

Can we hit the ultimate regime of thermal turbulence using large-scale LES simulations?

F. Xavier Trias¹

¹Heat and Mass Transfer Technological Center, Technical University of Catalonia

²Termo Fluids S.L. Carrer de Magí Colet 8, 08204 Sabadell (Barcelona), Spain

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`francesc.xavier.trias@upc.edu`

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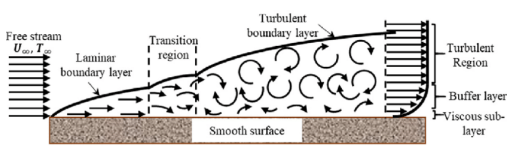
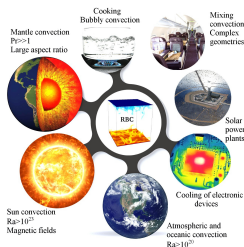
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- 2 Preserving symmetries at discrete level
- 3 Portability and beyond
- 4 LES of RBC
- 5 Conclusions

Motivation

Research question #1:

- Can we hit the ultimate regime of thermal turbulence ?

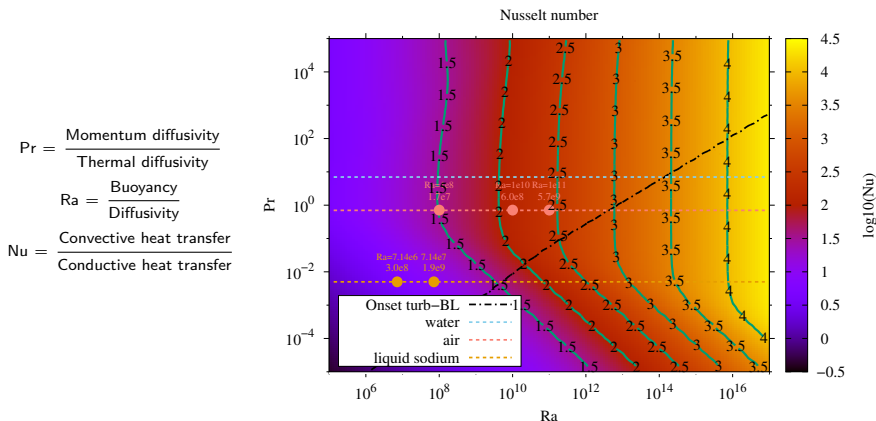


¹Robert H. Kraichnan. *Turbulent Thermal Convection at Arbitrary Prandtl Number*, *Physics of Fluids*, 1374–1389, 1962.

Motivation

Research question #1:

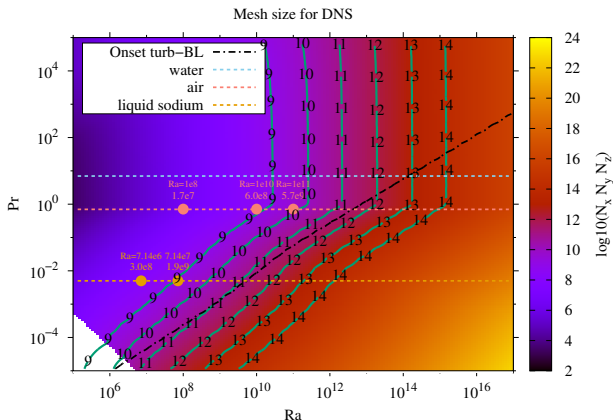
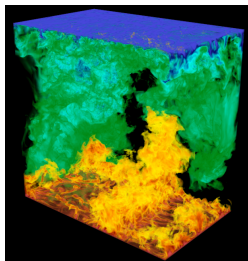
- Can we hit the ultimate regime of thermal turbulence ?



Motivation

Research question #1:

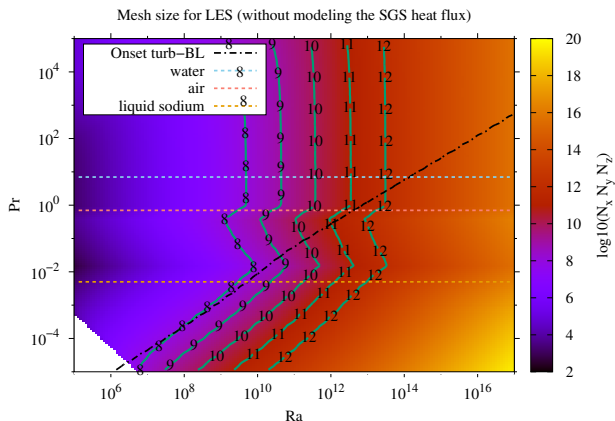
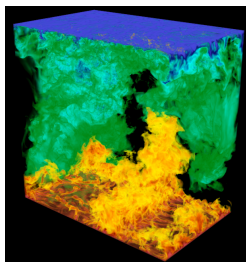
- Can we hit the ultimate regime of thermal turbulence with **DNS**?



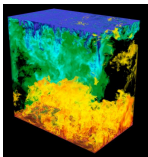
Motivation

Research question #1:

- Can we hit the ultimate regime of thermal turbulence with **LES**?

