

F.Xavier Trias¹, Jesús Ruano¹, Enrico Di Lavore¹ Alexev Duben². Andrev Gorobets²

¹Heat and Mass Transfer Technological Center, Technical University of Catalonia ²Keldysh Institute of Applied Mathematics of RAS, Russia







A rational length scale for large-eddy simulations on anisotropic grids

F.Xavier Trias¹. Jesús Ruano¹. Enrico Di Lavore¹ Alexey Duben², Andrey Gorobets²

¹Heat and Mass Transfer Technological Center, Technical University of Catalonia ²Keldysh Institute of Applied Mathematics of RAS, Russia

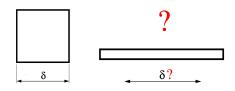




A rational length scale for large-eddy simulations on anisotropic grids

<u>F.Xavier Trias</u>¹, Jesús Ruano¹, Enrico Di Lavore¹ Alexey Duben², Andrey Gorobets²

 1 Heat and Mass Transfer Technological Center, Technical University of Catalonia 2 Keldysh Institute of Applied Mathematics of RAS, Russia



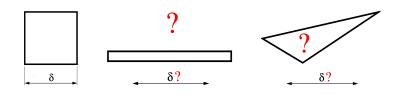




A rational length scale for large-eddy simulations on anisotropic grids

F.Xavier Trias¹. Jesús Ruano¹. Enrico Di Lavore¹ Alexey Duben², Andrey Gorobets²

¹Heat and Mass Transfer Technological Center, Technical University of Catalonia ²Keldysh Institute of Applied Mathematics of RAS, Russia



Contents

- Motivation
- 2 Revisiting FVM
- 3 A rational length scale for LES
- 4 Results
- Conclusions

General motivation: (very) large-scale DNS/LES



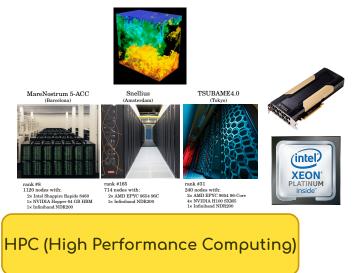


Motivation •00

DNS

•00

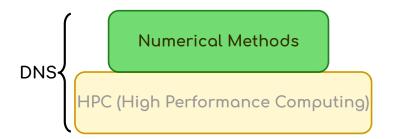
Results



General motivation: (very) large-scale DNS/LES

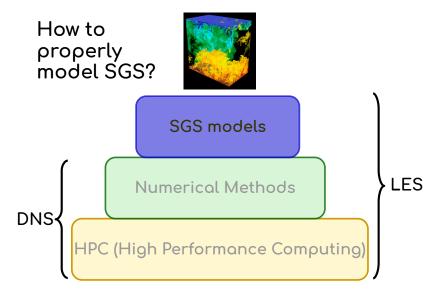


How to properly discretize NS?



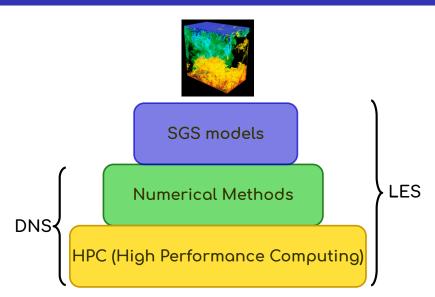
000

General motivation: (very) large-scale DNS/LES



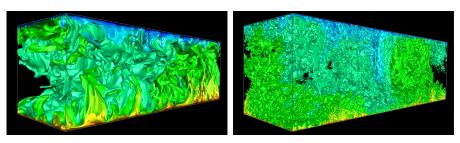
•00

General motivation: (very) large-scale DNS/LES



Research question #1:

• What are we indeed solving with finite volume method?



DNS¹ of air-filled Rayleigh-Bénard convection at $Ra = 10^8$ and 10^{10}

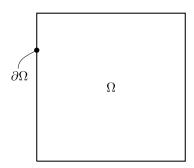
¹B.Sanderse, F.X.Trias. Energy-consistent discretization of viscous dissipation with application to natural convection flow. Computers & Fluids, 286:106473, 2025

Motivation

000

Research question #2:

• What are we interpolating? What is the correct interpretation?

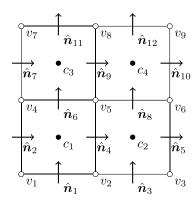


Motivation

000

Research question #2:

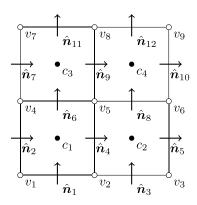
• What are we interpolating? What is the correct interpretation?

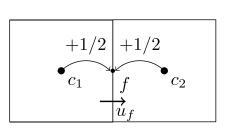


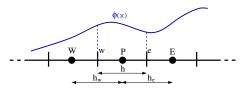
000

Research question #2:

• What are we interpolating? What is the correct interpretation?



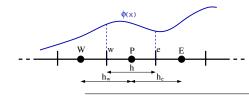




Box filter:

$$\overline{\phi}(x) = \frac{1}{h} \int_{x-h/2}^{x+h/2} \phi dx$$

$$\partial_{\mathsf{X}} \overline{\phi}|_{\mathsf{e}} = \overline{\partial_{\mathsf{X}} \phi}|_{\mathsf{e}} = (\phi_{\mathsf{E}} - \phi_{\mathsf{P}})/h_{\mathsf{e}}$$

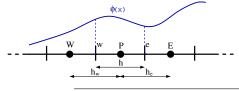


Box filter:

$$\overline{\phi}(x) = \frac{1}{h} \int_{x-h/2}^{x+h/2} \phi dx$$

$$\partial_{\mathsf{x}}\overline{\phi}|_{\mathsf{e}} = \overline{\partial_{\mathsf{x}}\phi}|_{\mathsf{e}} = (\phi_{\mathsf{E}} - \phi_{\mathsf{P}})/h_{\mathsf{e}}$$

$$\frac{\partial \phi}{\partial t} + \frac{\partial (\mathbf{u}\phi)}{\partial x} = \nu \frac{\partial^2 \phi}{\partial x^2}$$



Box filter:

