



Centre Tecnològic de Transferència de Calor
UNIVERSITAT POLITÈCNICA DE CATALUNYA



On a proper tensorial subgrid heat flux model

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(ECCOMAS 2020)

Virtual Congress 11-15 January, 2021

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- 2 Modeling the subgrid heat flux
- 3 Building proper models
- 4 Results
- 5 Conclusions

Motivation

Research question:

- Can we find a nonlinear SGS heat flux model with **good physical and numerical properties**, such that we can obtain satisfactory predictions for a turbulent Rayleigh-Bénard convection?

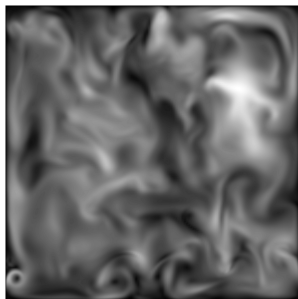
DNS of an air-filled Rayleigh-Bénard convection at $Ra = 10^{10}$

¹F.Dabbagh, F.X.Trias, A.Gorobets, A.Oliva. *On the evolution of flow topology in turbulent Rayleigh-Bénard convection*, **Physics of Fluids**, 28:115105, 2016.

Motivation

Air-filled RB: $Pr = 0.7$

$$Ra = 10^8$$



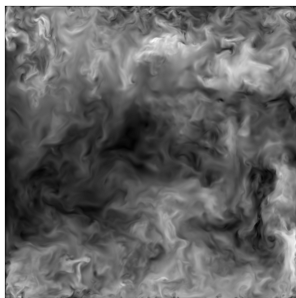
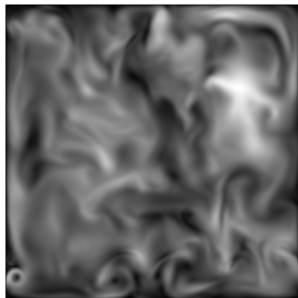
²F.Dabbagh, F.X.Trias, A.Gorobets, A.Oliva. *Flow topology dynamics in a 3D phase space for turbulent Rayleigh-Bénard convection*, **Phys.Rev.Fluids**, 5:024603, 2020.

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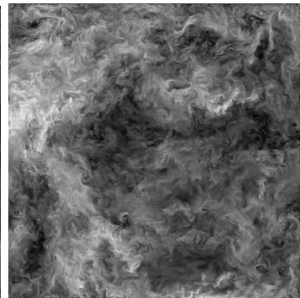
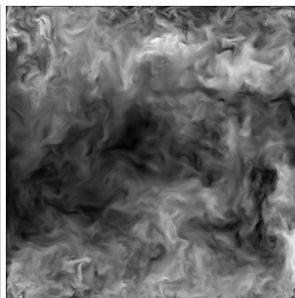
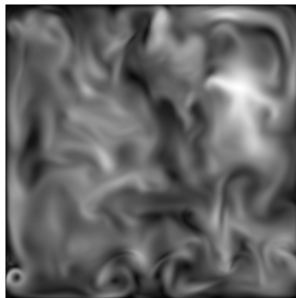
Motivation

Air-filled RB: $Pr = 0.7$

$$Ra = 10^8$$

$$Ra = 10^{10}$$

$$Ra = 10^{11}$$



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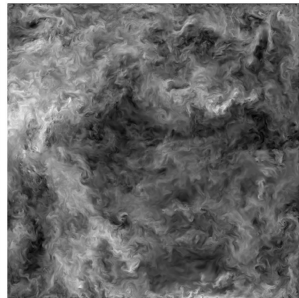
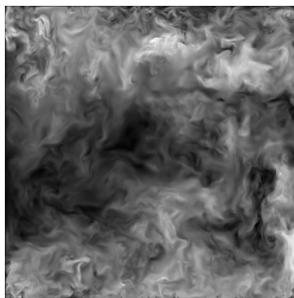
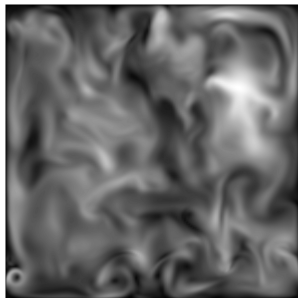
Motivation

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$208 \times 208 \times 400$

17.5M

$768 \times 768 \times 1024$

607M

$1662 \times 1662 \times 2048$

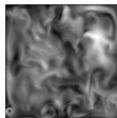
5600M

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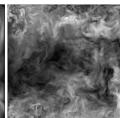
Motivation

DNS: $208 \times 208 \times 400$

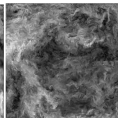
$Ra = 10^8$



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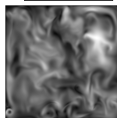


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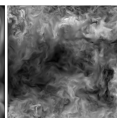
DNS: $208 \times 208 \times 400$

LES: $80 \times 80 \times 120$

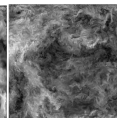
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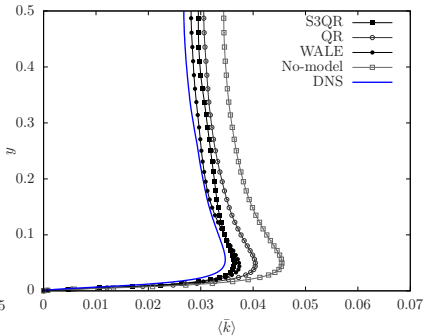
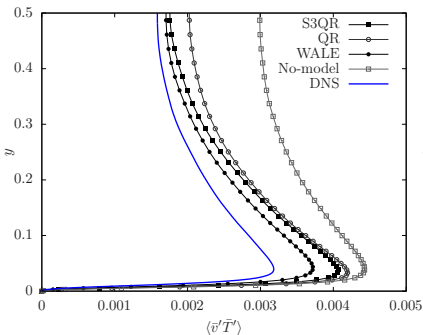
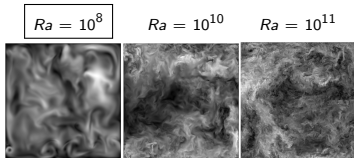
$Ra = 10^{11}$