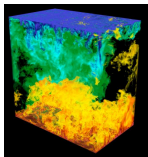


Motivation



DNS {

Motivation



HAWK



Rank #27
5,632 nodes with:
2 AMD EPYC 7742
(64 cores each)

MareNostrum 4



Rank #82
3456 nodes with:
2x Intel Xeon 8160
1x Intel Omni-Path

Marconi100



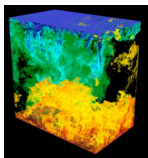
Rank #21
980 nodes with:
2 IBM Power9
4 NVIDIA Volta V100



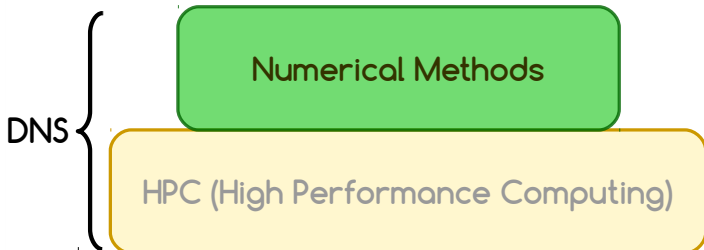
DNS {

HPC (High Performance Computing)

Motivation

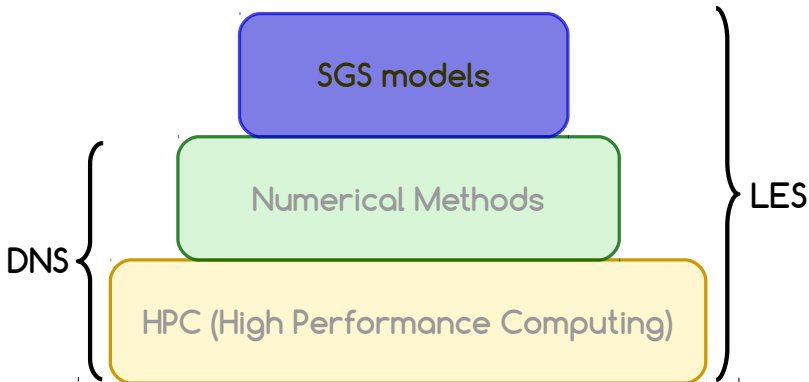
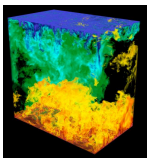


How to properly discretize NS?

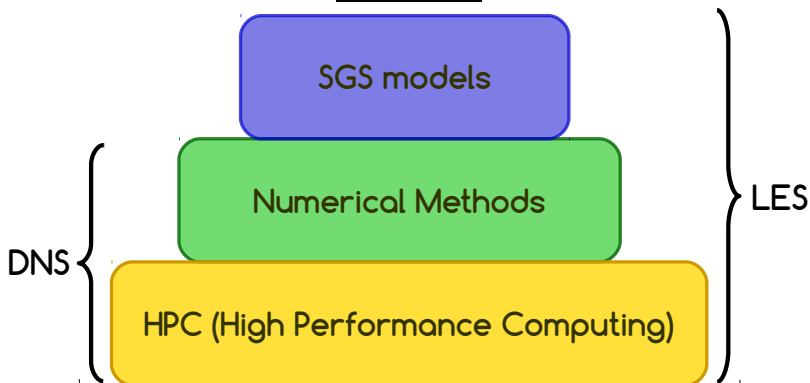
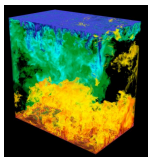


Motivation

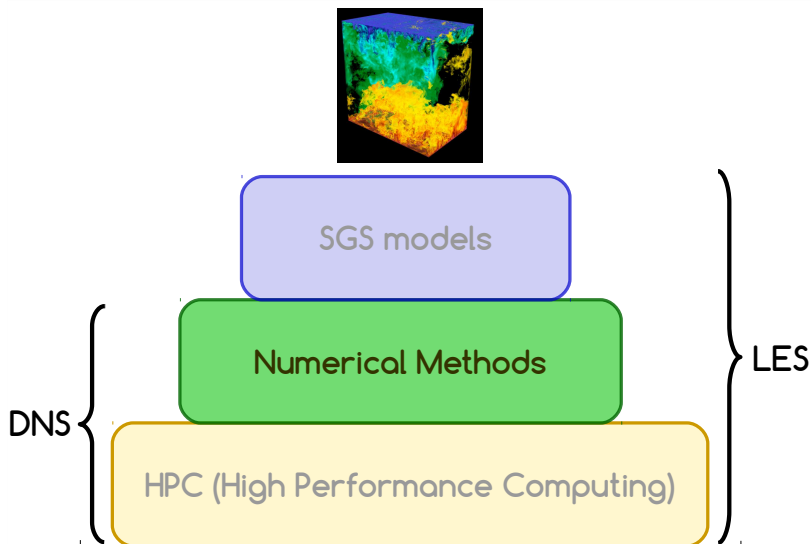
How to properly model SGS?



Motivation



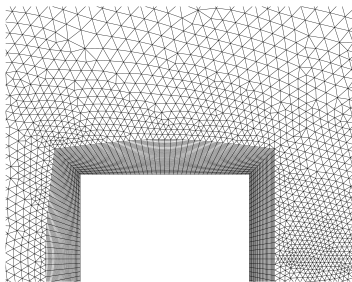
Numerical methods for DNS/LES



Numerical methods for DNS/LES

Research question #2:

- Can we construct numerical discretizations of the Navier-Stokes equations suitable for **complex geometries**, such that the **symmetry properties** are exactly preserved?