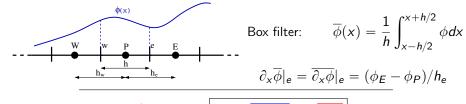
$$\overline{\phi}(x) = \frac{1}{h} \int_{x-h/2}^{x+h/2} \phi dx$$

$$\partial_x \overline{\phi}|_e = \overline{\partial_x \phi}|_e = (\phi_E - \phi_P)/h_e$$

$$\frac{\partial \phi}{\partial t} + \frac{\partial (\mathbf{u}\phi)}{\partial x} = \nu \frac{\partial^2 \phi}{\partial x^2}$$

$$\boxed{\frac{\partial \overline{\phi}}{\partial t} + \overline{\frac{\partial (\mathbf{u}\phi)}{\partial x}} = \nu \overline{\frac{\partial^2 \phi}{\partial x^2}}}$$

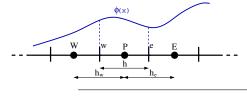
Exact FVM eq!



$$\partial_{\mathbf{x}}\overline{\phi}|_{\mathbf{e}} = \overline{\partial_{\mathbf{x}}\phi}|_{\mathbf{e}} = (\phi_{\mathbf{F}} - \phi_{\mathbf{P}})/h_{\mathbf{e}}$$

$$\frac{\partial \phi}{\partial t} + \frac{\partial (\mathbf{u}\phi)}{\partial x} = \nu \frac{\partial^2 \phi}{\partial x^2} \longrightarrow \boxed{\frac{\partial \overline{\phi}}{\partial t} + \overline{\frac{\partial (\mathbf{u}\phi)}{\partial x}} = \nu \overline{\frac{\partial^2 \phi}{\partial x^2}}} \text{ Exact FVM eq!}$$

$$h\frac{\partial\overline{\phi}_{P}}{\partial t} + (\mathbf{u}\phi)_{e} - (\mathbf{u}\phi)_{w} = \nu \left(\frac{\partial\phi}{\partial x} \Big|_{e} - \frac{\partial\phi}{\partial x} \Big|_{w} \right)$$



Box filter:

$$\overline{\phi}(x) = \frac{1}{h} \int_{x-h/2}^{x+h/2} \phi dx$$

$$\partial_x \overline{\phi}|_e = \overline{\partial_x \phi}|_e = (\phi_E - \phi_P)/h_e$$

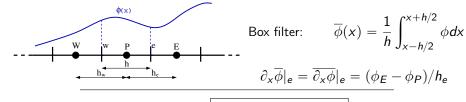
$$\frac{\partial \phi}{\partial t} + \frac{\partial (\mathbf{u}\phi)}{\partial x} = \nu \frac{\partial^2 \phi}{\partial x^2} \longrightarrow \left| \frac{\partial \overline{\phi}}{\partial t} + \frac{\overline{\partial (\mathbf{u}\phi)}}{\partial x} \right| = \nu \frac{\overline{\partial^2 \phi}}{\partial x^2} \quad \text{Exact FVM eq!}$$

$$h\frac{\partial\overline{\phi}_{P}}{\partial t} + (\mathbf{u}\phi)_{e} - (\mathbf{u}\phi)_{w} = \nu \left(\frac{\partial\phi}{\partial x} \Big|_{e} - \frac{\partial\phi}{\partial x} \Big|_{w} \right)$$

$$h\frac{\partial \overline{\phi}_{P}}{\partial x} + =$$

$$\left| \frac{\partial \overline{\phi}}{\partial t} + \right| =$$

Revisiting FVM

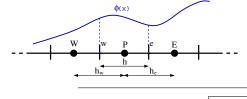


$$\frac{\partial \phi}{\partial t} + \frac{\partial (\mathbf{u}\phi)}{\partial x} = \nu \frac{\partial^2 \phi}{\partial x^2} \longrightarrow \left| \frac{\partial \overline{\phi}}{\partial t} + \frac{\overline{\partial (\mathbf{u}\phi)}}{\partial x} \right| = \nu \frac{\overline{\partial^2 \phi}}{\partial x^2}$$
 Exact FVM eq!

$$h\frac{\partial\overline{\phi}_{P}}{\partial t} + (\mathbf{u}\phi)_{e} - (\mathbf{u}\phi)_{w} = \nu \left(\frac{\partial\phi}{\partial x}\Big|_{e} - \frac{\partial\phi}{\partial x}\Big|_{w}\right)$$

$$= \nu \left(\frac{\overline{\phi}_{E} - \overline{\phi}_{P}}{h_{e}} - \frac{\overline{\phi}_{P} - \overline{\phi}_{W}}{h_{w}}\right)$$

$$\boxed{\frac{\partial \overline{\phi}}{\partial t} + = \nu \frac{\partial^2 \overline{\overline{\phi}}}{\partial x^2}}$$



Box filter:

$$\overline{\phi}(x) = \frac{1}{h} \int_{x-h/2}^{x+h/2} \phi dx$$

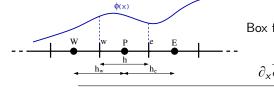
$$\partial_x \overline{\phi}|_e = \overline{\partial_x \phi}|_e = (\phi_E - \phi_P)/h_e$$

Results

$$\frac{\partial \phi}{\partial t} + \frac{\partial (\mathbf{u}\phi)}{\partial x} = \nu \frac{\partial^2 \phi}{\partial x^2} \longrightarrow \left| \frac{\partial \overline{\phi}}{\partial t} + \frac{\overline{\partial (\mathbf{u}\phi)}}{\partial x} \right| = \nu \frac{\overline{\partial^2 \phi}}{\partial x^2} \quad \text{Exact FVM eq!}$$

$$\begin{split} h\frac{\partial\overline{\phi}_{P}}{\partial t} + (\mathbf{u}\phi)_{e} - (\mathbf{u}\phi)_{w} &= \nu \left(\left. \frac{\partial\phi}{\partial x} \right|_{e} - \left. \frac{\partial\phi}{\partial x} \right|_{w} \right) \\ h\frac{\partial\overline{\phi}_{P}}{\partial t} + \mathbf{u}_{e} \frac{\overline{\phi}_{P} + \overline{\phi}_{E}}{2} - \mathbf{u}_{w} \frac{\overline{\phi}_{W} + \overline{\phi}_{P}}{2} &= \nu \left(\left. \frac{\overline{\phi}_{E} - \overline{\phi}_{P}}{h_{c}} - \frac{\overline{\phi}_{P} - \overline{\phi}_{W}}{h_{w}} \right) \end{split}$$

$$\boxed{\frac{\partial \overline{\phi}}{\partial t} + \frac{\overline{\partial (\mathbf{u}\overline{\phi})}}{\partial x} = \nu \overline{\frac{\partial^2 \overline{\overline{\phi}}}{\partial x^2}}}$$



$$\overline{\phi}(x) = \frac{1}{h} \int_{x-h/2}^{x+h/2} \phi dx$$

$$\frac{\partial_{\mathsf{x}}\overline{\phi}|_{\mathsf{e}} = \overline{\partial_{\mathsf{x}}\phi}|_{\mathsf{e}} = (\phi_{\mathsf{E}} - \phi_{\mathsf{P}})/h_{\mathsf{e}}}{}$$

