

ON THE EVALUATION OF LARGE EDDY SIMULATION OF A WIND-TURBINE ARRAY BOUNDARY LAYER

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Introduction

Evaluate the performance of S3PR Large Eddy Simulation model on boundary layer and wind farm cases through different resolution meshes

Spatially filtered incompressible Navier-Stokes equations

$$\begin{aligned}\partial_t \bar{\mathbf{u}} + C(\bar{\mathbf{u}}, \bar{\mathbf{u}}) &= D(\bar{\mathbf{u}}) - \nabla p - \nabla \cdot \boldsymbol{\tau}(\bar{\mathbf{u}}); \\ \nabla \cdot \bar{\mathbf{u}} &= 0\end{aligned}$$

$\boldsymbol{\tau}(\bar{\mathbf{u}}) \approx -2\nu_e S(\bar{\mathbf{u}})$ is the LES closure

$S(\bar{\mathbf{u}}) = 1/2(\nabla \bar{\mathbf{u}} + \nabla \bar{\mathbf{u}}^T)$ is the rate-of-strain tensor

ν_e is the eddy viscosity for each model



Quick S3PQR theory review

Besides the trace, several mathematical invariants can be calculated from the gradient tensor $G = \nabla \bar{u}$, for example:

$$Q_G = (1/2)(tr^2(G) - tr(G^2))$$

$$R_G = det(G)$$

$$Q_S = (1/2)(tr^2(S) - tr(S^2))$$

$$R_S = det(S)$$

$$V_G^2 = 4(tr(S^2\Omega^2) - 2Q_S Q_\Omega)$$

$S = 1/2(G + G^T)$ and $\Omega = 1/2(G - G^T)$ are the symmetric and the skew-symmetric parts of the gradient tensor



S3PQR

The symmetric tensor GG^T formally based on the lowest-order approximation of the subgrid stress tensor is

$$\tau(\bar{\mathbf{u}}) = \frac{\Delta^2}{12} GG^T + \mathcal{O}(\Delta^4)$$

Three invariants of this tensor can be defined and are directly related to the previous ones

$$P_{GG^T} = \text{tr}(GG^T) = 2(Q_\Omega - Q_S)$$

$$Q_{GG^T} = 2(Q_\Omega - Q_S)^2 - Q_G^2 + 4\text{tr}(S^2\Omega^2)$$

$$R_{GG^T} = \det(GG^T) = \det(G)\det(G^T) = R_G^2$$



S3PQR

S3PQR: combination of two invariants of GG^T (Trias et al. (2015))

$$\nu_e^{S3PQ} = (C_{s3pq}\Delta)^2 P_{GG^T}^{-5/2} Q_{GG^T}^{3/2}$$

$$\nu_e^{S3PR} = (C_{s3pr}\Delta)^2 P_{GG^T}^{-1} R_{GG^T}^{1/2}$$

$$\nu_e^{S3QR} = (C_{s3qr}\Delta)^2 Q_{GG^T}^{-1} R_{GG^T}^{5/6}$$

where Δ is the subgrid characteristic length.

Two ways to determine the model constant C_{s3pq} :

1. Less or equal dissipation than Vreman's model.

$$C_{s3pq} = C_{s3pr} = C_{s3qr} = \sqrt{3}C_{Vr} \approx 0.458$$

2. The averaged dissipation of the models is equal to that of the Smagorinsky model.

$$C_{s3pq} = 0.572, C_{s3pr} = 0.709, C_{s3qr} = 0.762$$



Boundary layer and wind farm algorithm characteristics

- S3PQR LES model
- Semi-infinite domain that requires scaling procedure $y_\infty = L \frac{1+y}{1-y}$
- Pseudospectral: **Chebyshev polynomials** for Dirichlet and Newman boundary conditions

Main drawback (again): time-step of $O(1/N^2)$ for the convective term and $O(1/N^4)$ for the diffusive term!

High-resolution mesh computations are not feasible using fully explicit methods



Another algorithm details

- $Re_{\delta^*} = 1000$, where δ^* is the displacement thickness.
- Growing terms $GT(\bar{\mathbf{u}}, \bar{U})$, Spalart and Leonard (1987)
- Wind-turbine model, Calaf et al. (2010)
- We will test the **zero mean pressure gradient** case

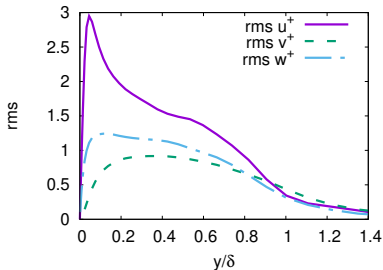
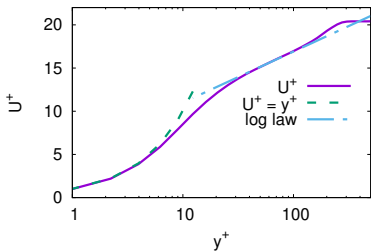
Previous computations (Folch et al. (2023)):

Boundary layer and wind farm comparison between LES models:
Smagorinsky, Verstappen, WALE, Vreman, and all the S3PQR.