Motivation

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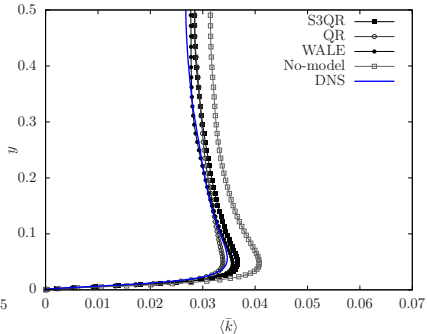
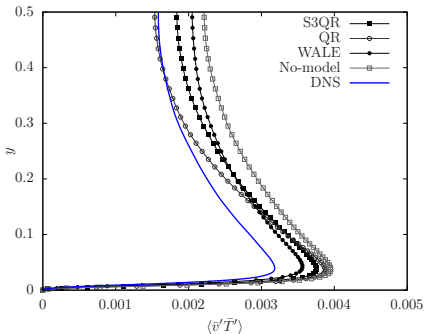
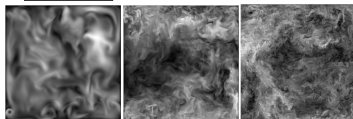
DNS: $208 \times 208 \times 400$

LES: $110 \times 110 \times 168$

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How to model the subgrid heat flux in LES?

$$\partial_t \bar{u} + (\bar{u} \cdot \nabla) \bar{u} = \nu \nabla^2 \bar{u} - \nabla \bar{p} - \nabla \cdot \tau(\bar{u}) ; \quad \nabla \cdot \bar{u} = 0$$

eddy-viscosity $\rightarrow \tau(\bar{u}) = -2\nu_t S(\bar{u})$

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$Pr_t?$

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$$G \equiv \nabla \bar{u} \quad \mathbf{q} = -\frac{\delta^2}{12} G \nabla \bar{T} + \mathcal{O}(\delta^4)$$

A priori alignment trends³

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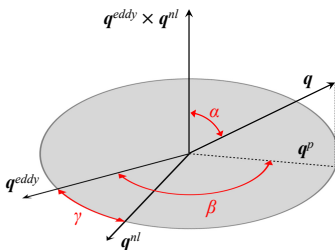
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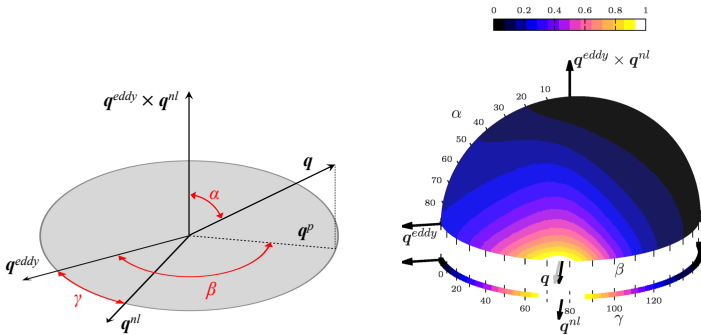


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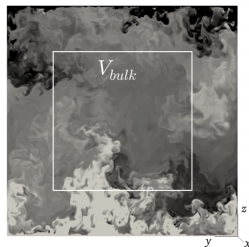
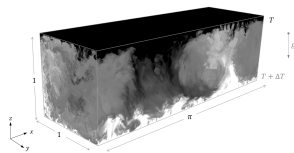
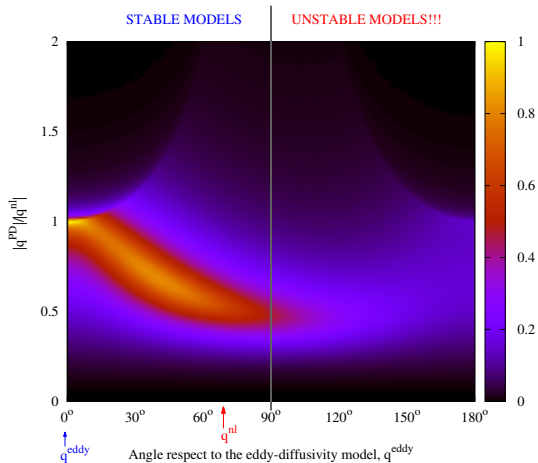
Peng&Davidson⁴ $\rightarrow \mathbf{q} \approx -\frac{\delta^2}{12} S \nabla \bar{T} \quad (\equiv q^{PD})$

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A priori alignment trends

$$q^{nl} \equiv -\frac{\delta^2}{12} G \nabla \bar{T}$$

$$q^{PD} \equiv -\frac{\delta^2}{12} S \nabla \bar{T}$$



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mixed model $\longrightarrow \mathbf{q} \approx q^{nl} + \sigma q^{eddy} \quad (\equiv q^{mix})$

⁵B.J.Daly and F.H.Harlow. **Physics of Fluids**, 13:2634, 1970.

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$$\text{Daly\&Harlow}^5 \longrightarrow \mathbf{q} \approx -T_{SGS} \frac{\delta^2}{12} \mathbf{GG}^T \nabla \bar{T} \quad (\equiv \mathbf{q}^{DH})$$

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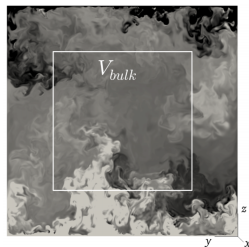
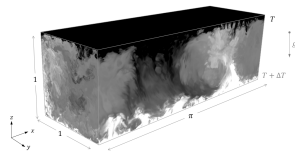
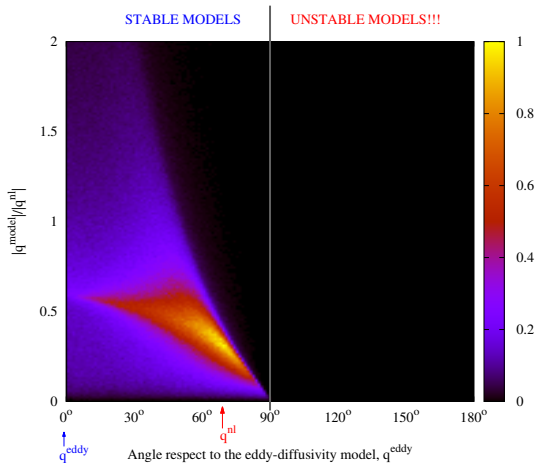
$$\mathcal{T}_{SGS} = 1/|S|$$

