheme:

$$\Pi_{c \rightarrow s} = \Delta_s^{-1} \Delta_{sc}^T, \quad (8)$$

where  $\Delta_s \in \mathbb{R}^{m \times m}$  is a diagonal matrix containing the projected distances between two adjacent control volumes, and  $\Delta_{sc} \in \mathbb{R}^{m \times n}$  is a matrix containing the projected distances between an adjacent cell node and its corresponding face.

This problem was widely addressed in: *D. Santos, J. Muela, N. Valle, F.X. Trias, On the Interpolation Problem for the Poisson Equation on Collocated Meshes. 14th WCCM-ECCOMAS Congress 2020, DOI: 10.23967/wccm-eccomas.2020.257.*

## 5. Test case: Turbulent channel flow $Re_\tau = 395$

In order to check the stability and accuracy of the method, some tests have been carried out:

- $Re_\tau = 395$ .
- $y^+$  around 1.
- Domain  $5\pi \times 2 \times \pi$ , periodic x and z directions.

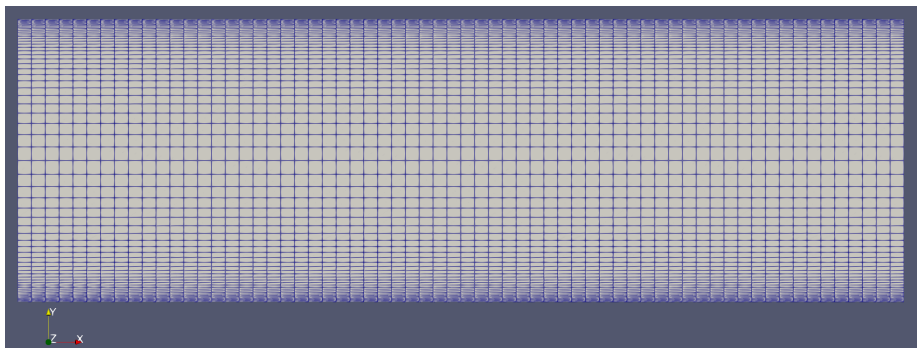
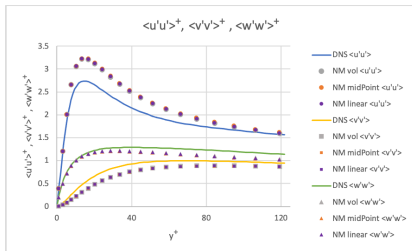
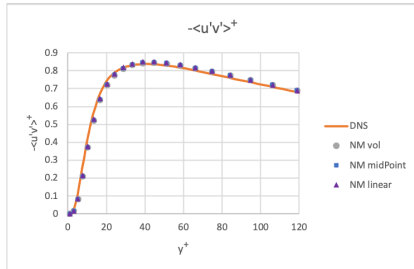
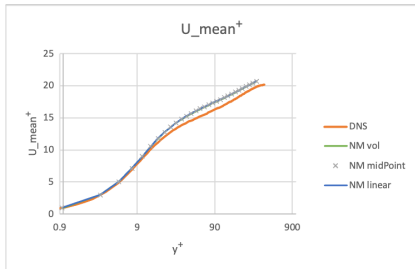


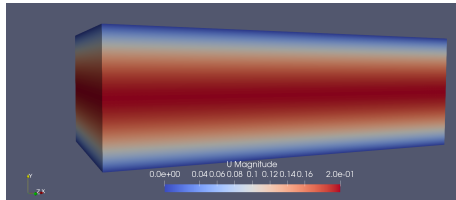
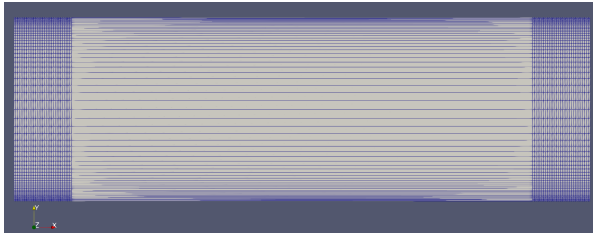
Figure 2:  $64 \times 64 \times 64$  mesh used.

