



Centre Tecnològic de Transferència de Calor  
UNIVERSITAT POLITÈCNICA DE CATALUNYA



# Exa, zetta, yotta and beyond

Àdel Alsalti-Baldellou<sup>1,2</sup>, F. Xavier Trias<sup>1</sup>, Assensi Oliva<sup>1</sup>

<sup>1</sup>Heat and Mass Transfer Technological Center, Technical University of Catalonia

<sup>2</sup>Termo Fluids S.L. Carrer de Magí Colet 8, 08204 Sabadell (Barcelona), Spain





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# Exa, zetta, yotta and beyond: On the evolution of Poisson solvers for extreme-scale simulations

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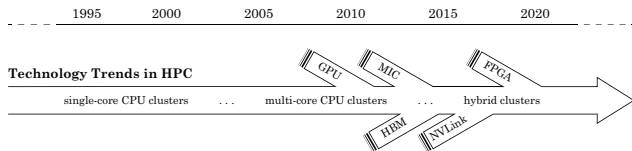
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- 2 Two competing effects
- 3 Residual of Poisson's equation
- 4 Solver convergence
- 5 Results
- 6 Conclusions

# Motivation

## Research question #1:

- How can we develop **portable** and **efficient** CFD codes for large-scale simulations on modern supercomputers?



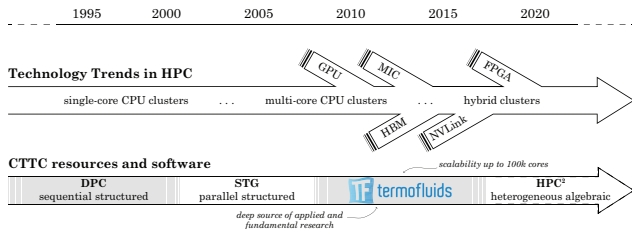
<sup>1</sup>X.Álvarez, A.Gorobets, F.X.Trias. *A hierarchical parallel implementation for heterogeneous computing. Application to algebra-based CFD simulations on hybrid supercomputers.* **Computers & Fluids**, 214:104768, 2021.

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# Motivation

## Research question #1:

- How can we develop **portable** and **efficient** CFD codes for large-scale simulations on modern supercomputers?

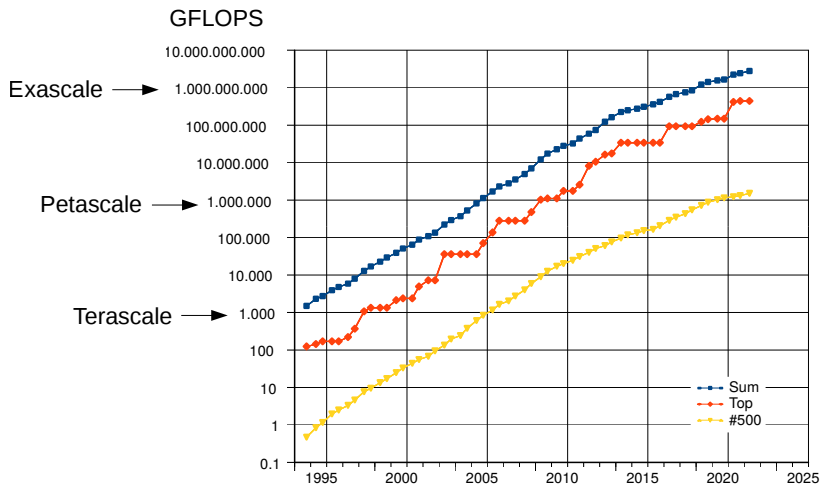


**HPC<sup>2</sup>:** portable, algebra-based framework for heterogeneous computing is being developed<sup>1</sup>. Traditional stencil-based data and sweeps are replaced by algebraic structures (sparse matrices and vectors) and kernels. SpMM-based strategies to increase the arithmetic intensity are being considered<sup>2</sup>.

<sup>1</sup>X.Álvarez, A.Gorobets, F.X.Trias. *A hierarchical parallel implementation for heterogeneous computing. Application to algebra-based CFD simulations on hybrid supercomputers.* **Computers & Fluids**, 214:104768, 2021.