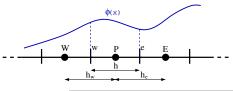
$$\frac{\partial q}{\partial t}$$

$$\frac{\partial \phi}{\partial t} + \frac{\partial (\mathbf{u}\phi)}{\partial x} = \nu \frac{\partial^2 \phi}{\partial x^2} \longrightarrow \left| \frac{\partial \overline{\phi}}{\partial t} + \frac{\overline{\partial (\mathbf{u}\phi)}}{\partial x} \right| = \nu \frac{\overline{\partial^2 \phi}}{\partial x^2} \quad \text{Exact FVM eq!}$$

$$\begin{split} h\frac{\partial\overline{\phi}_{P}}{\partial t} + (\mathbf{u}\phi)_{e} - (\mathbf{u}\phi)_{w} &= \nu \left(\left. \frac{\partial\phi}{\partial x} \right|_{e} - \left. \frac{\partial\phi}{\partial x} \right|_{w} \right) \\ h\frac{\partial\overline{\phi}_{P}}{\partial t} + \mathbf{u}_{e} \frac{\overline{\phi}_{P} + \overline{\phi}_{E}}{2} - \mathbf{u}_{w} \frac{\overline{\phi}_{W} + \overline{\phi}_{P}}{2} &= \nu \left(\frac{\overline{\phi}_{E} - \overline{\phi}_{P}}{h_{c}} - \frac{\overline{\phi}_{P} - \overline{\phi}_{W}}{h_{cw}} \right) \end{split}$$

$$\overline{\frac{\partial \overline{\phi}}{\partial t} + \frac{\overline{\partial (\mathbf{u} \widetilde{\phi})}}{\partial x}} = \nu \frac{\overline{\partial^2 \overline{\phi}}}{\partial x^2}$$

 $\frac{\partial \overline{\phi}}{\partial z} + \frac{\partial (\mathbf{u} \overline{\phi})}{\partial z} = \nu \frac{\partial^2 \overline{\phi}}{\partial z^2} \bigg| \qquad \text{where } \widetilde{\phi}_e = \frac{\overline{\phi}_P + \overline{\phi}_E}{2} = \overline{\phi}_e + \mathcal{O}(h^2)$



Box filter:

$$\overline{\phi}(x) = \frac{1}{h} \int_{x-h/2}^{x+h/2} \phi dx$$

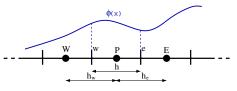
$$\partial_{\mathsf{x}}\overline{\phi}|_{\mathsf{e}} = \overline{\partial_{\mathsf{x}}\phi}|_{\mathsf{e}} = (\phi_{\mathsf{E}} - \phi_{\mathsf{P}})/h_{\mathsf{e}}$$

In summary (assuming that u is known):

$$\frac{\partial \overline{\phi}}{\partial t} + \frac{\overline{\partial (\mathbf{u}\phi)}}{\partial x} = \nu \frac{\overline{\partial^2 \phi}}{\partial x^2}$$

Instead, we are solving (in a 2^{nd} -order symmetry-preserving discretization):

$$\frac{\partial \overline{\phi}}{\partial t} + \frac{\overline{\partial (\mathbf{u}\overline{\phi})}}{\partial x} = \nu \frac{\overline{\partial^2 \overline{\phi}}}{\partial x^2}$$



Box filter:
$$\varphi \equiv \overline{\phi}(x) = \frac{1}{h} \int_{x-h/2}^{x+h/2} \phi dx$$

$$\partial_{\mathsf{x}} \overline{\phi}|_{\mathsf{e}} = \overline{\partial_{\mathsf{x}} \phi}|_{\mathsf{e}} = (\phi_{\mathsf{E}} - \phi_{\mathsf{P}})/h_{\mathsf{e}}$$

In summary (assuming that u is known):

$$\frac{\partial \overline{\phi}}{\partial t} + \frac{\overline{\partial (\mathbf{u}\phi)}}{\partial x} = \nu \frac{\overline{\partial^2 \phi}}{\partial x^2} \longrightarrow \boxed{\frac{\partial \varphi}{\partial t} + \frac{\partial (\mathbf{u}\varphi)}{\partial x} = \nu \frac{\partial^2 \varphi}{\partial x^2}} \boxed{\text{Exact FVM eq!}}$$

$$\boxed{\frac{\partial \varphi}{\partial t} + \frac{\partial (\mathbf{u}\varphi)}{\partial x} = \nu \frac{\partial^2 \varphi}{\partial x^2}}$$

Instead, we are solving (in a 2^{nd} -order symmetry-preserving discretization):

$$\frac{\partial \overline{\phi}}{\partial t} + \frac{\overline{\partial (\mathbf{u}\overline{\phi})}}{\partial x} = \nu \frac{\overline{\partial^2 \overline{\phi}}}{\partial x^2} \longrightarrow \boxed{\frac{\partial \varphi}{\partial t} + \overline{\frac{\partial (\mathbf{u}\overline{\varphi})}{\partial x}} = \nu \overline{\frac{\partial^2 \overline{\varphi}}{\partial x^2}}} \quad \text{Approx FVM eq!}$$

$$\boxed{\frac{\partial \varphi}{\partial t} + \frac{\overline{\partial (\mathbf{u}\overline{\varphi})}}{\partial x} = \nu \frac{\overline{\partial^2 \overline{\varphi}}}{\partial x^2}}$$

$$\partial_{t}\overline{\boldsymbol{u}} + (\overline{\boldsymbol{u}} \cdot \nabla)\overline{\boldsymbol{u}} = \nabla^{2}\overline{\boldsymbol{u}} - \nabla\overline{\boldsymbol{p}} - \nabla \cdot \boldsymbol{\tau}(\overline{\boldsymbol{u}}) ; \quad \nabla \cdot \overline{\boldsymbol{u}} = 0$$
eddy-viscosity $\longrightarrow \boldsymbol{\tau} (\overline{\boldsymbol{u}}) = -2\nu_{t}S(\overline{\boldsymbol{u}})$

²F.X.Trias, D.Folch, A.Gorobets, A.Oliva. **Physics of Fluids**, 27: 065103, 2015.

