

# ON THE CONDITIONS FOR A STABLE PROJECTION METHOD ON COLLOCATED UNSTRUCTURED GRIDS

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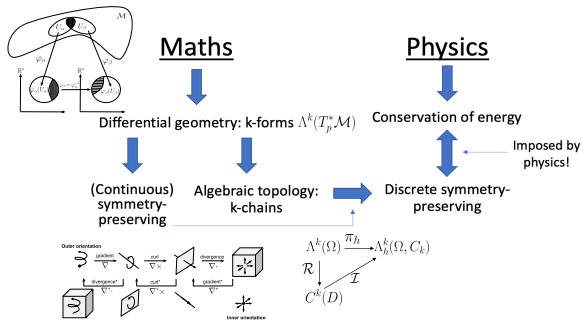
03-07 June 2024, ECCOMAS 2024, Lisbon, Portugal



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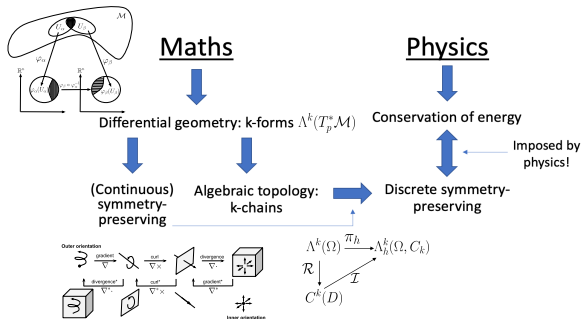
- 1 Symmetry-Preserving unconditionally stable discretization of NS equations on collocated unstructured grids.
- 2 Conservation of global kinetic energy
- 3 Summary and conclusions

# Motivation



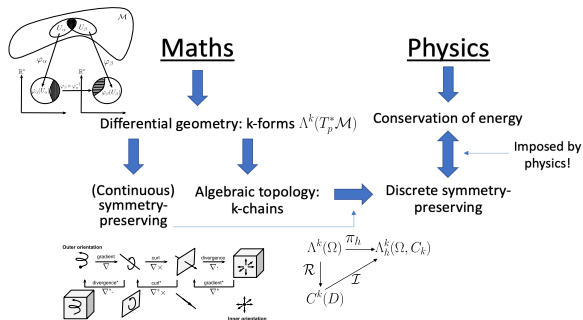
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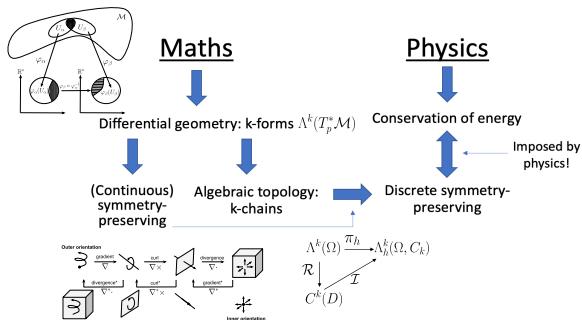
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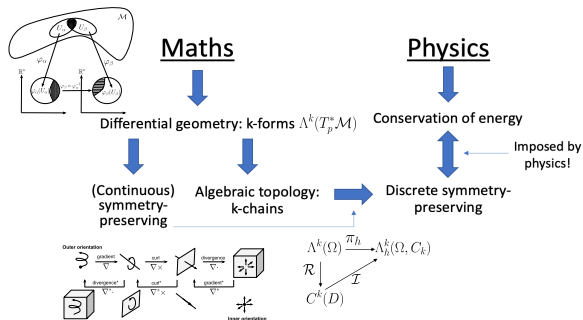
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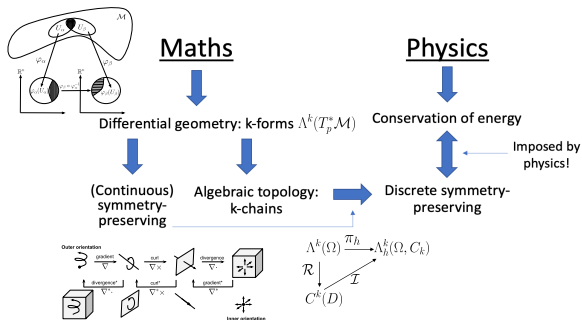
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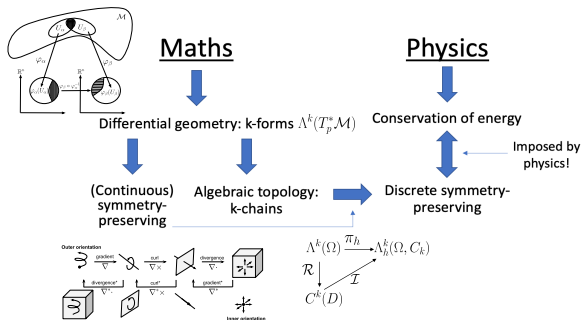


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- Easily portable (to other codes, platforms...)
- Unravel why using midpoint in the convective term.
- Explain why the staggered metric is used at faces.