DNS¹ of the turbulent flow around a square cylinder at $Re = 22000$

¹F.X.Trias, A.Gorobets, A.Oliva. *Turbulent flow around a square cylinder at Reynolds number 22000: a DNS study*, **Computers&Fluids**, 123:87-98, 2015.

Symmetry-preserving discretization on unstructured grids³

Continuous

$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{C}(\mathbf{u}, \mathbf{u}) = \nu \nabla^2 \mathbf{u} - \nabla p$$

$$\nabla \cdot \mathbf{u} = 0$$

³F.X.Trias, O.Lehmkuhl, A.Oliva, C.D.Pérez-Segarra, R.W.C.P.Verstappen.
Symmetry-preserving discretization of Navier-Stokes equations on collocated unstructured grids, **Journal of Computational Physics**, 258 (1): 246-267, 2014.

Symmetry-preserving discretization on unstructured grids³

Continuous

Discrete

$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{C}(\mathbf{u}, \mathbf{u}) = \nu \nabla^2 \mathbf{u} - \nabla p$$

$$\nabla \cdot \mathbf{u} = 0$$

$$\Omega \frac{d\mathbf{u}_h}{dt} + \mathbf{C}(\mathbf{u}_h) \mathbf{u}_h = \mathbf{D}\mathbf{u}_h - \mathbf{G}\mathbf{p}_h$$

$$\mathbf{M}\mathbf{u}_h = \mathbf{0}_h$$

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$$\langle \mathbf{a}_h, \mathbf{b}_h \rangle_h = \mathbf{a}_h^T \Omega \mathbf{b}_h$$

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$$\langle \mathbf{C}(\mathbf{u}, \varphi_1), \varphi_2 \rangle = - \langle \mathbf{C}(\mathbf{u}, \varphi_2), \varphi_1 \rangle$$

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Symmetry-preserving discretization on unstructured grids³

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$$\mathbf{D} = \mathbf{D}^T \quad \text{def -}$$

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Why collocated arrangements are so popular?

- STAR-CCM+



SIEMENS



- ANSYS-FLUENT



- Code-Saturne



- OpenFOAM

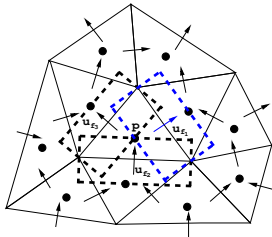
OpenFOAM®



$$\Omega_s \frac{d\mathbf{u}_s}{dt} + \mathbf{C}(\mathbf{u}_s) \mathbf{u}_s = \mathbf{D}\mathbf{u}_s - \mathbf{G}p_c; \quad \mathbf{M}\mathbf{u}_s = \mathbf{0}_c$$

In staggered meshes

- p - \mathbf{u}_s coupling is naturally solved ✓
- $\mathbf{C}(\mathbf{u}_s)$ and \mathbf{D} difficult to discretize ✗



Why collocated arrangements are so popular?

Everything is easy except the pressure-velocity coupling

- STAR-CCM+



SIEMENS



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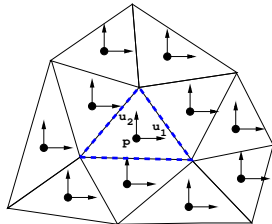
OpenFOAM®



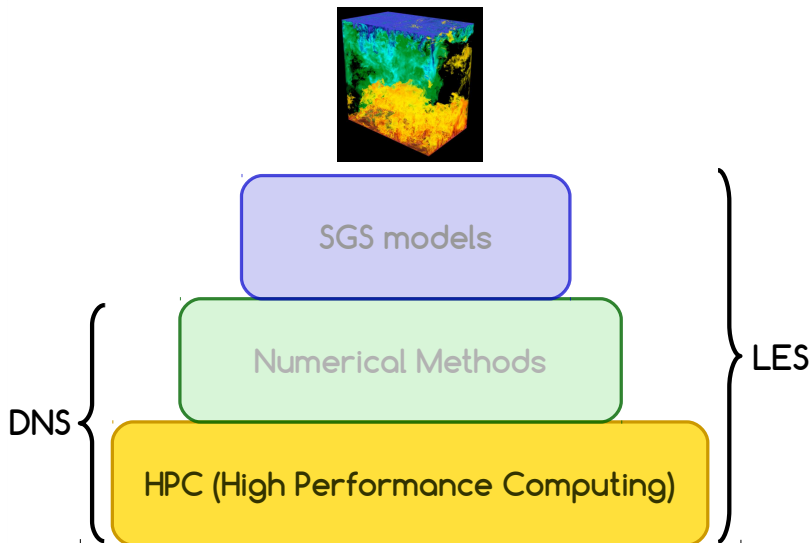
$$\Omega_c \frac{d\mathbf{u}_c}{dt} + \mathbf{C}(\mathbf{u}_s) \mathbf{u}_c = \mathbf{D} \mathbf{u}_c - \mathbf{G}_c \mathbf{p}_c; \quad \mathbf{M}_c \mathbf{u}_c = \mathbf{0}_c$$

In collocated meshes

- $p-\mathbf{u}_c$ coupling is cumbersome \times
- $\mathbf{C}(\mathbf{u}_s)$ and \mathbf{D} easy to discretize \checkmark
- Cheaper, less memory, ... \checkmark



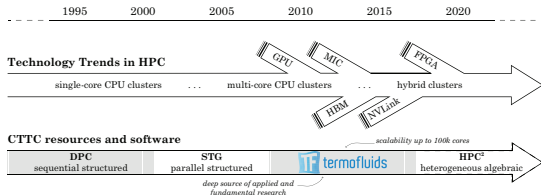
HPC on modern supercomputers



HPC on modern supercomputers

Research question #3:

- How can we develop **portable** and **efficient** CFD codes for large-scale simulations on modern supercomputers?