³M.H.Silvis, R.A.Remmerswaal, R.Verstappen, **Physics of Fluids**, 29: 015105, 2017.

Subgrid characteristic length for LES: state of the art

$$\begin{split} \partial_t \overline{\boldsymbol{u}} + (\overline{\boldsymbol{u}} \cdot \nabla) \overline{\boldsymbol{u}} &= \nabla^2 \overline{\boldsymbol{u}} - \nabla \overline{\boldsymbol{p}} - \nabla \cdot \boldsymbol{\tau}(\overline{\boldsymbol{u}}) \; ; \quad \nabla \cdot \overline{\boldsymbol{u}} = 0 \\ \text{eddy-viscosity} &\longrightarrow \boldsymbol{\tau} \; (\overline{\boldsymbol{u}}) = -2\nu_t S(\overline{\boldsymbol{u}}) \end{split}$$
$$\boxed{\nu_t = (C_m \delta)^2 D_m(\overline{\boldsymbol{u}})}$$

²F.X.Trias, D.Folch, A.Gorobets, A.Oliva. **Physics of Fluids**, 27: 065103, 2015.

³M.H.Silvis, R.A.Remmerswaal, R.Verstappen, **Physics of Fluids**, 29: 015105, 2017.

Subgrid characteristic length for LES: state of the art

$$\partial_t \overline{\boldsymbol{u}} + (\overline{\boldsymbol{u}} \cdot \nabla) \overline{\boldsymbol{u}} = \nabla^2 \overline{\boldsymbol{u}} - \nabla \overline{\boldsymbol{p}} - \nabla \cdot \boldsymbol{\tau}(\overline{\boldsymbol{u}}) \; ; \quad \nabla \cdot \overline{\boldsymbol{u}} = 0$$
 eddy-viscosity $\longrightarrow \boldsymbol{\tau} \; (\overline{\boldsymbol{u}}) = -2\nu_t S(\overline{\boldsymbol{u}})$

$$\nu_t = (C_m \delta)^2 D_m(\overline{\boldsymbol{u}})$$

 $D_m(\overline{u}) \longrightarrow \text{Smagorinsky (1963), WALE (1999), Vreman (2004),}$ QR-model (2011), σ -model (2011), S3PQR² (2015), vortex-stretching-based model³ (2017)

²F.X.Trias, D.Folch, A.Gorobets, A.Oliva. **Physics of Fluids**, 27: 065103, 2015.

³M.H.Silvis, R.A.Remmerswaal, R.Verstappen, **Physics of Fluids**, 29: 015105, 2017.

Subgrid characteristic length for LES: state of the art

$$\partial_t \overline{\boldsymbol{u}} + (\overline{\boldsymbol{u}} \cdot \nabla) \overline{\boldsymbol{u}} = \nabla^2 \overline{\boldsymbol{u}} - \nabla \overline{\boldsymbol{p}} - \nabla \cdot \boldsymbol{\tau}(\overline{\boldsymbol{u}}) \; ; \quad \nabla \cdot \overline{\boldsymbol{u}} = 0$$
eddy-viscosity $\longrightarrow \boldsymbol{\tau} \; (\overline{\boldsymbol{u}}) = -2\nu_t S(\overline{\boldsymbol{u}})$

$$\boxed{\nu_t = (C_m \delta)^2 D_m(\overline{\boldsymbol{u}})}$$

- $D_m(\overline{u}) \longrightarrow \text{Smagorinsky (1963), WALE (1999), Vreman (2004),}$ QR-model (2011), σ -model (2011), S3PQR² (2015), vortex-stretching-based model³ (2017)
- $C_m \longrightarrow Germano's dynamic model (1991), Lagrangian dynamic (1995),$ Global dynamic approach (2006)

²F.X.Trias, D.Folch, A.Gorobets, A.Oliva. **Physics of Fluids**, 27: 065103, 2015.

³M.H.Silvis, R.A.Remmerswaal, R.Verstappen, **Physics of Fluids**, 29: 015105, 2017.

$$\partial_{t}\overline{\boldsymbol{u}} + (\overline{\boldsymbol{u}} \cdot \nabla)\overline{\boldsymbol{u}} = \nabla^{2}\overline{\boldsymbol{u}} - \nabla\overline{p} - \nabla \cdot \boldsymbol{\tau}(\overline{\boldsymbol{u}}) \; ; \qquad \nabla \cdot \overline{\boldsymbol{u}} = 0$$
 eddy-viscosity $\longrightarrow \boldsymbol{\tau} \; (\overline{\boldsymbol{u}}) = -2\nu_{\boldsymbol{t}}S(\overline{\boldsymbol{u}})$

$$\boxed{\nu_t = (C_m \delta)^2 D_m(\overline{\boldsymbol{u}})}$$

$$D_m(\overline{u}) \longrightarrow \text{Smagorinsky (1963), WALE (1999), Vreman (2004),}$$

QR-model (2011), σ -model (2011), S3PQR² (2015),

vortex-stretching-based model³ (2017)

 $C_m \longrightarrow \text{Germano's dynamic model (1991), Lagrangian dynamic (1995),}$ Global dynamic approach (2006)



²F.X.Trias, D.Folch, A.Gorobets, A.Oliva. **Physics of Fluids**, 27: 065103, 2015.

³M.H.Silvis, R.A.Remmerswaal, R.Verstappen, **Physics of Fluids**, 29: 015105, 2017.

• In the context of LES, most popular (by far) is:

$$\begin{array}{|c|c|}\hline \delta_{\rm vol} = (\Delta x \Delta y \Delta z)^{1/3} & \longleftarrow \text{Deardorff (1970)}\\ \delta_{\rm Sco} = f(\textit{a}_1,\textit{a}_2)\delta_{\rm vol}, & \delta_{\textit{L}^2} = \sqrt{(\Delta x^2 + \Delta y^2 + \Delta z^2)/3} \end{array}$$

A rational length scale for LES

In the context of LES, most popular (by far) is:

$$\begin{split} \delta_{\mathrm{vol}} &= (\Delta x \Delta y \Delta z)^{1/3} \\ &\iff \text{Deardorff (1970)} \\ \delta_{\mathrm{Sco}} &= f(a_1, a_2) \delta_{\mathrm{vol}}, \qquad \delta_{L^2} = \sqrt{(\Delta x^2 + \Delta y^2 + \Delta z^2)/3} \\ \delta_{\mathrm{lsq}} &= \sqrt{\frac{\hat{\mathsf{G}} \hat{\mathsf{G}}^T : \mathsf{G} \mathsf{G}^T}{\mathsf{G} \mathsf{G}^T : \mathsf{G} \mathsf{G}^T}} \quad \Longleftrightarrow \text{Trias et al. (2017)} \end{split}$$