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A priori alignment trends

$$q^{nl} \equiv -\frac{\delta^2}{12} G \nabla \bar{T}$$

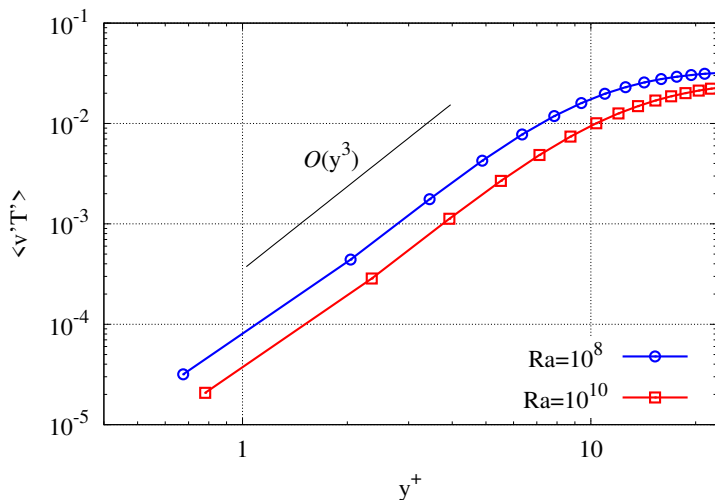
$$q^{DH} \equiv -\mathcal{T}_{SGS} \frac{\delta^2}{12} GG^T \nabla \bar{T}$$



What about near-wall scaling?

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⇒ Answer: it should be $\mathcal{O}(y^3)$



Near-wall scaling for DH model?

$$q^{DH} \equiv -\mathcal{T}_{SGS} \frac{\delta^2}{12} GG^T \nabla \bar{T}; \quad \mathcal{T}_{SGS} = 1/|S|$$

Near-wall scaling for DH model?

⇒ Answer: it is $\mathcal{O}(y^1)$ instead of $\mathcal{O}(y^3)$

$$q^{DH} \equiv -\mathcal{T}_{SGS} \frac{\delta^2}{12} GG^T \nabla \bar{T}; \quad \mathcal{T}_{SGS} = 1/|S|$$

$$u = ay + \mathcal{O}(y^2); \quad v = by^2 + \mathcal{O}(y^3); \quad w = cy + \mathcal{O}(y^2); \quad T = dy + \mathcal{O}(y^2)$$

$$G = \begin{pmatrix} y & 1 & y \\ y^2 & y & y^2 \\ y & 1 & y \end{pmatrix}; \quad \nabla \bar{T} = \begin{pmatrix} y \\ 1 \\ y \end{pmatrix} \Rightarrow GG^T \nabla \bar{T} = \begin{pmatrix} y \\ y^2 \\ y \end{pmatrix} = \mathcal{O}(y^1)$$

$$\mathcal{T}_{SGS} = 1/|S| = \mathcal{O}(y^0)$$

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Idea: build a \mathcal{T}_{SGS} with the proper $\mathcal{O}(y^2)$ scaling!!!

Building proper models for the subgrid heat flux

Let us consider models that are based on the invariants of the tensor GG^T

$$\mathbf{q} \approx -C_M \left(P_{GG^T}^p Q_{GG^T}^q R_{GG^T}^r \right) \frac{\delta^2}{12} GG^T \nabla \bar{T} \quad (\equiv \mathbf{q}^{S2})$$

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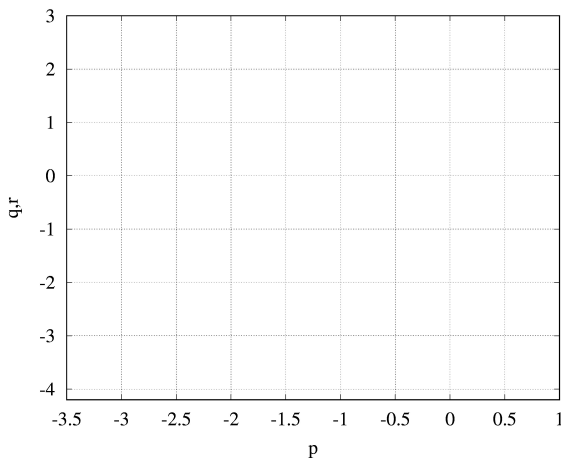
	P_{GG^T}	Q_{GG^T}	R_{GG^T}
Formula	$2(Q_\Omega - Q_S)$	$V^2 + Q_G^2$	R_G^2
Wall-behavior	$\mathcal{O}(y^0)$	$\mathcal{O}(y^2)$	$\mathcal{O}(y^6)$
Units	$[T^{-2}]$	$[T^{-4}]$	$[T^{-6}]$

$$-6r - 4q - 2p = 1 \quad [T]; \quad 6r + 2q = s,$$

where s is the slope for the asymptotic near-wall behavior, *i.e.* $\mathcal{O}(y^s)$.