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# Tera, Peta, Exa,..., Zetta, Yotta

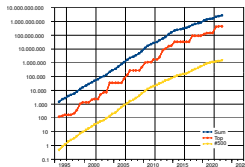


# Tera, Peta, Exa,..., Zetta, Yotta

~10 years



	PetaFLOPS		#1 in LINPACK	#1 in HPCG	Cutting-edge CFD simulation	'Routine' CFD simulation
Zetta	$10^6$					
Exa	$10^3$	11 years 14 years	2022 (Frontier)			
Peta	1		2008 (Roadrunner)	2018 (Summit)		
Tera	$10^3$		1997 (ASCI Red)	No data		

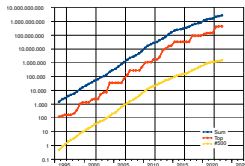


# Tera, Peta, Exa,..., Zetta, Yotta

~10 years



	PetaFLOPS		#1 in LINPACK	#1 in HPCG	Cutting-edge CFD simulation	'Routine' CFD simulation
Zetta	$10^6$		2037	2047		
Exa	$10^3$	14 years ↑	2022 (Frontier)	2032		
Peta	1	11 years ↑	2008 (Roadrunner)	2018 (Summit)		
Tera	$10^3$		1997 (ASCI Red)	No data		





# Tera, Peta, Exa, ..., Zetta, Yotta

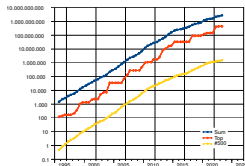
~10 years



~5 years



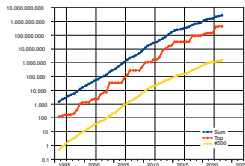
	PetaFLOPS		#1 in LINPACK	#1 in HPCG	Cutting-edge CFD simulation	'Routine' CFD simulation
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Tera	$10^3$		1997 (ASCI Red)	No data		



# Tera, Peta, Exa, ..., Zetta, Yotta

~10 years ~5 years ~10 years  
➡ ➡ ➡

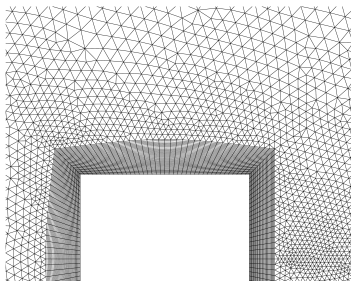
	PetaFLOPS		#1 in LINPACK	#1 in HPCG	Cutting-edge CFD simulation	'Routine' CFD simulation
Zetta	$10^6$		2037	2047	2052	2062
Exa	$10^3$	11 years 14 years	2022 (Frontier)	2032	2037	2047
Peta	1		2008 (Roadrunner)	2018 (Summit)	2023	2033
Tera	$10^3$		1997 (ASCI Red)	No data		



# Motivation

## Research question #2:

- Will the **complexity** of numerically solving **Poisson's equation** **increase** or **decrease** for **very large scale DNS/LES** simulations of incompressible turbulent flows?

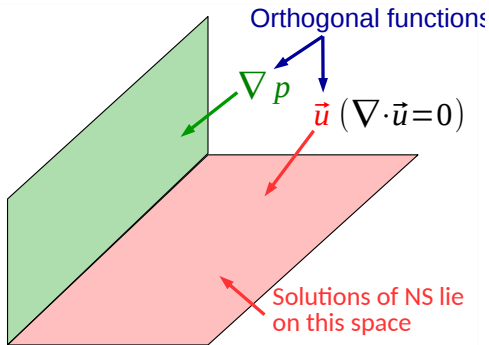


DNS<sup>3</sup> of the turbulent flow around a square cylinder at  $Re = 22000$

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<sup>3</sup>F.X.Trias, A.Gorobets, A.Oliva. *Turbulent flow around a square cylinder at Reynolds number 22000: a DNS study*, **Computers&Fluids**, 123:87-98, 2015.

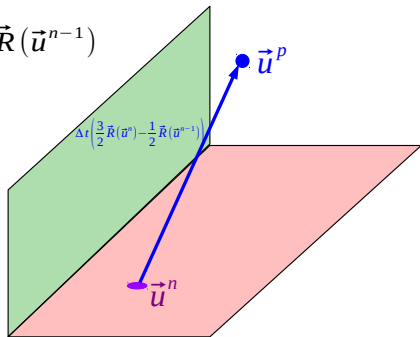
# Poisson's equation: a quick reminder



$$\text{Semi-discrete (just in time) NS equations} \left\{ \begin{array}{l} \frac{\vec{u}^{n+1} - \vec{u}^n}{\Delta t} = \frac{3}{2} \vec{R}(\vec{u}^n) - \frac{1}{2} \vec{R}(\vec{u}^{n-1}) - \nabla p^{n+1} \\ \nabla \cdot \vec{u}^{n+1} = 0 \end{array} \right.$$