A rational length scale for LES

In the context of LES, most popular (by far) is:

• In the context of DES:

$$\delta_{\max} = \max(\Delta x, \Delta y, \Delta z)$$
 Sparlart et al. (1997)

Flow-dependant definitions

$$\delta_{\omega} = \sqrt{(\omega_{x}^{2} \Delta y \Delta z + \omega_{y}^{2} \Delta x \Delta z + \omega_{z}^{2} \Delta x \Delta y)/|\omega|^{2}} \iff \text{Chauvet et al. (2007)}$$

$$\tilde{\delta}_{\omega} = \frac{1}{\sqrt{3}} \max_{n,m=1,\dots,8} |\mathbf{I}_{n} - \mathbf{I}_{m}| \iff \text{Mockett et al. (2015)}$$

 $\delta_{SLA} = \tilde{\delta}_{\alpha}, F_{KH}(VTM)$

 \leftarrow Shur et al. (2015) $_{9/21}$

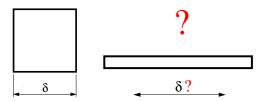
Research question #3:

• Can we establish a simple, robust, and easily implementable definition of δ for any type of grid that minimizes the impact of mesh anisotropies on the performance of subgrid-scale models?



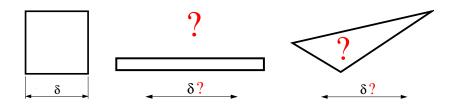
Research question #3:

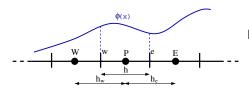
• Can we establish a simple, robust, and easily implementable definition of δ for any type of grid that minimizes the impact of mesh anisotropies on the performance of subgrid-scale models?



Research question #3:

• Can we establish a simple, robust, and easily implementable definition of δ for any type of grid that minimizes the impact of mesh anisotropies on the performance of subgrid-scale models?





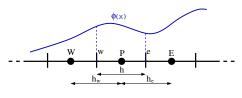
Box filter:
$$\overline{\phi}(x) = \frac{1}{h} \int_{x-h/2}^{x+h/2} \phi dx$$

Results

$$\partial_{\mathsf{x}}\overline{\phi} = \overline{\partial_{\mathsf{x}}\phi} = (\phi_{\mathsf{e}} - \phi_{\mathsf{w}})/h$$

Remark #1: the actual filter length, δ , when computing the face derivative is h_e , i.e., the distance between the adjacent nodes P and E.

A rational length scale for LES

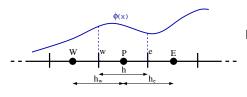


Box filter: $\overline{\phi}(x) = \frac{1}{h} \int_{x-h/2}^{x+h/2} \phi dx$

$$\partial_{\mathsf{x}}\overline{\phi} = \overline{\partial_{\mathsf{x}}\phi} = (\phi_{\mathsf{e}} - \phi_{\mathsf{w}})/h$$

The diffusive term in a FVM framework is approximated as follows

$$\left. \frac{\partial}{\partial x} \left(\Gamma \frac{\partial \phi}{\partial x} \right) \right|_{P} \approx \left. \frac{1}{h} \left(\Gamma \frac{\partial \phi}{\partial x} \right|_{e} - \left. \Gamma \frac{\partial \phi}{\partial x} \right|_{w} \right) = \left. \frac{\overline{\partial}}{\partial x} \left(\Gamma \frac{\partial \phi}{\partial x} \right) \right|_{P}$$



Box filter: $\overline{\phi}(x) = \frac{1}{h} \int_{-\pi/2}^{x+h/2} \phi dx$

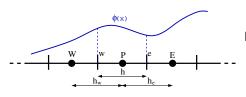
Results

$$\partial_{\mathsf{x}}\overline{\phi} = \overline{\partial_{\mathsf{x}}\phi} = (\phi_{\mathsf{e}} - \phi_{\mathsf{w}})/h$$

The diffusive term in a FVM framework is approximated as follows

$$\frac{\partial}{\partial x} \left(\Gamma \frac{\partial \phi}{\partial x} \right) \Big|_{P} \approx \frac{1}{h} \left(\Gamma \frac{\partial \phi}{\partial x} \Big|_{e} - \Gamma \frac{\partial \phi}{\partial x} \Big|_{w} \right) = \frac{\overline{\partial}}{\partial x} \left(\Gamma \frac{\partial \phi}{\partial x} \right) \Big|_{P}$$

$$\approx \frac{1}{h} \left(\Gamma_{e} \frac{\phi_{E} - \phi_{P}}{h_{e}} - \Gamma_{w} \frac{\phi_{P} - \phi_{W}}{h_{w}} \right) = \overline{\frac{\partial}{\partial x} \left(\Gamma \frac{\overline{\partial \phi}}{\partial x} \right)} \Big|_{P}$$



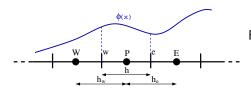
 $\overline{\phi}(x) = \frac{1}{h} \int_{-\pi/2}^{x+n/2} \phi dx$

Results

$$\partial_x \overline{\phi} = \overline{\partial_x \phi} = (\phi_e - \phi_w)/h$$

The diffusive term in a FVM framework is approximated as follows

$$\begin{split} \frac{\partial}{\partial x} \left(\Gamma \frac{\partial \phi}{\partial x} \right) \bigg|_{P} &\approx \frac{1}{h} \left(\Gamma \frac{\partial \phi}{\partial x} \bigg|_{e} - \Gamma \frac{\partial \phi}{\partial x} \bigg|_{w} \right) = \overline{\frac{\partial}{\partial x} \left(\Gamma \frac{\partial \phi}{\partial x} \right)} \bigg|_{P} \\ &\approx \frac{1}{h} \left(\Gamma_{e} \frac{\phi_{E} - \phi_{P}}{h_{e}} - \Gamma_{w} \frac{\phi_{P} - \phi_{W}}{h_{w}} \right) = \overline{\frac{\partial}{\partial x} \left(\Gamma \frac{\overline{\partial \phi}}{\partial x} \right)} \bigg|_{P} \\ &\approx \frac{1}{h} \left(\overline{\frac{\Gamma_{E} + \Gamma_{P}}{2} \frac{\phi_{E} - \phi_{P}}{h_{e}} - \overline{\frac{\Gamma_{P} + \Gamma_{W}}{2} \frac{\phi_{P} - \phi_{W}}{h_{w}}} \right) = \overline{\frac{\partial}{\partial x} \left(\overline{\Gamma} \frac{\overline{\partial \phi}}{\partial x} \right)} \bigg|_{P} \end{split}$$



Box filter:

A rational length scale for LES

$$\overline{\phi}(x) = \frac{1}{h} \int_{x-h/2}^{x+h/2} \phi dx$$

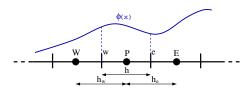
$$\partial_{\mathsf{x}}\overline{\phi} = \overline{\partial_{\mathsf{x}}\phi} = (\phi_{\mathsf{e}} - \phi_{\mathsf{w}})/h$$

Remark #1: the actual filter length, δ , when computing the face derivative is h_e , *i.e.*, the distance between the adjacent nodes P and E.