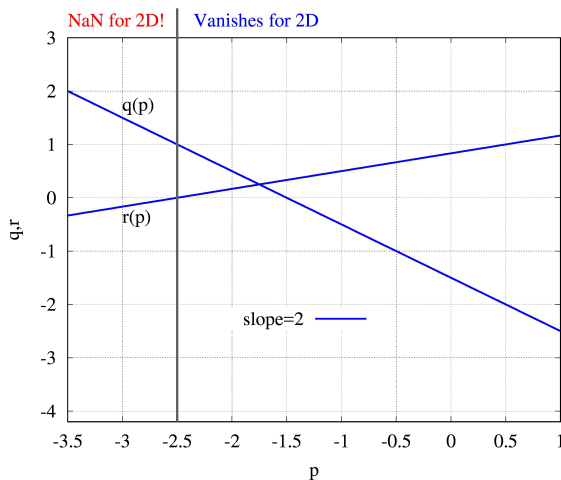
Building proper models for the subgrid heat flux

Solutions: $q(p, s) = -(1 + s)/2 - p$ and $r(p, s) = (2s + 1)/6 + p/3$



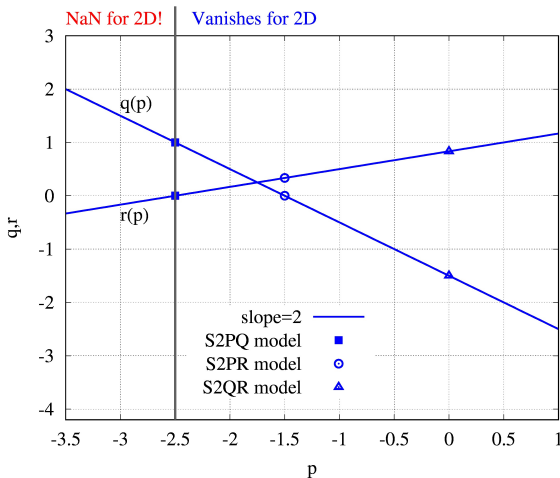
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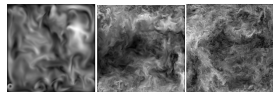
A priori analysis in the bulk

Estimation of the model constant, C_{s2pqr}

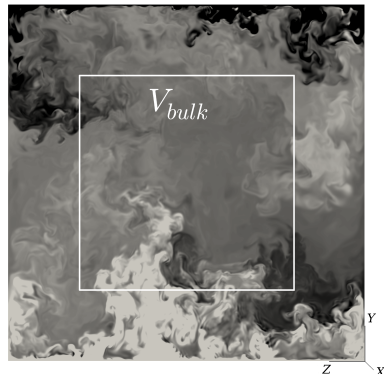
Playing with exponent $p...$

$$q^{S2PQR} \approx -C_{s2pqr} \left(P_{GGT}^p Q_{GGT}^{-(p+3/2)} R_{GGT}^{(2p+5)/6} \right) \frac{\delta^2}{12} GG^T \nabla \bar{T}$$

$Ra = 10^8$ $Ra = 10^{10}$ $Ra = 10^{11}$



$$\frac{\langle |q^{S2PQR}| \rangle_{bulk}}{\langle |q| \rangle_{bulk}} = 1 \longrightarrow C_{s2pqr}$$



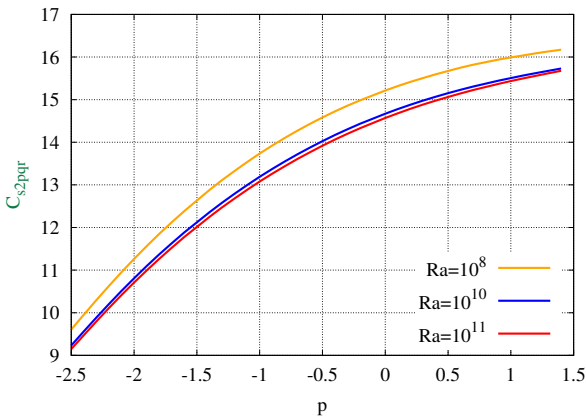
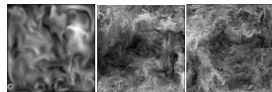
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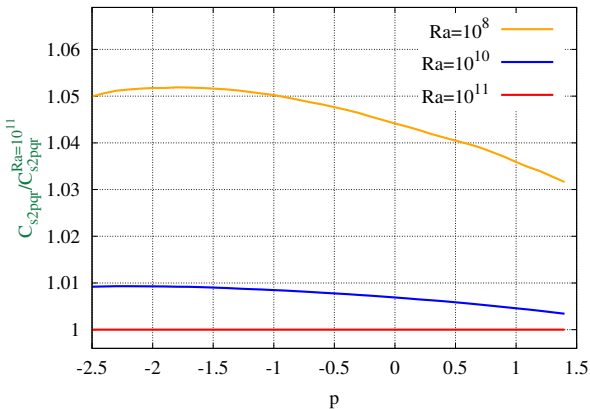
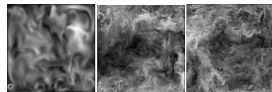
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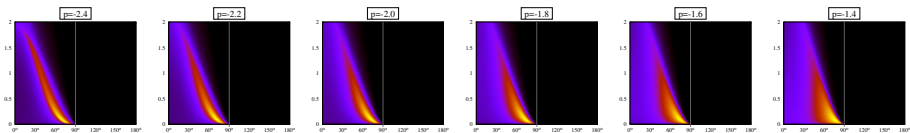
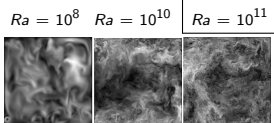


A priori analysis in the bulk

Alignment trends...

Playing with exponent $p...$

$$q^{S2PQR} \approx -C_s 2^{pqr} \left(P_{GGT}^p Q_{GGT}^{-(p+3/2)} R_{GGT}^{(2p+5)/6} \right) \frac{\delta^2}{12} GG^T \nabla \bar{T}$$

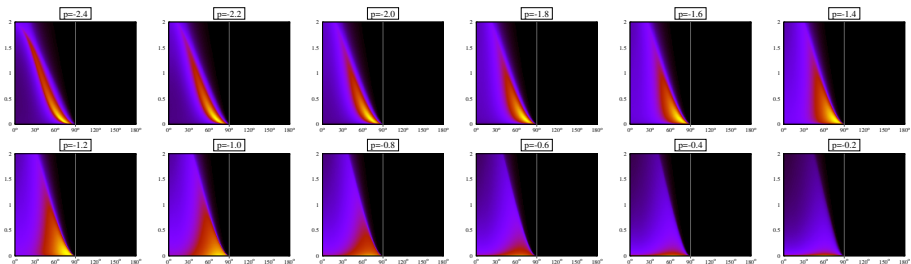
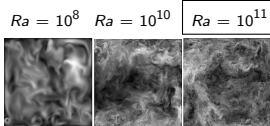