# Poisson's equation: a quick reminder

$$\text{Step 1: } \frac{\vec{u}^p - \vec{u}^n}{\Delta t} = \frac{3}{2} \vec{R}(\vec{u}^n) - \frac{1}{2} \vec{R}(\vec{u}^{n-1})$$

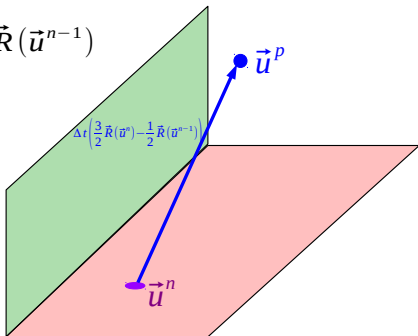


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Step 1: 
$$\frac{\vec{u}^p - \vec{u}^n}{\Delta t} = \frac{3}{2} \vec{R}(\vec{u}^n) - \frac{1}{2} \vec{R}(\vec{u}^{n-1})$$

Step 2: 
$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^p$$



Semi-discrete (just in time) NS equations

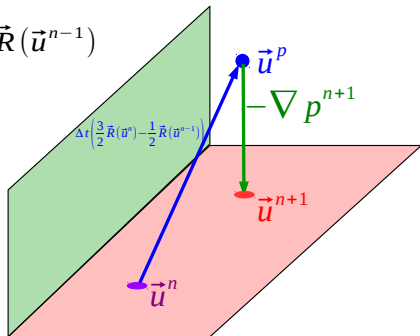
$$\left\{ \begin{array}{l} \frac{\vec{u}^{n+1} - \vec{u}^n}{\Delta t} = \frac{3}{2} \vec{R}(\vec{u}^n) - \frac{1}{2} \vec{R}(\vec{u}^{n-1}) - \nabla p^{n+1} \\ \nabla \cdot \vec{u}^{n+1} = 0 \end{array} \right.$$

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Step 2:  $\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^p$

Step 3:  $\vec{u}^{n+1} = \vec{u}^p - \Delta t \nabla p^{n+1}$



Semi-discrete (just in time) NS equations

$$\left\{ \begin{array}{l} \frac{\vec{u}^{n+1} - \vec{u}^n}{\Delta t} = \frac{3}{2} \vec{R}(\vec{u}^n) - \frac{1}{2} \vec{R}(\vec{u}^{n-1}) - \nabla p^{n+1} \\ \nabla \cdot \vec{u}^{n+1} = 0 \end{array} \right.$$

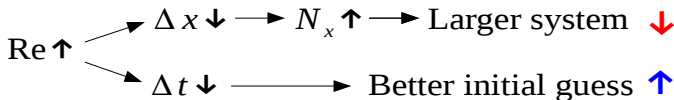
# Poisson's equation: getting more tough or not?

## Research question #2:

- Will the **complexity** of numerically solving **Poisson's equation** **increase** or **decrease** for **very large scale DNS/LES** simulations of incompressible turbulent flows?

$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^p$$

**Two competing effects:** who (if any) will eventually win?



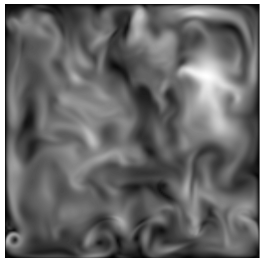


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$$Ra = 10^8$$



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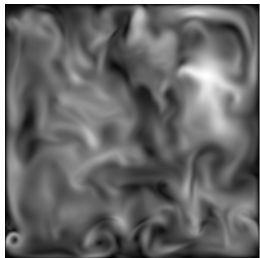
<sup>4</sup>F.Dabbagh, F.X.Trias, A.Gorobets, A.Oliva. *Flow topology dynamics in a 3D phase space for turbulent Rayleigh-Bénard convection*, **Phys.Rev.Fluids**, 5:024603, 2020.