Remark #2: two filtering operations are performed when computing the diffusive term:

- the calculation of the face derivative
- ullet the cell-to-face interpolation of Γ

$$\left. \frac{\partial}{\partial x} \left(\overline{\Gamma} \frac{\overline{\partial \phi}}{\partial x} \right) \right|_{P}$$



Box filter:

$$\overline{\phi}(x) = \frac{1}{h} \int_{x-h/2}^{x+h/2} \phi dx$$

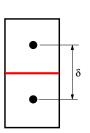
$$\partial_x \overline{\phi} = \overline{\partial_x \phi} = (\phi_e - \phi_w)/h$$

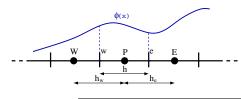
Remark #1: the actual filter length, δ , when computing the face derivative is h_e , i.e., the distance between the adjacent nodes P and E.

Remark #2: two filtering operations are performed when computing the diffusive term:

- the calculation of the face derivative
- the cell-to-face interpolation of Γ

Both filtering operators share the same filter length, δ : namely, the distance between the nodes adjacent to the corresponding face.





A rational length scale for LES

$$\overline{\phi}(x) = \frac{1}{h} \int_{x-h/2}^{x+h/2} \phi dx$$

$$\partial_{\mathsf{x}}\overline{\phi} = \overline{\partial_{\mathsf{x}}\phi} = (\phi_{\mathsf{e}} - \phi_{\mathsf{w}})/h$$

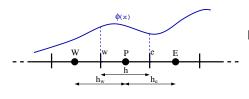
Remark #1: the actual filter length, δ , when computing the face derivative is h_e , i.e., the distance between the adjacent nodes P and E.

Remark #2: two filtering operations are performed when computing the diffusive term:

- the calculation of the face derivative
- the cell-to-face interpolation of Γ



Both filtering operators share the same filter length, δ ; namely, the distance between the nodes adjacent to the corresponding face.



Box filter:

A rational length scale for LES

$$\overline{\phi}(x) = \frac{1}{h} \int_{x-h/2}^{x+h/2} \phi dx$$

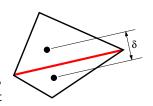
$$\partial_{\mathsf{x}}\overline{\phi} = \overline{\partial_{\mathsf{x}}\phi} = (\phi_{\mathsf{e}} - \phi_{\mathsf{w}})/h$$

Remark #1: the actual filter length, δ , when computing the face derivative is h_e , i.e., the distance between the adjacent nodes P and E.

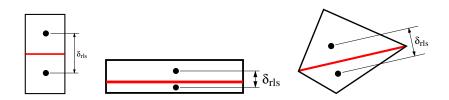
Remark #2: two filtering operations are performed when computing the diffusive term:

- the calculation of the face derivative
- the cell-to-face interpolation of Γ

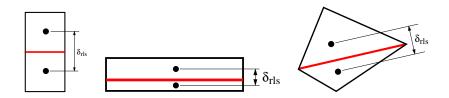
Both filtering operators share the same filter length, δ ; namely, the distance between the nodes adjacent to the corresponding face.



Properties of new definition, $\delta_{\rm rls}$

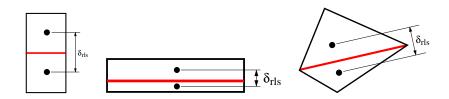


Properties of new definition, $\delta_{\rm rls}$



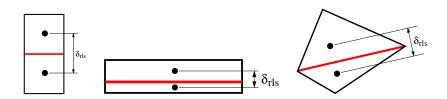
Locally defined

Properties of new definition, $\delta_{\rm rls}$



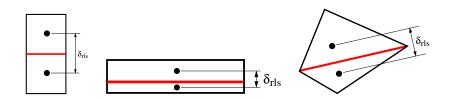
- Locally defined
- Well-bounded: $\Delta x \leqslant \delta_{\rm rls} \leqslant \Delta z$ (assuming $\Delta x \leqslant \Delta y \leqslant \Delta z$)

Properties of new definition, $\delta_{\rm rls}$



- Locally defined
- Well-bounded: $\Delta x \leq \delta_{\rm rls} \leq \Delta z$ (assuming $\Delta x \leq \Delta y \leq \Delta z$)
- Sensitive to flow orientation, e.g. shear layers

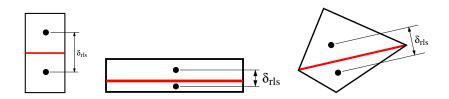
Properties of new definition, $\delta_{\rm rls}$



A rational length scale for LES

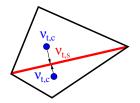
- Locally defined
- Well-bounded: $\Delta x \leq \delta_{\rm rls} \leq \Delta z$ (assuming $\Delta x \leq \Delta y \leq \Delta z$)
- Sensitive to flow orientation, e.g. shear layers
- Applicable to unstructured grids

Properties of new definition, $\delta_{\rm rls}$

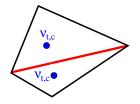


- Locally defined
- Well-bounded: $\Delta x \leq \delta_{\rm rls} \leq \Delta z$ (assuming $\Delta x \leq \Delta y \leq \Delta z$)
- Sensitive to flow orientation, e.g. shear layers
- Applicable to unstructured grids
- Easy and cheap

$$m{
u}_{t,c} \stackrel{\mathsf{interpolation}}{-\!\!\!\!-\!\!\!\!-\!\!\!\!-\!\!\!\!-\!\!\!\!-} m{
u}_{t,s}$$



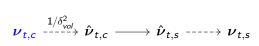
$$\boldsymbol{\nu}_{t,c} \dashrightarrow \hat{\boldsymbol{\nu}}_{t,c} \longrightarrow \hat{\boldsymbol{\nu}}_{t,s} \dashrightarrow \boldsymbol{\nu}_{t,s}$$