Size domain for all of them $N_x \times N_y \times N_z = 32 \times 64 \times 32$ for
streamwise, normal, and spanwise directions



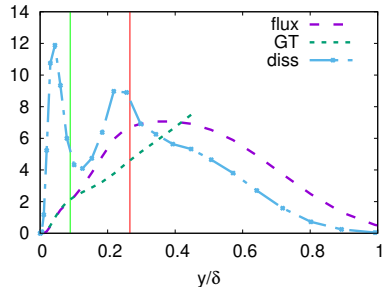
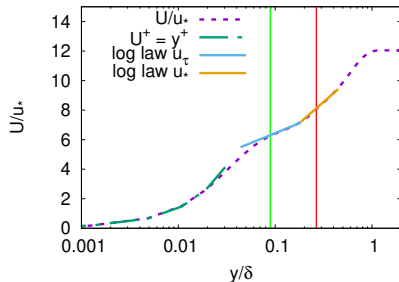
Boundary Layer



Case PR. **Left:** normalized average streamwise velocity profile, U^+ ; log law; $U^+ = y^+$. **Right:** $\text{rms } u^+$; $\text{rms } v^+$; $\text{rms } w^+$; δ is the boundary layer thickness



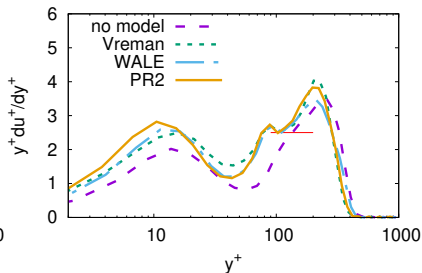
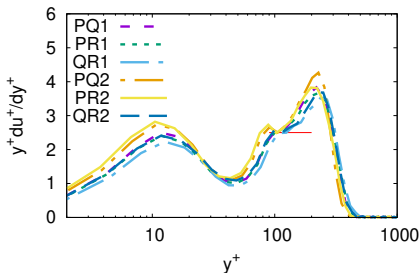
Wind farm



PR. Left: velocity. Right: Normalized mean kinetic energy contributions: **flux**, $\delta\Phi = -\langle uv \rangle U/u_*^3$; **GT**, normalized growing terms; **diss**, $-\langle uv \rangle \partial_y U/(u_*^3/\delta)$



Wind farm velocity derivative



Left: S3PQR models. Right: other LES models.



Evaluation

We will test: **the NO MODEL vs S3PR LES** algorithm
For completeness, the physical interpretation of PR:

$$\nu_e^{S3PR} = (C_{s3pr} \Delta)^2 P_{GG^T}^{-1} R_{GG^T}^{1/2}$$

$$\nu_e^{S3PR} \propto \frac{|\det(G)|}{2((1/4)\Omega_i \Omega_i + (1/2)S_{ij} S_{ij})} \sim \frac{|\partial_t GG^T|}{|w|^2 + |\epsilon|}$$

Mesh sizes $N_x \times N_y \times N_z$:

$32 \times 64 \times 32 \rightarrow 64 \times 64 \times 64 \rightarrow 96 \times 96 \times 96 \rightarrow 128 \times 128 \times 128$

For 128^3 : $\Delta x^+ \approx 20$, $\Delta z^+ \approx 6.7$ in wall units, and for the y-direction, 11 points within 9 wall units of the wall.



Semi-implicit algorithm

Recall:

- S3PQR yields **non uniform (and non constant)** eddy viscosity.
- The time step for the diffusive term in explicit schemes goes as $O(1/N^4)$

Solution: to compute **explicitly the convective** term and **implicitly the diffusive** term.

Caution: **At every step, we should compute a triple convolution sum** such as:

$\nu_e \times \text{Scaling} \times \text{Chebyshev derivative coefficients}$



Semi-implicit solution

A general class of two-step methods:

$$\text{Diffusion} = (\nu_p + \nu_e) \nabla^2 (\theta u^{n+1} + (1 - \theta) u^n)$$

where ν_p is the prescribed viscosity of the case

We will make a slight modification:

$$\text{Diffusion} = (0.5\nu_p) \nabla^2 u^{n+1} + \nabla \cdot ((0.5\nu_p + \nu_e) (\nabla u^n + \nabla (u^n)^T))$$

1. We calculate the matrix operator **only once** at the beginning with uniform and constant ν_p
2. Change in time-step: from $O(1/N^4)$ to $O(1/N^2)$