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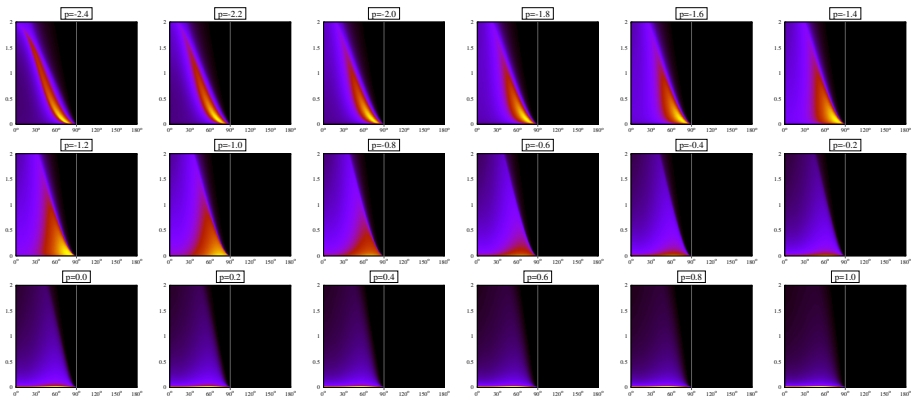
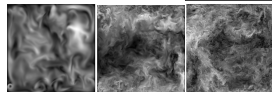
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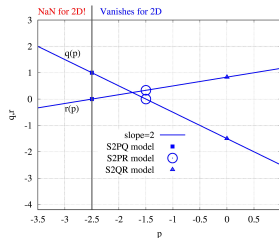
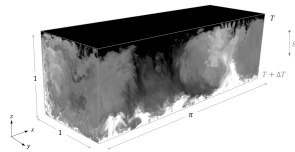
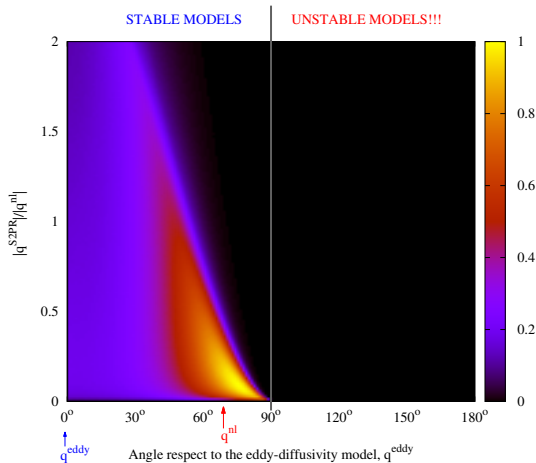
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A priori alignment trends of S2PR in the bulk

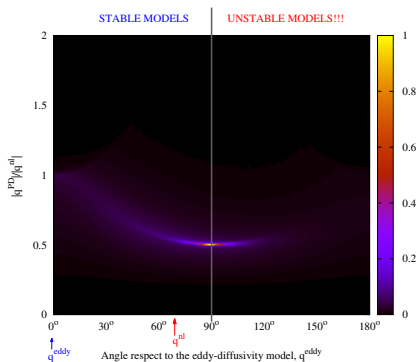
$$q^{nl} \equiv -\frac{\delta^2}{12} G \nabla \bar{T}$$

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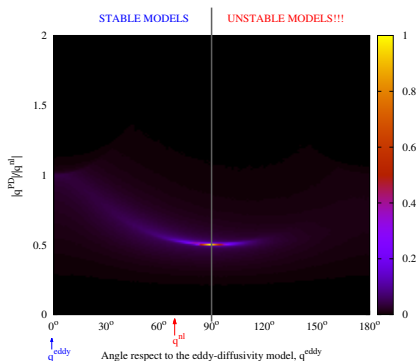
A priori alignment trends of S2PR in the near-wall region

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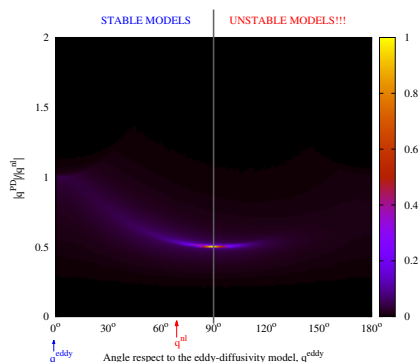
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$$\cos \beta^{model} = \frac{q^{eddy} \cdot q^{model}}{|q^{eddy}| |q^{model}|}$$

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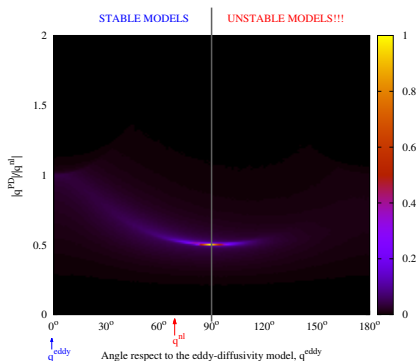
$$\lim_{y \rightarrow 0^+} \frac{|q^{PD}|}{|q^{nl}|} = \lim_{y \rightarrow 0^+} \frac{|S \nabla \bar{T}|}{|G \nabla \bar{T}|} = \frac{1}{2}$$

$$\cos \beta^{model} = \frac{q^{eddy} \cdot q^{model}}{|q^{eddy}| |q^{model}|}$$

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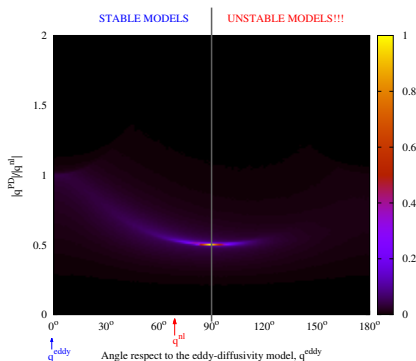