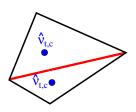


$$m{
u}_{t,c} \xrightarrow{----} \hat{m{
u}}_{t,c} \xrightarrow{----} \hat{m{
u}}_{t,s} \xrightarrow{----} m{
u}_{t,s}$$

0000000

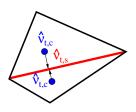
A rational length scale





Implementation and an alternative definition

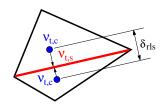
$$m{
u}_{t,c} \stackrel{1/\delta_{vol}^2}{---} \hat{m{
u}}_{t,c} \stackrel{ ext{interpolation}}{\longrightarrow} \hat{m{
u}}_{t,s} \stackrel{\cdots}{\longrightarrow} m{
u}_{t,s}$$



A rational length scale for LES

0000000

$$\boldsymbol{\nu}_{t,c} \xrightarrow{1/\delta_{vol}^2} \hat{\boldsymbol{\nu}}_{t,c} \xrightarrow{\text{interpolation}} \hat{\boldsymbol{\nu}}_{t,s} \xrightarrow{\delta_{rls}^2} \boldsymbol{\nu}_{t,s}$$



Implementation and an alternative definition

$$\begin{array}{c} \boldsymbol{\nu}_{t,c} \xrightarrow{\text{interpolation}} \boldsymbol{\nu}_{t,s} \\ \\ \boldsymbol{\nu}_{t,c} \xrightarrow{1/\delta_{vol}^2} \hat{\boldsymbol{\nu}}_{t,c} \xrightarrow{\text{interpolation}} \hat{\boldsymbol{\nu}}_{t,s} \xrightarrow{\delta_{rls}^2} \boldsymbol{\nu}_{t,s} \end{array}$$

We can also compute an equivalent filter length, $\tilde{\delta}_{rls},$ that leads to the same local dissipation

$$\tilde{\delta}_{rls}^2 \hat{\nu}_t \mathbf{G} : \mathbf{G} = \hat{\nu}_t \hat{\mathbf{G}} : \hat{\mathbf{G}}$$

where $\hat{G} \equiv G\Delta$ and $\Delta \equiv \operatorname{diag}(\Delta x, \Delta y, \Delta x)$.

Results

A rational length scale

Implementation and an alternative definition

$$\begin{array}{c} \boldsymbol{\nu}_{t,c} \xrightarrow{\text{interpolation}} \boldsymbol{\nu}_{t,s} \\ \\ \boldsymbol{\nu}_{t,c} \xrightarrow{1/\delta_{vol}^2} \hat{\boldsymbol{\nu}}_{t,c} \xrightarrow{\text{interpolation}} \hat{\boldsymbol{\nu}}_{t,s} \xrightarrow{\delta_{rls}^2} \boldsymbol{\nu}_{t,s} \end{array}$$

We can also compute an equivalent filter length, $\tilde{\delta}_{rls}$, that leads to the same local dissipation

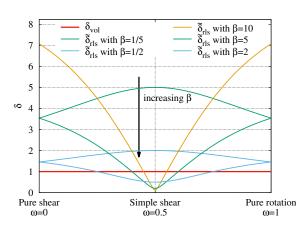
$$\tilde{\delta}_{\mathrm{rls}}^2 \hat{\nu}_t \mathsf{G} : \mathsf{G} = \hat{\nu}_t \hat{\mathsf{G}} : \hat{\mathsf{G}} \qquad \Longrightarrow \qquad \left| \tilde{\delta}_{\mathrm{rls}} = \sqrt{\frac{\hat{\mathsf{G}} : \hat{\mathsf{G}}}{\mathsf{G} : \mathsf{G}}} = \sqrt{\frac{\mathrm{tr}(\hat{\mathsf{G}}\hat{\mathsf{G}}^T)}{\mathrm{tr}(\mathsf{G}\mathsf{G}^T)}} \right|$$

where $\hat{G} \equiv G\Delta$ and $\Delta \equiv \operatorname{diag}(\Delta x, \Delta y, \Delta x)$.

Properties of new definition $\tilde{\delta}_{\rm rls}$

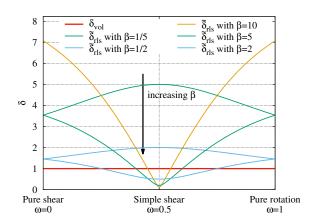
$$\Delta = \left(\begin{array}{cc} \Delta x & 0 \\ 0 & \Delta y \end{array}\right)$$

$$\mathsf{G} = \left(\begin{array}{cc} \partial_{\mathsf{x}} u & \partial_{\mathsf{y}} u \\ \partial_{\mathsf{y}} u & \partial_{\mathsf{y}} v \end{array} \right)$$



Properties of new definition $\tilde{\delta}_{\rm rls}$

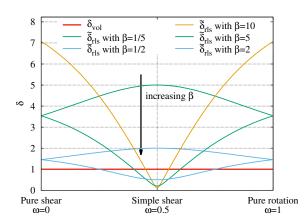
$$\Delta = \left(\begin{array}{cc} \Delta x & 0 \\ 0 & \Delta y \end{array} \right) = \left(\begin{array}{cc} \beta & 0 \\ 0 & \beta^{-1} \end{array} \right) \quad \mathsf{G} = \left(\begin{array}{cc} \partial_x u & \partial_y u \\ \partial_y u & \partial_y v \end{array} \right)$$



Results

Properties of new definition $\tilde{\delta}_{rls}$

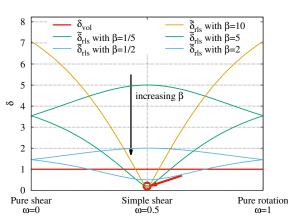
$$\Delta = \left(\begin{array}{cc} \Delta x & 0 \\ 0 & \Delta y \end{array} \right) \\ = \left(\begin{array}{cc} \beta & 0 \\ 0 & \beta^{-1} \end{array} \right) \\ = \left(\begin{array}{cc} \partial_x u & \partial_y u \\ \partial_y u & \partial_y v \end{array} \right) \\ = \left(\begin{array}{cc} 0 & 1 \\ 1 - 2\omega & 0 \end{array} \right)$$



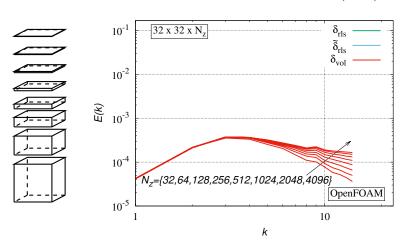
Properties of new definition $\tilde{\delta}_{rls}$