Evaluation results

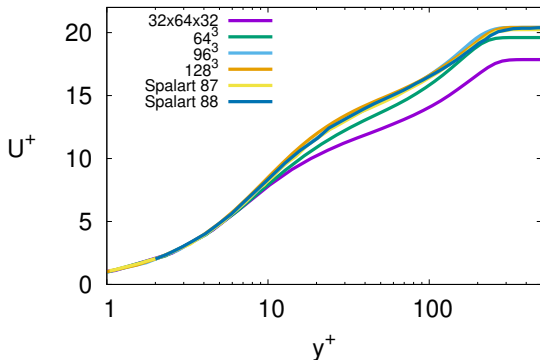
Boundary layer. Evolution of friction velocity:

Dimensions	No Model	S3PR
32x64x32	0.056	0.049
64 ³	0.051	0.048
96 ³	0.049	0.048
128 ³	0.049	0.048
Sp-Le DNS	0.049	

Reference: Spalart and Leonart (1987) 264x60x170 or Spalart (1988) 256x64x192



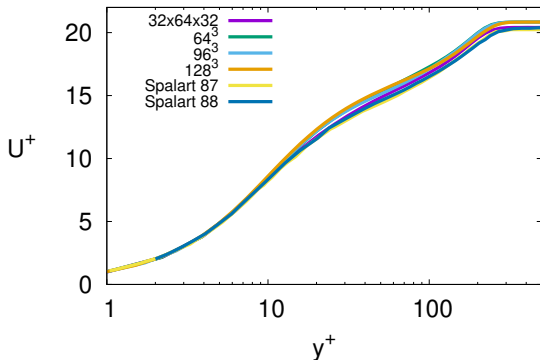
Boundary Layer. NO MODEL



Average streamwise velocity profile. Domain dimensions $N_x \times N_y \times N_z$



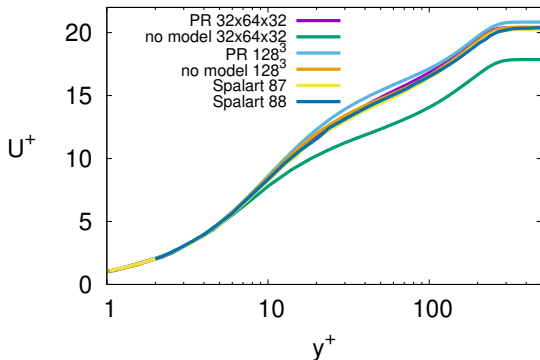
Boundary Layer. S3PR



Average streamwise velocity profile. Domain dimensions $N_x \times N_y \times N_z$



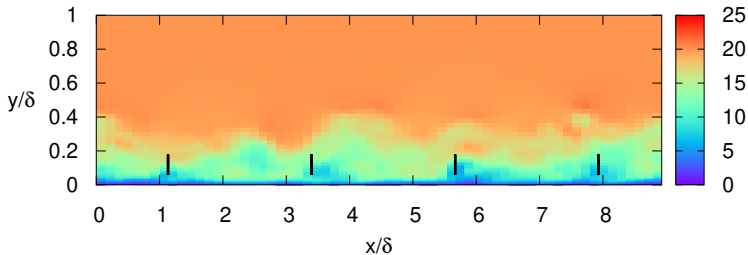
Boundary Layer. NO MODEL vs S3PR



Average streamwise velocity profile. Domain dimensions $N_x \times N_y \times N_z$



Wind farm. Instantaneous streamwise velocity

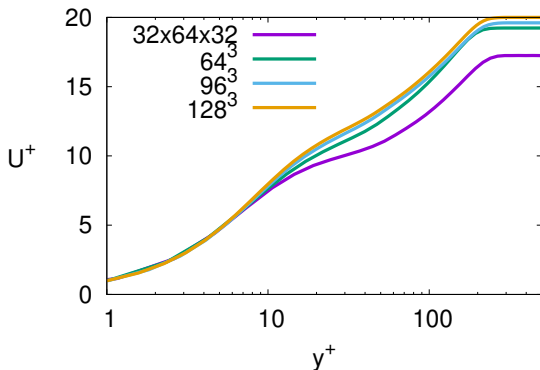


PR. Normalized streamwise velocity u^+ . $64 \times 64 \times 64$



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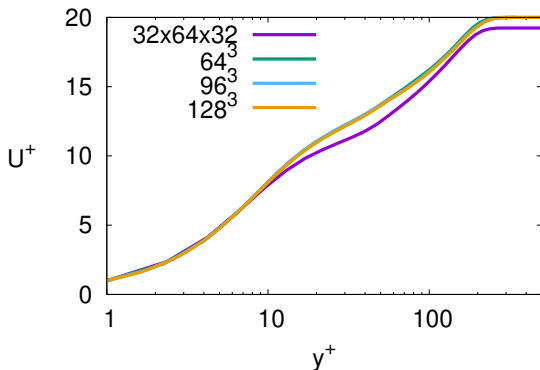
Wind farm. NO MODEL