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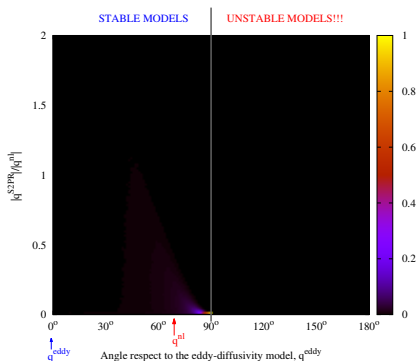
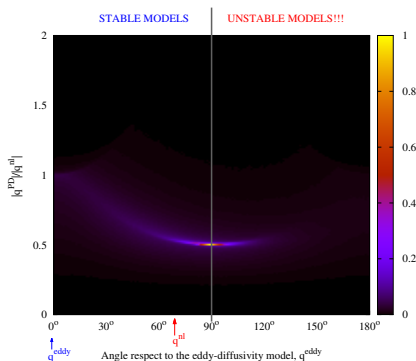
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A posteriori results?

$$\partial_t \bar{u} + (\bar{u} \cdot \nabla) \bar{u} = \nu \nabla^2 \bar{u} - \nabla \bar{p} + \bar{f} - \nabla \cdot \tau(\bar{u}) ; \quad \nabla \cdot \bar{u} = 0$$

$$\text{eddy-viscosity} \longrightarrow \tau(\bar{u}) = -2\nu_t S(\bar{u})$$

$$\nu_t \approx (C_m \delta)^2 D_m(\bar{u})$$

$$\partial_t \bar{T} + (\bar{u} \cdot \nabla) \bar{T} = \alpha \nabla^2 \bar{T} - \nabla \cdot \mathbf{q} \quad \text{where} \quad \mathbf{q} = \overline{uT} - \bar{u}\bar{T}$$

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⚠ But first we need to answer the following **research question**:

- Are **eddy-viscosity models** for momentum able to provide satisfactory results for turbulent Rayleigh-Bénard convection?

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Idea: let's do an LES for momentum and a DNS for temperature!

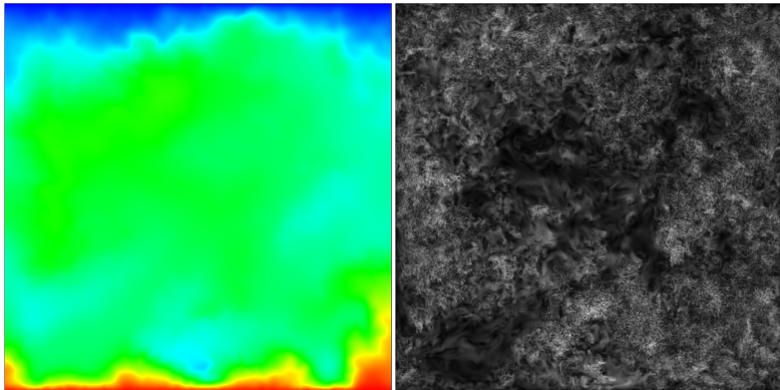
DNS at very low Pr number

Why? scale separation grows as $\eta_K/\eta_T = Pr^{3/4}$.

η_T : Obukhov-Corrsin scale; η_K : Kolmogorov scale

DNS at very low Pr number

Why? scale separation grows as $\eta_K/\eta_T = Pr^{3/4}$. Here: $\eta_T \approx 53.2\eta_K$
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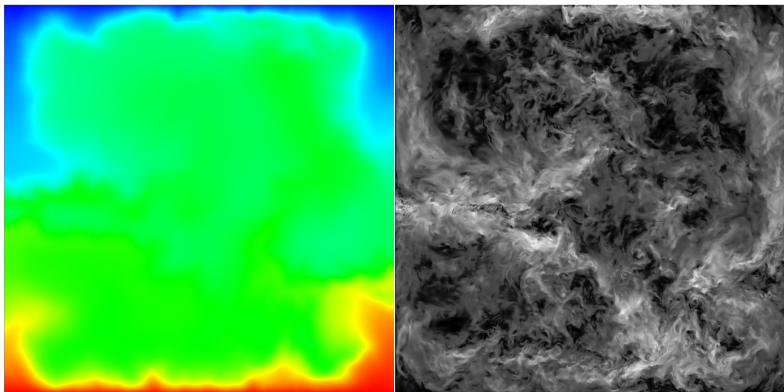


DNS of a RB at $Ra = 7.14 \times 10^6$ and $Pr = 0.005$ (liquid sodium)
 $488 \times 488 \times 1280 \approx \mathbf{305M}$

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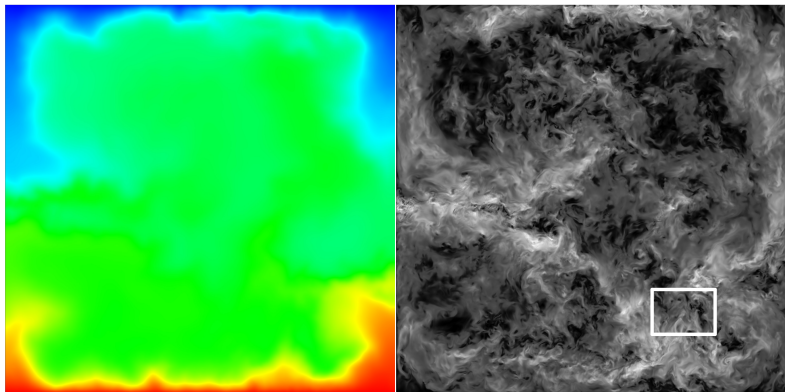


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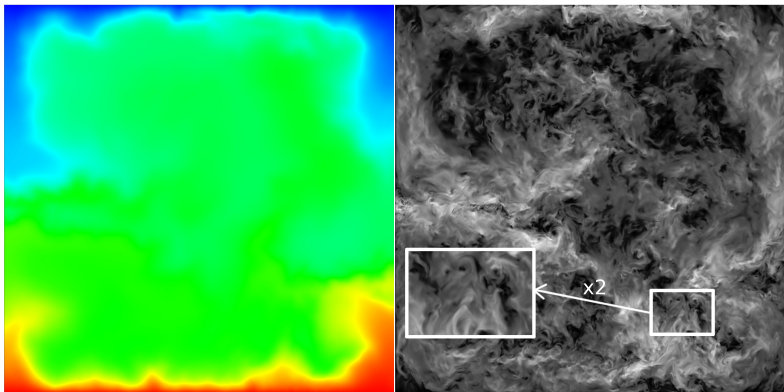