# 1. Definition of basic collocated operators

Let us suppose we have  $n$  control volumes and  $m$  faces.

Finite volume discretization of incompressible NS equations on an arbitrary collocated mesh

$$\Omega \frac{d\mathbf{u}_c}{dt} + C(\mathbf{u}_s)\mathbf{u}_c = -D\mathbf{u}_c - \Omega G_c \mathbf{p}_c, \quad (1)$$

$$M\mathbf{u}_s = \mathbf{0}_c. \quad (2)$$

- $\mathbf{p}_c = (p_1, \dots, p_n)^T \in \mathbb{R}^n$  is the cell-centered pressure.
- $\mathbf{u}_c = (\mathbf{u}_1, \mathbf{u}_2, \mathbf{u}_3)^T \in \mathbb{R}^{3n}$ , where  $\mathbf{u}_i = ((u_i)_1, \dots, (u_i)_n)^T$  are the vectors containing the velocity components corresponding to the  $x_i$ -spatial direction.
- $\mathbf{u}_s = ((u_s)_1, \dots, (u_s)_m)^T \in \mathbb{R}^m$  is the staggered velocity.
- The velocities are related via the interpolator from cells to faces  
 $\Gamma_{c \rightarrow s} \in \mathbb{R}^{m \times 3n} \implies \mathbf{u}_s = \Gamma_{c \rightarrow s} \mathbf{u}_c.$

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# Definition of basic collocated operators

The (3D) interpolator from cells to faces can be constructed as follows:

$$\Gamma_{c \rightarrow s} = N\Pi, \quad (3)$$

where

- $N = (N_{s,x} N_{s,y} N_{s,z}) \in \mathbb{R}^{m \times 3m}$  where  $N_{s,x}, N_{s,y}, N_{s,z} \in \mathbb{R}^{m \times m}$  are diagonal matrices containing the  $x_i$  spatial component of the face normal vectors.
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- $\Omega_c \in \mathbb{R}^{n \times n}$  is a diagonal matrix with the cell-centered volumes  
 $\implies \Omega = I_3 \otimes \Omega_c$ .
- $C_c(\mathbf{u}_s) \in \mathbb{R}^{n \times n}$  is the cell-centered convective operator for a discrete scalar field  
 $\implies C(\mathbf{u}_s) = I_3 \otimes C_c(\mathbf{u}_s)$ .
- $D_c \in \mathbb{R}^{n \times n}$  is the cell-centered diffusive operator for a discrete scalar field  
 $\implies D = I_3 \otimes D_c$ .

Finally,

- $G_c \in \mathbb{R}^{3n \times n}$  represents the discrete collocated gradient.
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# Mimicking continuous properties

**Mimicking Hilbert adjointness in  $L^2$  inner product**  $\rightarrow$

$$G = -\Omega_s^{-1} M^T,$$

**Mimicking continuous Laplacian**  $\rightarrow$

$$L = MG = -M\Omega_s^{-1} M^T,$$

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$$\Gamma_{s \rightarrow c} = \Omega^{-1} \Gamma_{c \rightarrow s}^T \Omega_s. \quad (4)$$

where  $G$  is the center-to-face staggered gradient,  $L$  is the Laplacian operator,  $L_c$  is the collocated-Laplacian operator and  $\Gamma_{s \rightarrow c}$  is the face-to-cell interpolator.

For more information about Symmetry-Preserving discretization consult: *F.X. Trias, O. Lehmkuhl, A. Oliva, C.D. Perez-Segarra, and R.W.C.P. Verstappen. Symmetry-preserving discretization of Navier-Stokes equations on collocated unstructured meshes. Journal of Computational Physics, 258:246–267, 2014.*

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# Reduction of 2-forms and divergence operator

- Operator Reduction  $\mathcal{R} \rightarrow$  FV discretization.
- 2-form  $\mathbf{u}_f$  can be "understood" as a vector at the faces.

$$\mathcal{R}\mathbf{u}_f \equiv \int_f \mathbf{u}_f \cdot \hat{n} dS = \mathbf{u}_s S_f \quad (5)$$