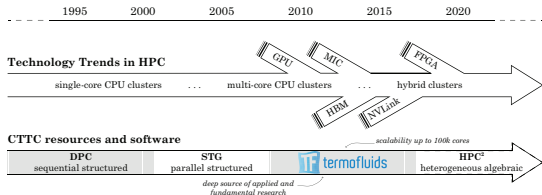
⁴X.Álvarez, A.Gorobets, F.X.Trias. *A hierarchical parallel implementation for heterogeneous computing. Application to algebra-based CFD simulations on hybrid supercomputers.* **Computers & Fluids**, 214:104768, 2021

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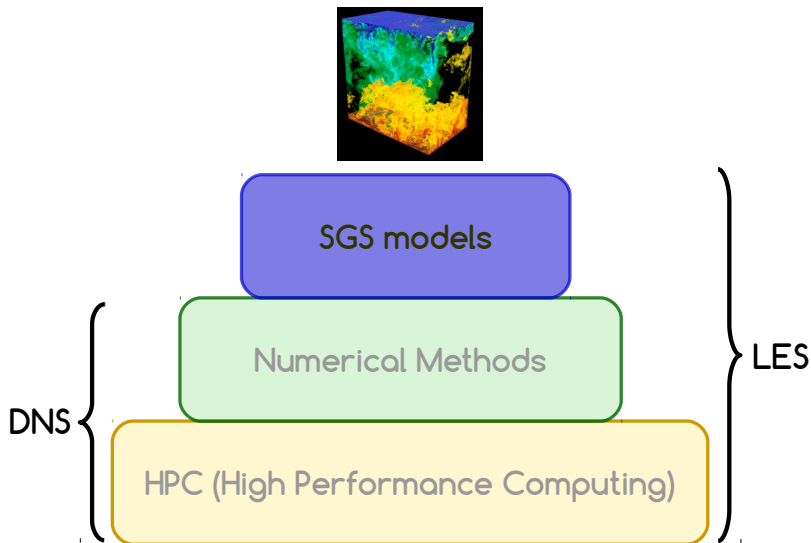


HPC²: portable, algebra-based framework for heterogeneous computing is being developed⁴. Traditional stencil-based data and sweeps are replaced by algebraic structures (sparse matrices and vectors) and kernels. SpMM-based strategies to increase the arithmetic intensity are under development⁵.

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LES of RBC



Problems to model the SGS heat flux⁶

$$\partial_t \bar{\mathbf{u}} + (\bar{\mathbf{u}} \cdot \nabla) \bar{\mathbf{u}} = \nu \nabla^2 \bar{\mathbf{u}} - \nabla \bar{p} - \nabla \cdot \boldsymbol{\tau}(\bar{\mathbf{u}}) ; \quad \nabla \cdot \bar{\mathbf{u}} = 0$$

$$\text{eddy-viscosity} \longrightarrow \boldsymbol{\tau}(\bar{\mathbf{u}}) = -2\nu_t \mathbf{S}(\bar{\mathbf{u}})$$

$$\nu_t \approx (C_m \delta)^2 D_m(\bar{\mathbf{u}}) \longrightarrow \{\text{WALE, Vreman, QR, Sigma, S3PQR, ...}\}$$

⁶F.Dabbagh, F.X.Trias, A.Gorobets, A.Oliva. *A priori study of subgrid-scale features in turbulent Rayleigh-Bénard convection*, **Physics of Fluids**, 29:105103, 2017.

Problems to model the SGS heat flux⁶

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eddy-diffusivity

gradient model

$$\mathbf{q} \approx -\alpha_t \nabla \bar{T} \quad (\equiv \mathbf{q}^{\text{eddy}})$$

$$\mathbf{q} \approx -\frac{\delta^2}{12} G \nabla \bar{T} \quad (\equiv \mathbf{q}^{nl})$$

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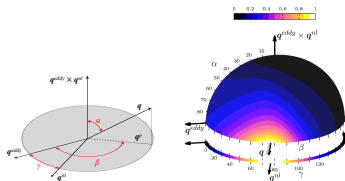
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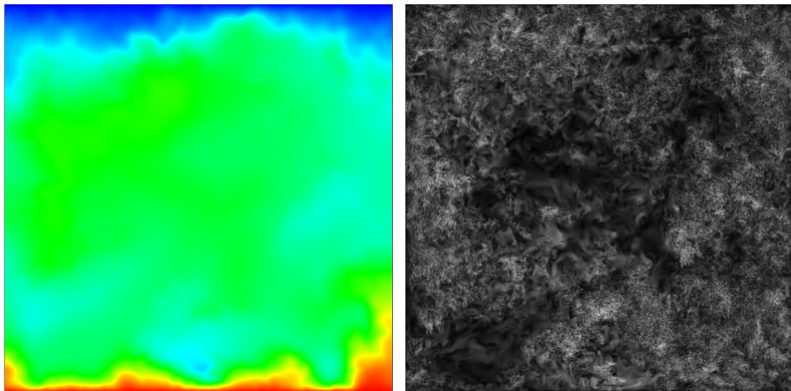
DNS results at very low Pr number

Why? scale separation grows as $\eta_K/\eta_T = Pr^{3/4}$.