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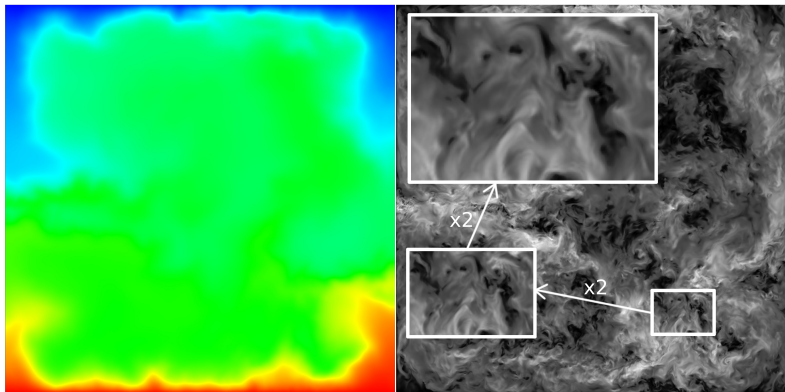


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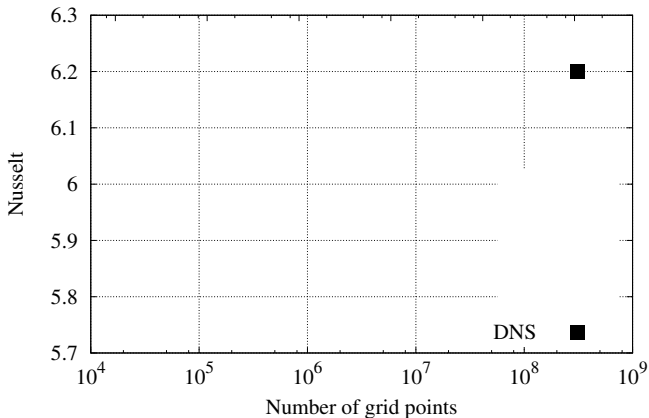
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LES⁶ results at very low Pr number

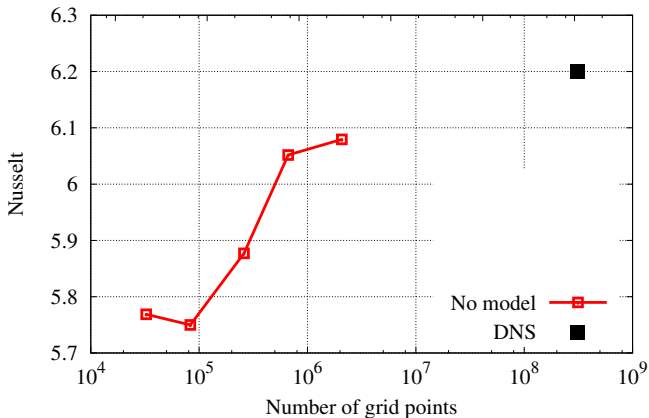
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⁶F.X.Trias, D.Folch, A.Gorobets, A.Oliva. *Building proper invariants for eddy-viscosity subgrid-scale models*, **Physics of Fluids**, 27: 065103, 2015.

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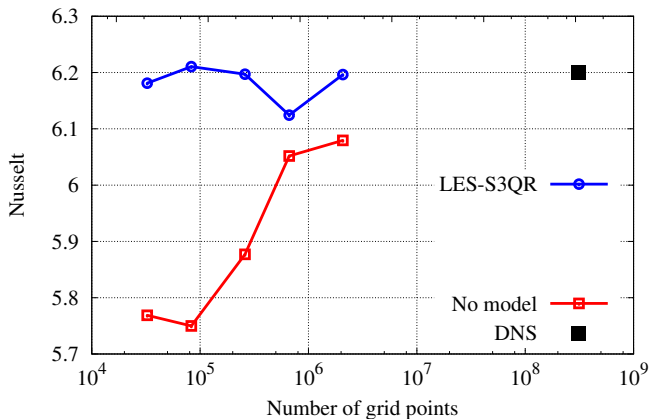
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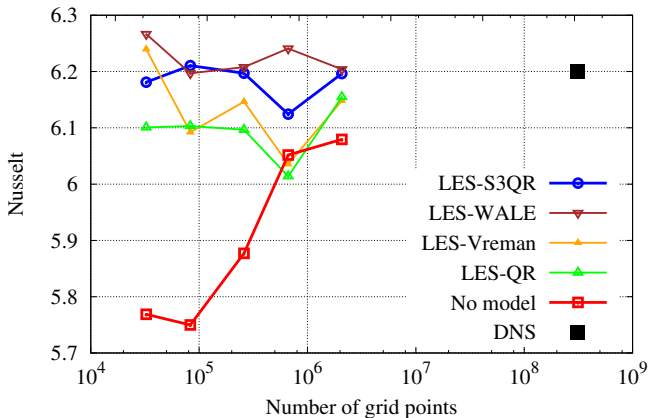
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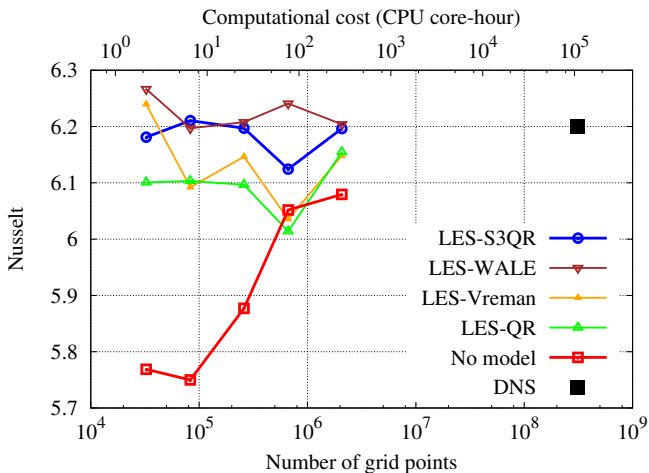
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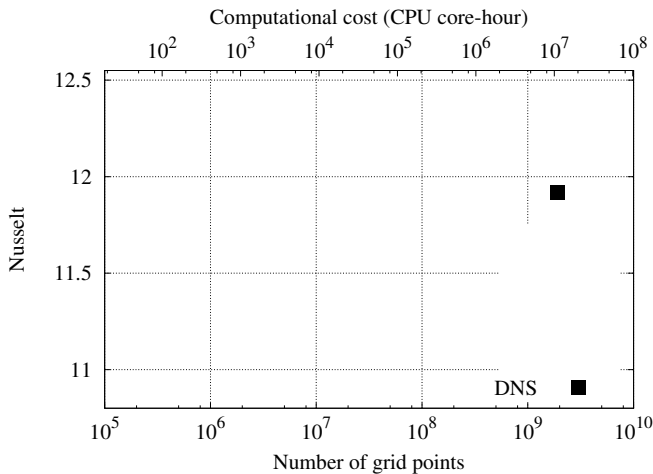
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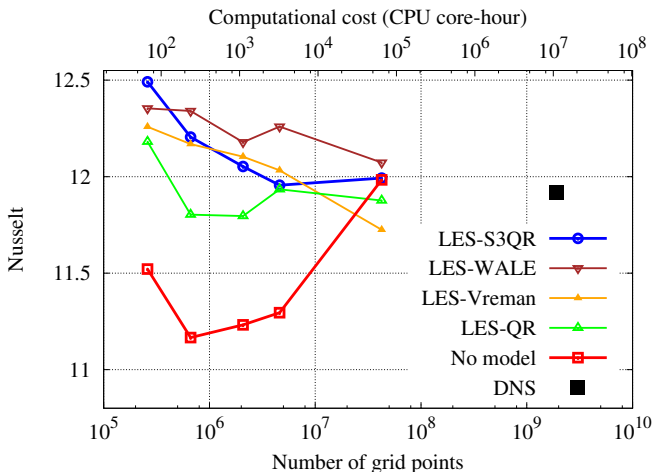
LES results at very low Pr number