Average streamwise velocity profile. Domain dimensions $N_x \times N_y \times N_z$



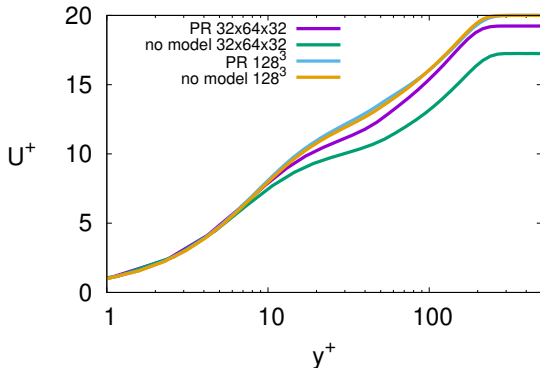
Wind farm. S3PR



Average streamwise velocity profile. Domain dimensions $N_x \times N_y \times N_z$



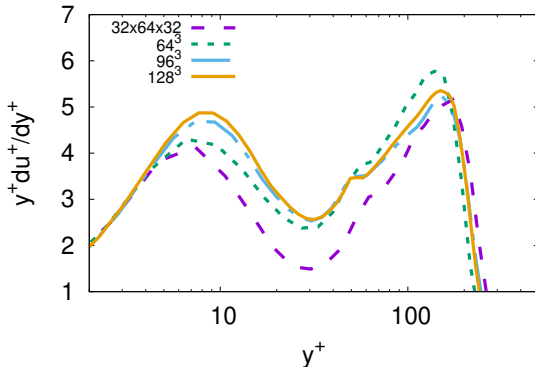
Wind farm. NO MODEL vs S3PR



Average streamwise velocity profile. Domain dimensions $N_x \times N_y \times N_z$



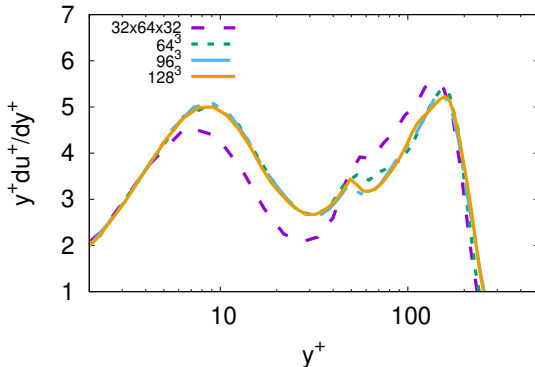
Wind farm. NO MODEL



Average streamwise velocity derivative. Domain dimensions
 $N_x \times N_y \times N_z$



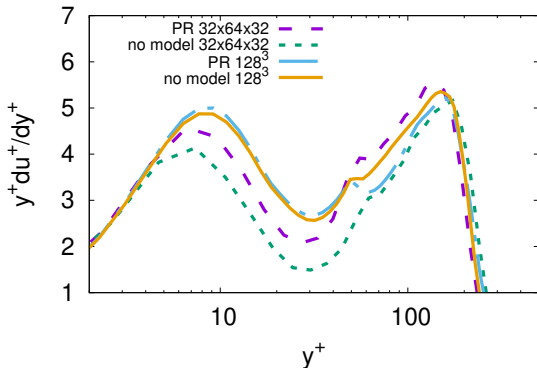
Wind farm. S3PR



Average streamwise velocity derivative. Domain dimensions
 $N_x \times N_y \times N_z$



Wind farm. NO MODEL vs S3PR



Average streamwise velocity derivative. Domain dimensions
 $N_x \times N_y \times N_z$



Conclusions

1. The no-model algorithm seems to approach an asymptotic profile for finer resolution.
2. The S3PR method gives the same asymptotic profile, even for coarse resolution.
3. The semi-implicit algorithm allows these higher-resolution computations.

Thank you for your attention.



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- **Spalart and Leonard (1987)**: P. R. Spalart, A. Leonard. Direct Numerical Simulation of Equilibrium Turbulent Boundary Layers. Turbulent Shear Flows 5. Springer, Berlin, Heidelberg.
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- **Calaf et al. (2010)**: M. Calaf, C. Meneveau, J. Meyers. Large eddy simulation study of fully developed wind-turbine array boundary layers. Physics of Fluids 22
- **Folch et al. (2023)**: D. Folch, F.X. Trias, A. Oliva. Assessment of LES models for a fully developed wind-turbine array boundary layer. International Symposium on Turbulence, Heat and Mass Transfer: Rome, Italy, 11-15 September 2023". Begell House, 2023