η_T : Obukhov-Corrsin scale; η_K : Kolmogorov scale

DNS results at very low Pr number

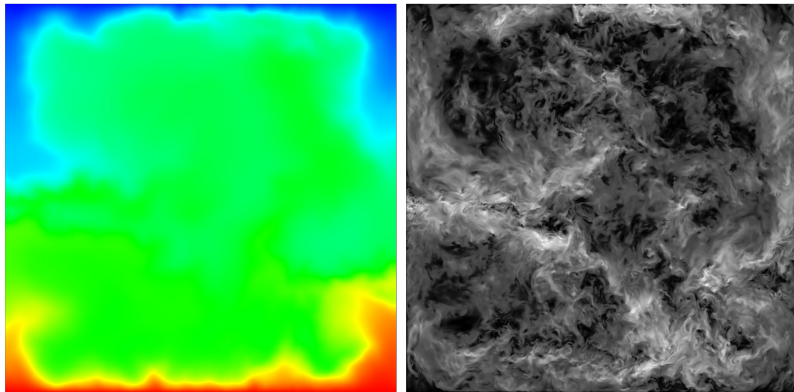
Why? scale separation grows as $\eta_K/\eta_T = Pr^{3/4}$. Here: $\eta_T \approx 53.2\eta_K$
 η_T : Obukhov-Corrsin scale; η_K : Kolmogorov scale



DNS of a RB at $Ra = 7.14 \times 10^6$ and $Pr = 0.005$ (liquid sodium)
 $488 \times 488 \times 1280 \approx 305M$

DNS results at very low Pr number

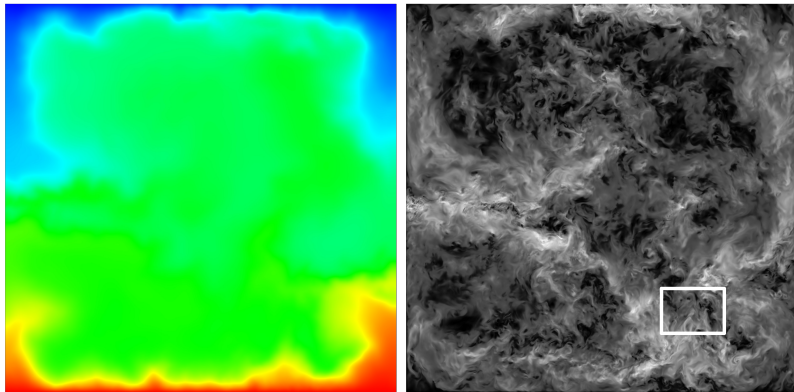
Why? scale separation grows as $\eta_K/\eta_T = Pr^{3/4}$. Here: $\eta_T \approx 53.2\eta_K$
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DNS of a RB at $Ra = 7.14 \times 10^7$ and $Pr = 0.005$ (liquid sodium)
 $966 \times 966 \times 2048 \approx \mathbf{1911M}$

DNS results at very low Pr number

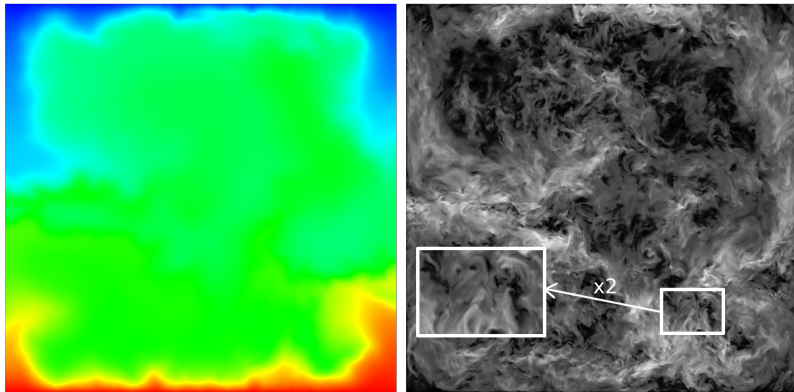
Why? scale separation grows as $\eta_K/\eta_T = Pr^{3/4}$. Here: $\eta_T \approx 53.2\eta_K$
 η_T : Obukhov-Corrsin scale; η_K : Kolmogorov scale



DNS of a RB at $Ra = 7.14 \times 10^7$ and $Pr = 0.005$ (liquid sodium)
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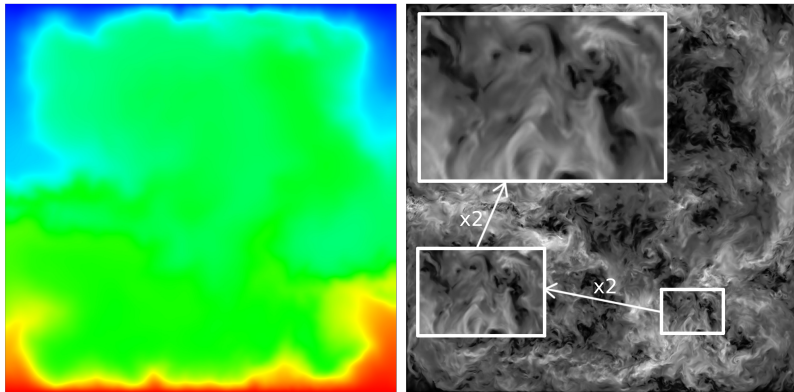
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LES of RBC Problems to model the SGS heat flux⁷