Properties of new definition
$$\sigma_{
m rls}$$

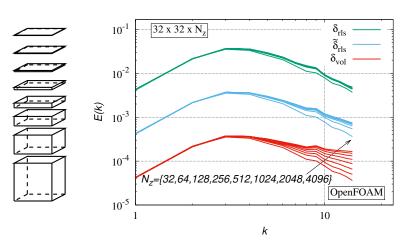
$$\Delta = \left(\begin{array}{cc} \Delta x & 0 \\ 0 & \Delta y \end{array} \right) = \left(\begin{array}{cc} \beta & 0 \\ 0 & \beta^{-1} \end{array} \right) \quad G = \left(\begin{array}{cc} \partial_x u & \partial_y u \\ \partial_y u & \partial_y v \end{array} \right) = \left(\begin{array}{cc} 0 & 1 \\ 0 & 0 \end{array} \right) \quad \boxed{\tilde{\delta}_{\rm rls} = \Delta y}$$



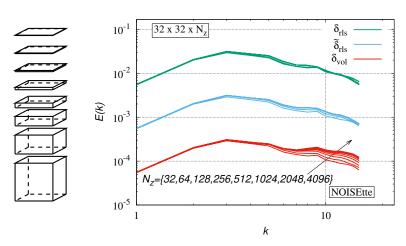
Isotropic turbulence on anisotropic grids



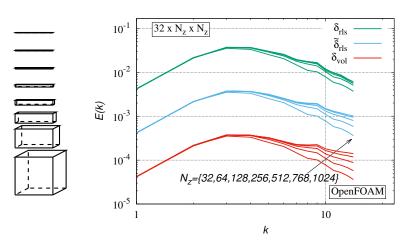
Isotropic turbulence on anisotropic grids

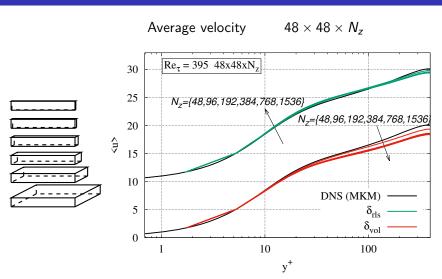


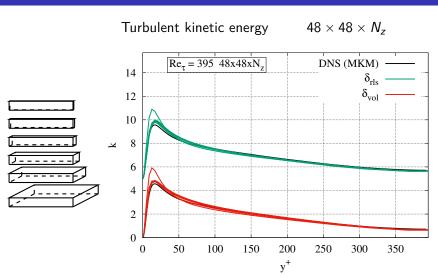
Isotropic turbulence on anisotropic grids

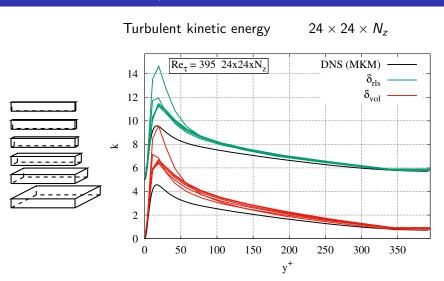


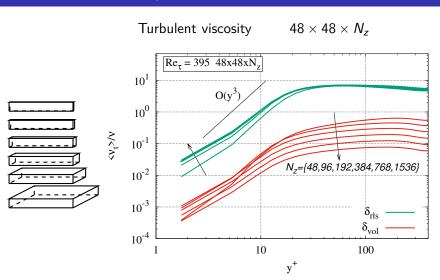
Isotropic turbulence on anisotropic grids



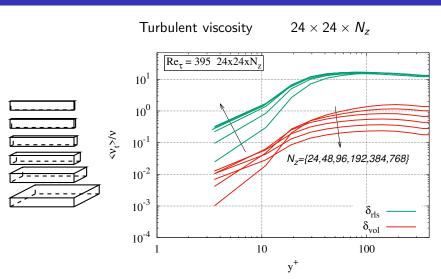






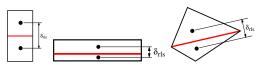


Turbulent channel flow at $Re_{\tau} = 395$



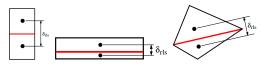
A rational length scale for LES

ullet A new definition for δ has been proposed



$$\tilde{\delta}_{\rm rls} = \sqrt{\frac{\hat{\mathsf{G}}:\hat{\mathsf{G}}}{\mathsf{G}:\mathsf{G}}}$$

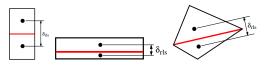
• A new definition for δ has been proposed



$$\tilde{\delta}_{
m rls} = \sqrt{rac{\hat{\mathsf{G}}:\hat{\mathsf{G}}}{\mathsf{G}:\mathsf{G}}}$$

- It is locally defined, well-bounded, cheap and easy to implement
- Suitable for unstructured grids

• A new definition for δ has been proposed



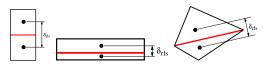
$$\tilde{\delta}_{\mathrm{rls}} = \sqrt{\frac{\hat{\mathsf{G}} : \hat{\mathsf{G}}}{\mathsf{G} : \mathsf{G}}}$$

It is locally defined, well-bounded, cheap and easy to implement

A rational length scale for LES

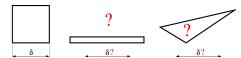
- Suitable for unstructured grids
- LES tests:
 - HIT
 - Turbulent channel flow √
 - Unstructured grids (on-going)

• A new definition for δ has been proposed



$$\tilde{\delta}_{
m rls} = \sqrt{rac{\hat{\sf G}:\hat{\sf G}}{{\sf G}:{\sf G}}}$$

- It is locally defined, well-bounded, cheap and easy to implement
- Suitable for unstructured grids
- LES tests:
 - HIT
 - Turbulent channel flow √
 - Unstructured grids (on-going)



Takeaway message:

• Definition of δ can have a big effect on simulation results





ABSTRACT

Due to the prohibitive cost of resolving all relevant scales, direct numerical simulations of turbulence remain unfeasible for most real-world applications. Consequently, dynamically simplified formulations are needed for coarse-grained simulations. In this regard, eddy-viscosity models for large-eddy simulation (LES) are widely used in both academia and industry. These models require a subgrid characteristic length, typically linked to the local grid size. While this length scale corresponds to the mesh step for isotropic grids, its definition for unstructured

https://github.com/jruanoperez/DHIT