# Poisson's equation: getting more tough or not?

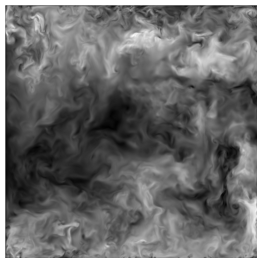
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$Ra = 10^8$



$Ra = 10^{10}$



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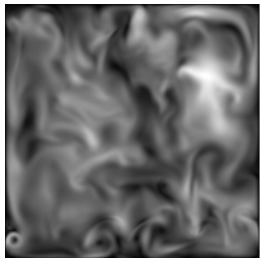
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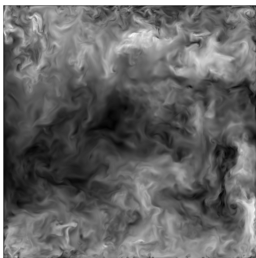
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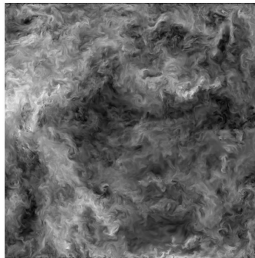
$Ra = 10^8$



$Ra = 10^{10}$



$Ra = 10^{11}$



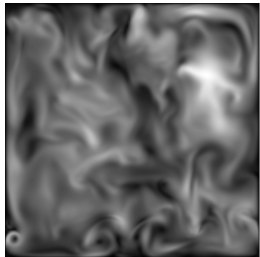
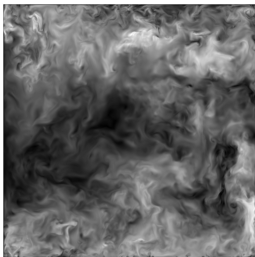
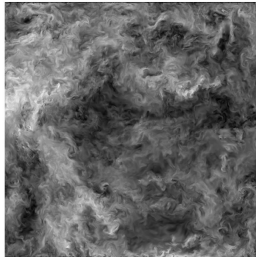
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# Poisson's equation: getting more tough or not?

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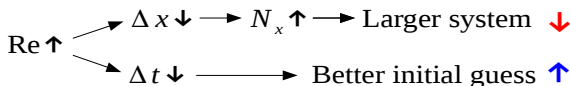
- Will the **complexity** of numerically solving **Poisson's equation** **increase** or **decrease** for **very large scale DNS/LES** simulations of incompressible turbulent flows?

 $Ra = 10^8$  $208 \times 208 \times 400$ **17.5M** $Ra = 10^{10}$  $768 \times 768 \times 1024$ **607M** $Ra = 10^{11}$  $1662 \times 1662 \times 2048$ **5600M**

<sup>4</sup>F.Dabbagh, F.X.Trias, A.Gorobets, A.Oliva. *Flow topology dynamics in a 3D phase space for turbulent Rayleigh-Bénard convection*, **Phys.Rev.Fluids**, 5:024603, 2020.

# Smaller and smaller, but how much?

Two competing effects: who (if any) will eventually win?



From classical  
K41 theory:

$$\frac{1}{N_x^{K41}} = \frac{\Delta x}{L_x} \sim \frac{\eta}{l} \propto Re^{-3/4}$$

$$\frac{u}{U} \propto Re^{-1/4}$$

$l$ : biggest eddies ( driving scale )  
 $\eta$ : smallest eddies ( Kolmogorov length scale )

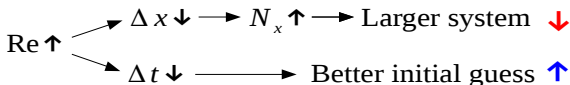
Question: ?

how  $\frac{\eta}{l}$  decreases with  $Re$ ?

$$\frac{1}{N_t^{K41}} = \frac{\Delta t}{t_{sim}} \sim \frac{t_n}{t_l} \propto \frac{\eta}{l} \frac{U}{u} \propto Re^{-3/4} Re^{1/4} = Re^{-1/2}$$

# Smaller and smaller, but how much?

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From classical  
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$$\frac{1}{N_t^{K41}} = \frac{\Delta t}{t_{\text{sim}}} \sim \frac{t_\eta}{t_l} \propto \frac{\eta}{l} \frac{U}{u} \propto \text{Re}^{-3/4} \text{Re}^{1/4} = \text{Re}^{-1/2}$$

From CFL condition:

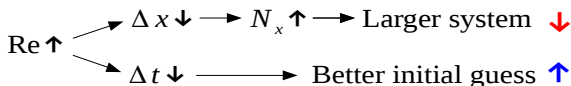
$$\Delta t^{\text{conv}} \sim \frac{\Delta x}{U} \quad \Delta t^{\text{diff}} \sim \frac{\Delta x^2}{\nu}$$

$$\frac{1}{N_t^{\text{conv}}} \sim \frac{\Delta t^{\text{conv}}}{t_l} \sim \frac{U}{l} \frac{l \text{Re}^{-3/4}}{U} = \text{Re}^{-3/4}$$

$$\frac{1}{N_t^{\text{diff}}} \sim \frac{\Delta t^{\text{diff}}}{t_l} \sim \frac{U}{l} \frac{l^2 (\text{Re}^{-3/4})^2}{\nu} = \text{Re}^{-1/2}$$

# Smaller and smaller, but how much?

**Two competing effects:** who (if any) will eventually win?



In summary:

$$\frac{1}{N_x^{K41}} = \frac{\Delta x}{L_x} \sim \frac{\eta}{l} \propto \text{Re}^{-3/4}$$

$\alpha = -1/2$  ( K41 or diffusion dominated )

$$\frac{\Delta t}{t_l} \sim \text{Re}^\alpha$$

$\alpha = -3/4$  ( convection dominated )

# Residual of Poisson's equation

$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^p$$

Initial guess  $\rightarrow p^n$

$$r^o = \nabla^2 p^n - \frac{1}{\Delta t} \nabla \cdot u^{p,n+1}$$



# Residual of Poisson's equation

$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^p$$

Initial guess  $\rightarrow p^n$

$$r^o = \nabla^2 p^n - \frac{1}{\Delta t} \nabla \cdot \vec{u}^{p,n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^{p,n} - \frac{1}{\Delta t} \nabla \cdot \vec{u}^{p,n+1} \approx \frac{\partial \nabla \cdot \vec{u}^p}{\partial t} = \nabla \cdot \frac{\partial \vec{u}^p}{\partial t}$$

## Residual of Poisson's equation

$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \tilde{\mathbf{u}}^p$$

Initial guess  $\rightarrow p^n$

$$r^o = \nabla^2 p^n - \frac{1}{\Delta t} \nabla \cdot \mathbf{u}^{p,n+1} = \frac{1}{\Delta t} \nabla \cdot \mathbf{u}^{p,n} - \frac{1}{\Delta t} \nabla \cdot \mathbf{u}^{p,n+1} \approx \frac{\partial \nabla \cdot \mathbf{u}^p}{\partial t} = \nabla \cdot \frac{\partial \mathbf{u}^p}{\partial t}$$

$$\tilde{r}^o = \nabla^2 \tilde{p}^n - \nabla \cdot \mathbf{u}^{p,n+1} \approx \nabla \cdot \mathbf{u}^{p,n} - \nabla \cdot \mathbf{u}^{p,n+1} \approx \Delta t \frac{\partial \nabla \cdot \mathbf{u}^p}{\partial t} = \Delta t \nabla \cdot \frac{\partial \mathbf{u}^p}{\partial t}$$

Initial guess  $\rightarrow \tilde{p}^n = \Delta t p^n$

$$\nabla^2 \tilde{p}^{n+1} = \nabla \cdot \tilde{\mathbf{u}}^p$$

# Residual of Poisson's equation

$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^p$$

Initial guess  $\rightarrow p^n$

$$r^o \approx \frac{\partial \nabla \cdot \vec{u}^p}{\partial t}$$

$$\tilde{r}^o \approx \Delta t \frac{\partial \nabla \cdot \vec{u}^p}{\partial t}$$

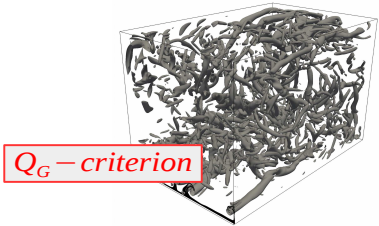
Initial guess  $\rightarrow \tilde{p}^n = \Delta t p^n$

$$\nabla^2 \tilde{p}^{n+1} = \nabla \cdot \vec{u}^p$$

# Residual of Poisson's equation

$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \tilde{u}^p$$

Initial guess  $\rightarrow p^n$



$Q_G$  - criterion

$$r^o \approx \frac{\partial \nabla \cdot u^p}{\partial t}$$

What is  $\nabla \cdot u^p$ ?