$$\partial_t \bar{\mathbf{u}} + (\bar{\mathbf{u}} \cdot \nabla) \bar{\mathbf{u}} = \nu \nabla^2 \bar{\mathbf{u}} - \nabla \bar{p} + \bar{\mathbf{f}} - \nabla \cdot \boldsymbol{\tau}(\bar{\mathbf{u}}) ; \quad \nabla \cdot \bar{\mathbf{u}} = 0$$

$$\text{eddy-viscosity} \longrightarrow \boldsymbol{\tau}(\bar{\mathbf{u}}) = -2\nu_t \mathbf{S}(\bar{\mathbf{u}})$$

$$\nu_t \approx (C_m \delta)^2 D_m(\bar{\mathbf{u}}) \longrightarrow \{\text{WALE, Vreman, QR, Sigma, S3PQR, ...}\}$$

$$\partial_t \bar{T} + (\bar{\mathbf{u}} \cdot \nabla) \bar{T} = \alpha \nabla^2 \bar{T} - \nabla \cdot \mathbf{q} \quad \text{where} \quad \mathbf{q} = \overline{\mathbf{u}T} - \bar{\mathbf{u}}\bar{T}$$

eddy-diffusivity

gradient model

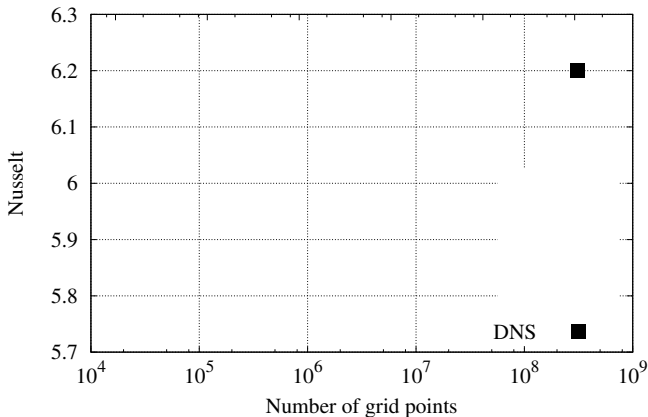
~~$$\mathbf{q} \approx -\alpha_t \nabla \bar{T} \quad (\equiv \mathbf{q}^{\text{eddy}})$$~~

~~$$\mathbf{q} \approx -\frac{\delta^2}{12} G \nabla \bar{T} \quad (\equiv \mathbf{q}^{nl})$$~~

⁷F.Dabbagh, F.X.Trias, A.Gorobets, A.Oliva. *A priori study of subgrid-scale features in turbulent Rayleigh-Bénard convection*, **Physics of Fluids**, 29:105103, 2017.

DNS vs LES results at very low Pr number⁸

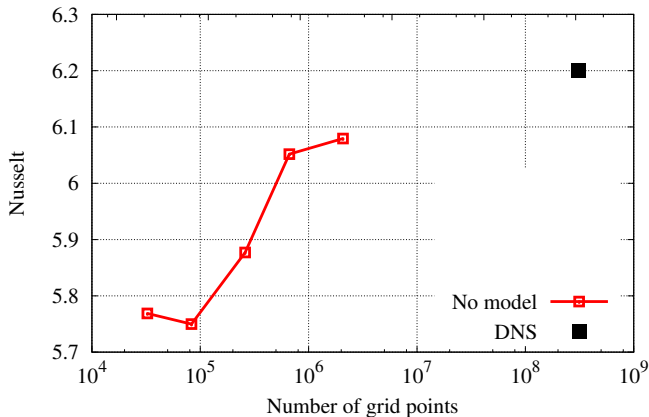
RB at $Ra = 7.14 \times 10^6$ and $Pr = 0.005$ (DNS $\rightarrow 488 \times 488 \times 1280 \approx 305M$)