Divergence operator:

$$\begin{aligned} \mathcal{R}(\nabla \cdot \mathbf{u}) &= \int_V \nabla \cdot \mathbf{u} dV = \int_{\partial V} \mathbf{u}_f \cdot \hat{n} dS = \sum_f \mathbf{u}_s S_f \rightarrow \\ &\rightarrow M = \mathbb{T}_{c \rightarrow s} S_f \rightarrow M\mathbf{u}_s = \delta \mathcal{R}\mathbf{u}_f \end{aligned} \quad (6)$$

# Skew-Symmetry of the convective operator

The convective operator in a FV context is constructed as:

$$C(\mathbf{u}_s)\mathbf{u}_c = M(\text{diag}(\mathbf{u}_s)\Pi_{c \rightarrow s}^{\text{mid}}\mathbf{u}_c) \quad (7)$$

From a DEC point of view, it is possible to decompose the continuous convective for 3-forms as follows:

$$\text{Conv}(\phi) = d\mathbf{u}\phi = d \star (\star\phi \wedge \mathbf{u}^b) \quad (8)$$

Assuming a local stencil, the construction of the discrete wedge product between 0-forms and 1-forms induces the coefficient  $1/2 \rightarrow$

**It is not an interpolation, but the construction of the wedge.**

Finally, this wedge combined with  $M = \mathbb{T}_{c \rightarrow s} S_f$  is automatically skew-symmetric.

## 2. Conservation of global kinetic energy

### Global kinetic energy equation

$$\begin{aligned} \frac{d\|\mathbf{u}_c\|^2}{dt} = & -\mathbf{u}_c^T (C(\mathbf{u}_s) + C^T(\mathbf{u}_s))\mathbf{u}_c - \mathbf{u}_c^T (D + D^T)\mathbf{u}_c \\ & - \mathbf{u}_c^T \Omega G_c \mathbf{p}_c - \mathbf{p}_c^T G_c^T \Omega^T \mathbf{u}_c. \end{aligned} \quad (9)$$

In the absence of diffusion, that is,  $D = 0$ , the global kinetic energy is conserved if:

- $C(\mathbf{u}_s) = -C^T(\mathbf{u}_s)$ , i.e, the convective operator should be skew-symmetric.
- $(-\Omega G_c)^T = M\Gamma_{c \rightarrow s}$ , because  $M\mathbf{u}_s = \mathbf{0}_c$ .

## Global kinetic energy equation with skew-symmetric convective operator

$$\frac{d\|\mathbf{u}_c\|^2}{dt} = -\mathbf{u}_c^T(D + D^T)\mathbf{u}_c - \mathbf{u}_c^T\Omega G_c \mathbf{p}_c - \mathbf{p}_c^T G_c^T \Omega^T \mathbf{u}_c.$$

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- $(-\Omega G_c)^T = M\Gamma_{c \rightarrow s}$ , because  $M\mathbf{u}_s = \mathbf{0}_c$  (But this relation is exact ONLY in staggered configurations!).

In collocated framework, we either solve:

$$M\mathbf{u}_s = 0 \rightarrow Lp_c = M\Gamma_{c \rightarrow s}\mathbf{u}_c^p \rightarrow \text{Kinetic Energy Error} \quad (10)$$

$$M_c\mathbf{u}_c = 0 \rightarrow L_cp_c = M\Gamma_{c \rightarrow s}\mathbf{u}_c^p \rightarrow \text{Checkerboard} \quad (11)$$

## Global kinetic energy equation with skew-symmetric convective operator

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In collocated framework and explicit time integration, the (artificial) kinetic energy added is given by:

$$-\mathbf{p}_c^T G_c^T \Omega^T \mathbf{u}_c = \mathbf{p}_c^T (L - L_c) \mathbf{p}_c \Delta t \quad (12)$$

# Volume-weighted interpolator

- The volume-weighted interpolator can be constructed as:

$$\Pi_{c \rightarrow s} = \Delta_s^{-1} \Delta_{sc}^T, \quad (13)$$

where  $\Delta_s \in \mathbb{R}^{m \times m}$  is a diagonal matrix containing the projected distances between two adjacent control volumes, and  $\Delta_{sc} \in \mathbb{R}^{m \times n}$  contains the projected distances between an adjacent cell node and its corresponding face.

**Volume-weighted interpolation:**  $\phi_f = \frac{\tilde{V}_{1,f}}{\tilde{V}_{1,f} + \tilde{V}_{2,f}} \phi_{c1} + \frac{\tilde{V}_{2,f}}{\tilde{V}_{1,f} + \tilde{V}_{2,f}} \phi_{c2}$ .