# Test case: Turbulent channel flow $Re_\tau = 395$ .



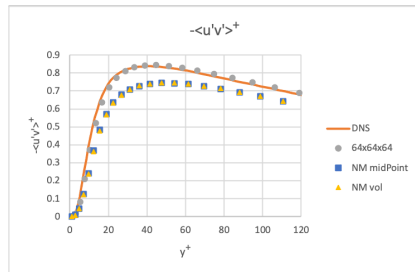
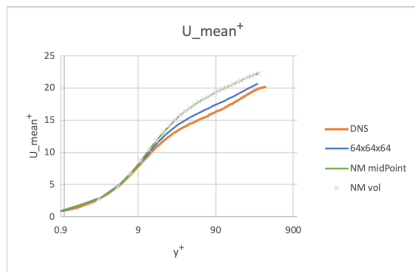
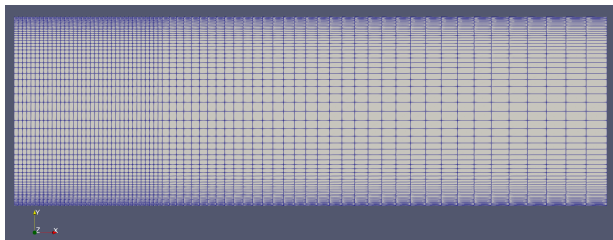


# Turbulent channel flow $Re_{\tau} = 395$ . High distorted mesh.



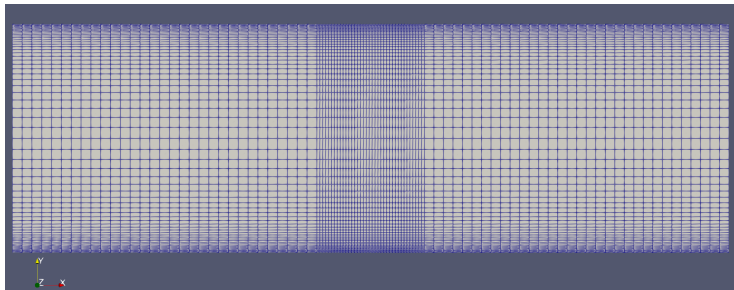
Pressure gradient interpolated using a volume weighted interpolator. MidPoint or linear are blowing up the simulation.

# Turbulent channel flow $Re_\tau = 395$ . Progressive mesh



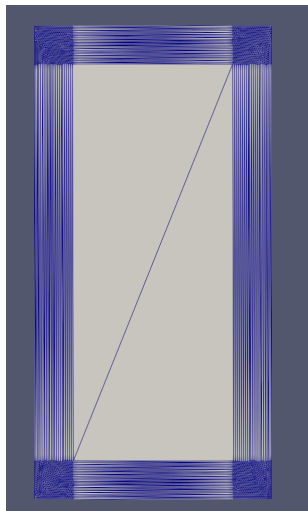
Linear doesn't trigger turbulence.

# Turbulent channel flow $Re_\tau = 395$ . Center refinement

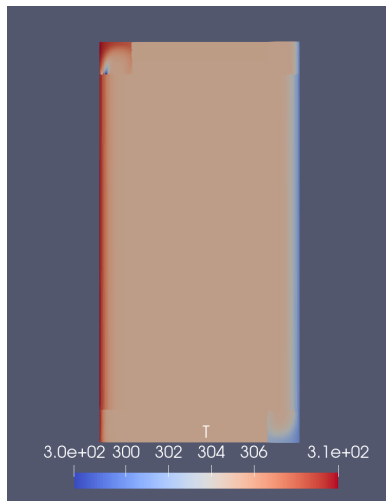


The three interpolators give the same (reasonable) results.

# Preliminary results on unstructured meshes.



Highly distorted mesh.



Temperature distribution.

## 6. Conclusions

### Accuracy and stability conclusions:

- Three interpolators have been tested for the interpolation of the pressure gradient from faces to cells: volume weighted, linear and midPoint interpolations.
- All three seem to have the same accuracy in high quality meshes.
- When distorting the mesh, the first one that loses accuracy and eventually blows up is the linear interpolation.
- The midPoint interpolation seems to be more stable than the linear, but it is still blowing up in highly distorted meshes.
- Volumetric interpolation is an unconditionally stable interpolation.

## General conclusions

- An energy-preserving unconditionally stable fractional step method on collocated grids has been presented.
- There are mathematical reasons beyond physical ones in order to preserve the underlying symmetries of the differential operators.
- The appearance of unphysical velocities is a common problem found in highly distorted meshes, and it solved by means of interpolating the pressure gradient using a volumetric scheme.

## Future work:

- Is there a moment when the midPoint interpolation blows up and the volumetric still gives good results?
- Test accuracy on unstructured meshes.
- Test the accuracy of the solution when varying progressively the distortion of the mesh.