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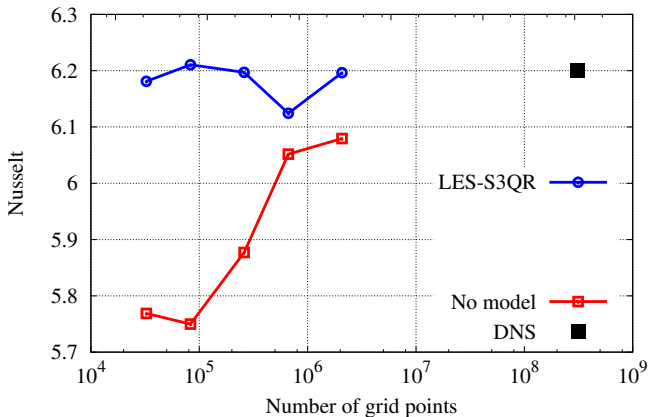
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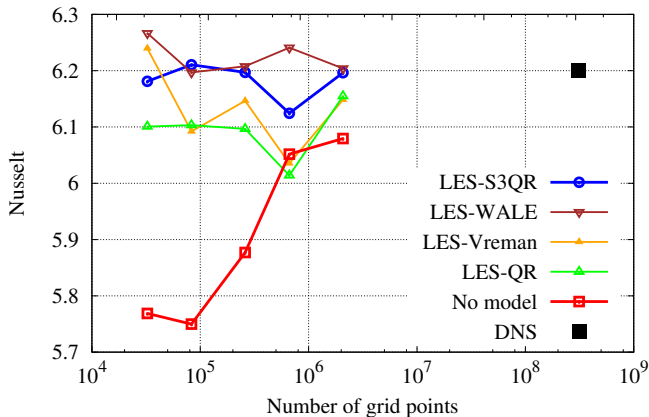
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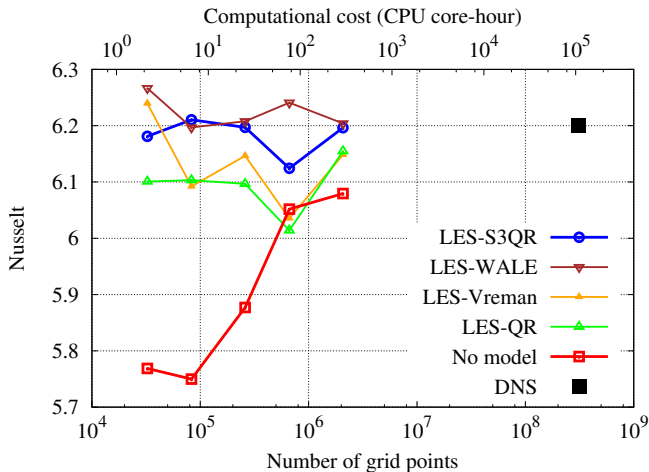
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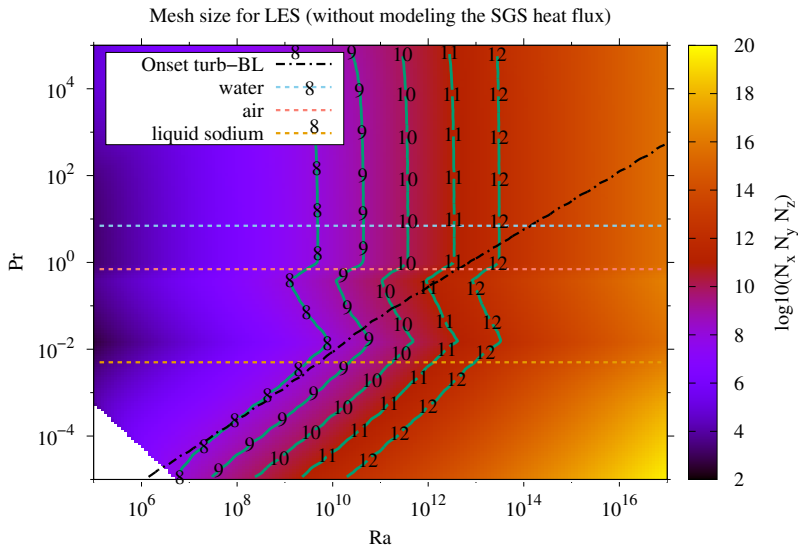
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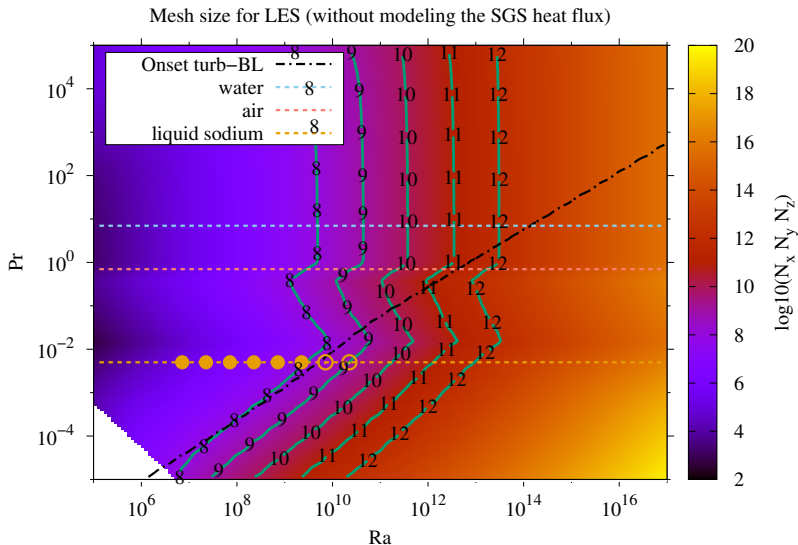
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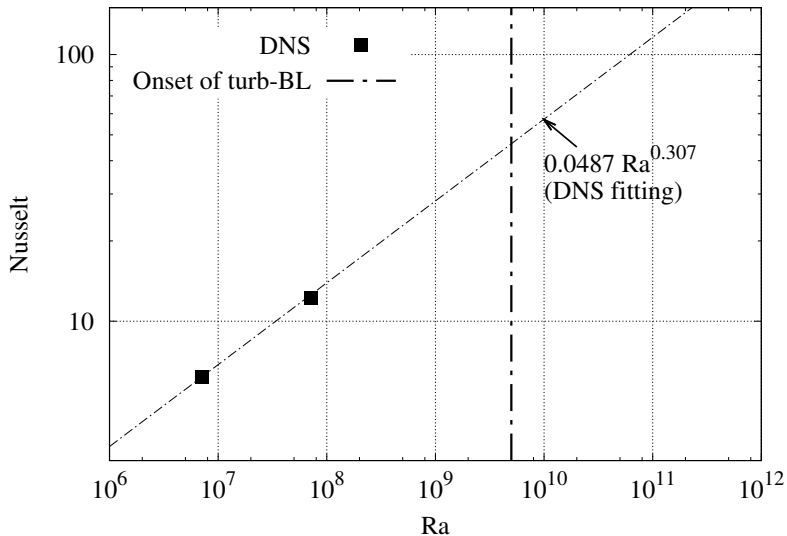