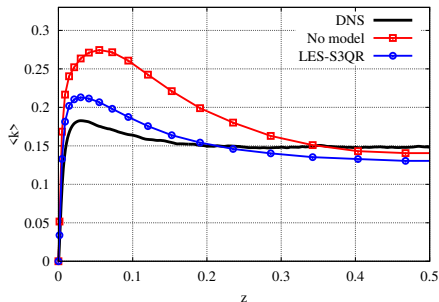
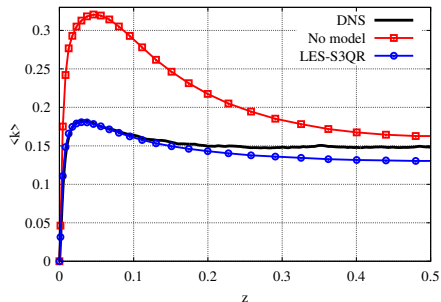
RB at $Ra = 7.14 \times 10^7$ and $Pr = 0.005$ (DNS $\rightarrow 966 \times 966 \times 2048 \approx 1911M$)



LES results at very low Pr number

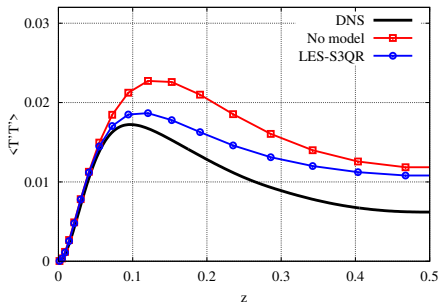
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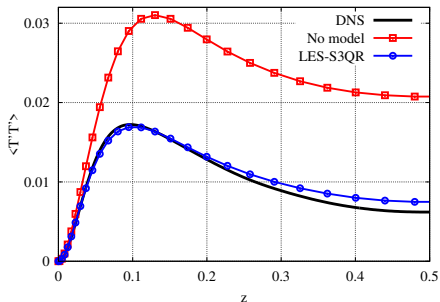
LES results at very low Pr numberRB at $Ra = 7.14 \times 10^6$ and $Pr = 0.005$ (DNS $\rightarrow 488 \times 488 \times 1280 \approx 305M$) $64 \times 32 \times 32$  $96 \times 52 \times 52$

LES results at very low Pr number

RB at $Ra = 7.14 \times 10^6$ and $Pr = 0.005$ (DNS $\rightarrow 488 \times 488 \times 1280 \approx 305M$)



$64 \times 32 \times 32$



$96 \times 52 \times 52$

Concluding remarks

- A new tensor-diffusivity model has been proposed⁷

$$q^{s2PR} \equiv -C_{s2pr} P_{GG^T}^{-3/2} R_{GG^T}^{1/3} \frac{\delta^2}{12} GG^T \nabla \bar{T}$$

- Locally defined, unconditionally stable and vanishes for 2D flows ✓
- Good *a priori* alignment trends and proper near-wall scaling ✓
- Eddy-viscosity models work for RB ✓

⁷F.X.Trias, F.Dabbagh, A.Gorobets, C.Oliet. *On a proper tensor-diffusivity model for LES of buoyancy-driven turbulence*, **Flow Turbul Combust**, 105:393-414, 2020.

⁸F.X.Trias, A.Gorobets, M.Silvis, R.Verstappen, A.Oliva. *A new subgrid characteristic length for turbulence simulations on anisotropic grids*, **Phys.Fluids**, 26:115109, 2017.

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Future research:

- *A posteriori* tests using q^{s2PR} for RB

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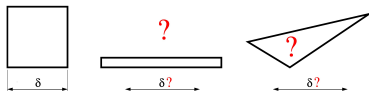
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Future research:

- *A posteriori* tests using q^{s2PR} for RB
- How δ should be defined for highly anisotropic grids⁸?



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Thank you for your virtual
attendance