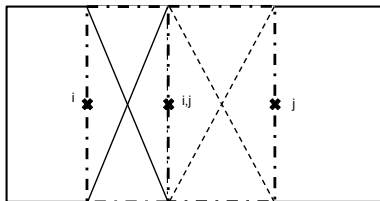


Figure 1: Volume-weighted volumes

# Volume-weighted interpolator

**Volume-weighted interpolation:**  $\phi_f = \frac{d_{1,f}}{d_{1,f}+d_{2,f}}\phi_{c1} + \frac{d_{2,f}}{d_{1,f}+d_{2,f}}\phi_{c2}$ .

Momentum is conserved when interpolated by the volume-weighted:

$$\begin{aligned}(\mathbf{u}_c, \mathbf{1}_c)_\Omega &= (\mathbf{u}_s, \mathbf{1}_s)_{\Omega_s} \rightarrow \\ \mathbf{u}_x^1 \tilde{V}_1 + \mathbf{u}_x^2 \tilde{V}_2 &= (a_1 \mathbf{u}_x^1 + a_2 \mathbf{u}_x^2) A_f = (d_1 \mathbf{u}_x^1 + d_2 \mathbf{u}_x^2) A_f \rightarrow \\ \mathbf{u}_x^1 \tilde{V}_1 + \mathbf{u}_x^2 \tilde{V}_2 &= \left( \frac{d_1}{d_1 + d_2} \mathbf{u}_x^1 + \frac{d_2}{d_1 + d_2} \mathbf{u}_x^2 \right) V_s\end{aligned}$$

And the staggered metric is induced automatically to the faces!

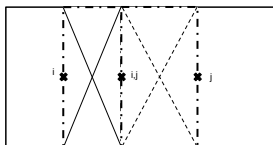


Figure 2: Volume-weighted volumes



# Conservation of global kinetic energy

Stable solutions  $\rightarrow$  Eigenvalues of  $L - L_c$  negative.

## Theorem

Assumptions:

# Conservation of global kinetic energy

Stable solutions  $\rightarrow$  Eigenvalues of  $L - L_c$  negative.

## Theorem

Assumptions:

- Our projection method adds a kinetic energy error of the form  $\mathbf{p}_c^T (L - L_c) \mathbf{p}_c^T$  (such as the FSM or PISO).

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## Corollary 1

Hexahedral (cuboid) meshes always give stable results when using the volume-weighted interpolator.

## Corollary 2

Triangular meshes give stable results when using the volume-weighted interpolator if the node is placed at the circumcenter.

# Numerical robustness

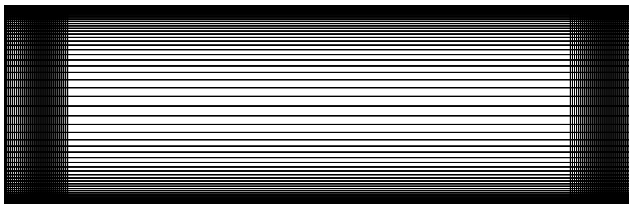


Figure 3: Highly distorted mesh in a  $Re_\tau = 395$  channel flow. Max aspect ratio is 250.

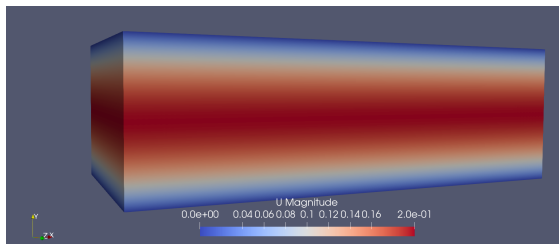


Figure 4: Velocity converged with the volume-weighted interpolator.

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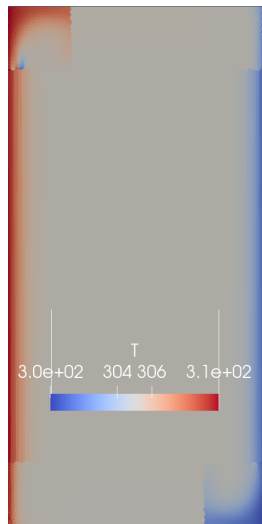
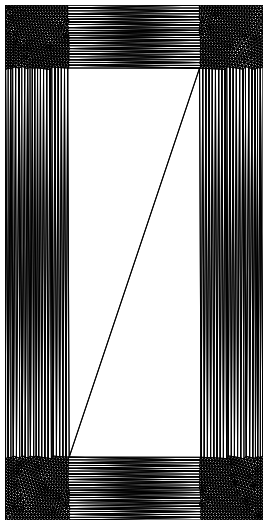


Figure 5: Test of the method's robustness in a  $Ra = 10^6$  differentially heated cavity.

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## Ongoing work

- Find the conditions for tetrahedral meshes in order to satisfy the geometrical conditions of the theorem.