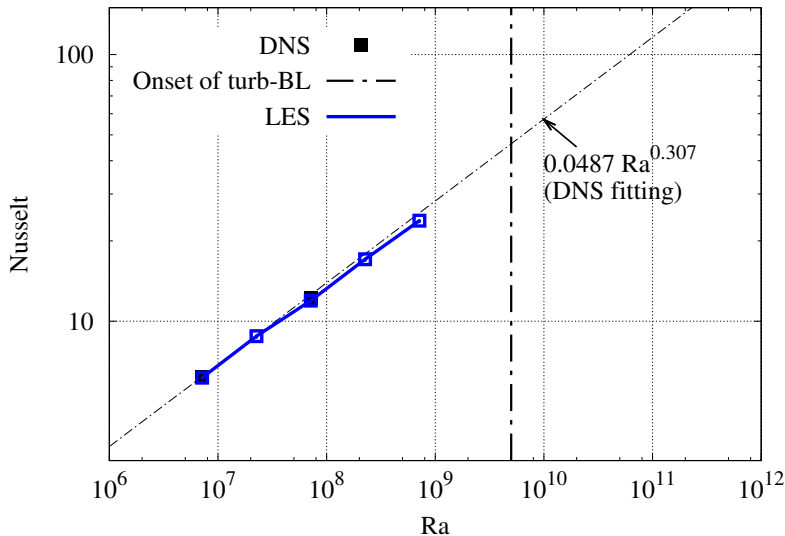
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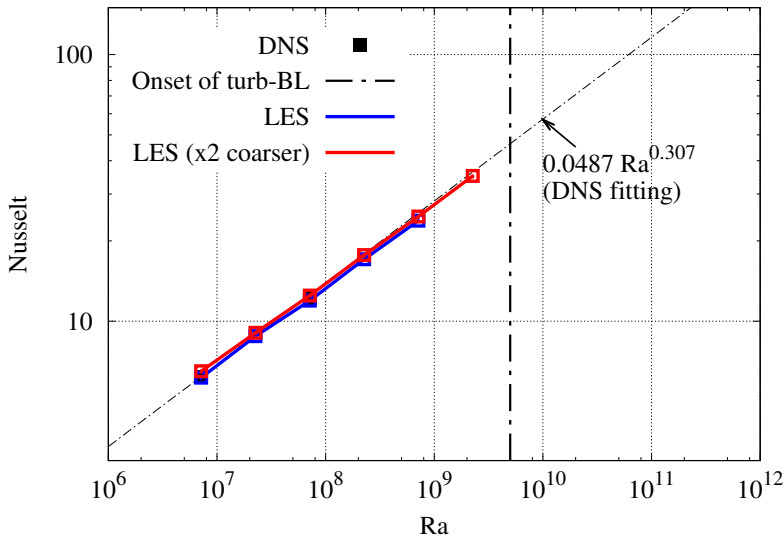
LES results at very low Pr number



LES results at very low Pr number (on-going)

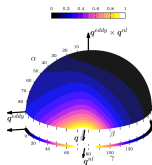
LES results at very low Pr number (on-going)

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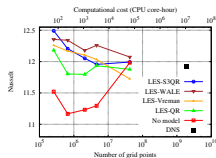
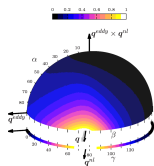
Concluding remarks

- Modeling the SGS heat flux, q , is the main difficulty for LES of buoyancy-driven flows



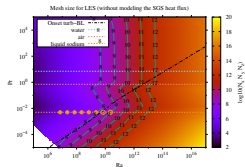
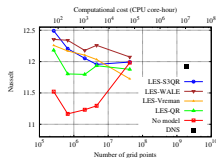
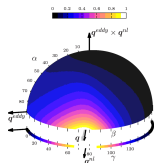
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- Ultimate regime of turbulence may be reached with LES at low-Pr ✓



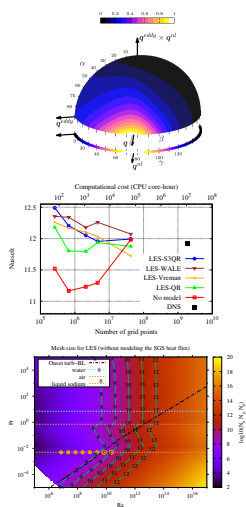
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On-going research:

- LES simulations at low- Pr and very large Ra
- Re-thinking standard CFD operators (e.g flux limiters^a, boundary conditions, CFL,...) to adapt them into an algebraic framework

^aN.Valle, X.Álvarez, A.Gorobets, J.Castro, A.Oliva, F.X.Trias. *On the implementation of flux limiters in algebraic frameworks*. **Computer Physics Communications**, 271:108230, 2022.



Thank you for your attendance