~~$$\nabla \cdot u^p = \nabla \cdot u^n - \Delta t \nabla \cdot (u^n \cdot \nabla u^n) + \nu \Delta t \nabla \cdot \nabla^2 u^n = 2 \Delta t Q_G$$~~

$$\tilde{r}^o \approx \Delta t \frac{\partial \nabla \cdot u^p}{\partial t}$$

$$Q_G = -\frac{1}{2} \text{tr}(G^2) \text{ where } G = \nabla u^n$$

Initial guess  $\rightarrow \tilde{p}^n = \Delta t p^n$

$$\nabla^2 \tilde{p}^{n+1} = \nabla \cdot \tilde{u}^p$$

## Residual of Poisson's equation

$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^p$$

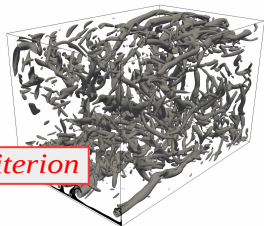
Initial guess  $\rightarrow p^n$

$$r^o \approx \frac{\partial \nabla \cdot \vec{u}^p}{\partial t} = 2 \Delta t \frac{\partial Q_G}{\partial t}$$

$$R_G = \det(G) = \frac{1}{3} \text{tr}(G^3)$$

$$\tilde{r}^o \approx \Delta t \frac{\partial \nabla \cdot \vec{u}^p}{\partial t} = 2 \Delta t^2 \frac{\partial Q_G}{\partial t}$$

$$Q_G = -\frac{1}{2} \text{tr}(G^2) \quad \text{where } G = \nabla u^n$$



$Q_G$  - criterion

*Exact equations for restricted Euler :*

$$\frac{dQ_G}{dt} = -3 R_G \longrightarrow \frac{\partial Q_G}{\partial t} = -(u \cdot \nabla) Q_G - 3 R_G$$

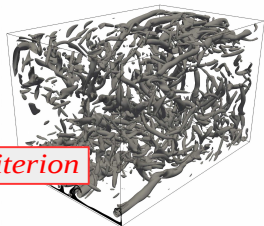
S.B.Pope, 2000.

## Residual of Poisson's equation

$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \vec{u}^p$$

Initial guess  $\rightarrow p^n$

$Q_G$  - criterion



$$r^o \approx \frac{\partial \nabla \cdot \vec{u}^p}{\partial t} = 2 \Delta t \frac{\partial Q_G}{\partial t} \approx -2 \Delta t \{ (\vec{u} \cdot \nabla) Q_G + 3 R_G \}$$

$$\tilde{r}^o \approx \Delta t \frac{\partial \nabla \cdot \vec{u}^p}{\partial t} = 2 \Delta t^2 \frac{\partial Q_G}{\partial t} \approx -2 \Delta t^2 \{ (\vec{u} \cdot \nabla) Q_G + 3 R_G \}$$

Exact equations for restricted Euler :

$$\frac{dQ_G}{dt} = -3 R_G \longrightarrow \frac{\partial Q_G}{\partial t} = -(\vec{u} \cdot \nabla) Q_G - 3 R_G$$

S.B.Pope, 2000.

## Residual of Poisson's equation

$$\nabla^2 p^{n+1} = \frac{1}{\Delta t} \nabla \cdot \tilde{u}^p$$

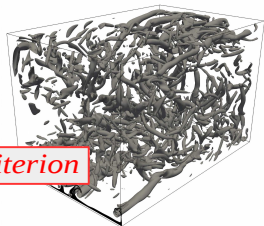
Initial guess  $\rightarrow p^n$

$$r^o \approx \frac{\partial \nabla \cdot u^p}{\partial t} = 2 \Delta t \frac{\partial Q_G}{\partial t} \approx -2 \Delta t \{ (u \cdot \nabla) Q_G + 3 R_G \}$$

$$\tilde{r}^o \approx \Delta t \frac{\partial \nabla \cdot u^p}{\partial t} = 2 \Delta t^2 \frac{\partial Q_G}{\partial t} \approx -2 \Delta t^2 \{ (u \cdot \nabla) Q_G + 3 R_G \}$$

Initial guess  $\rightarrow \tilde{p}^n = \Delta t p^n$

$$\nabla^2 \tilde{p}^{n+1} = \nabla \cdot \tilde{u}^p$$



## Residual of Poisson's equation in Fourier space

In summary:

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$$\|r\|^2 = \int_{\Omega} r^2 dV = \int_1^{k_{\max}} \hat{r}_k^2 dk$$

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# Solver convergence

$$\|r^n\|^2 = \int_1^{k_{max}} (\hat{\omega}_k^n \hat{r}_k^0)^2 dk \approx \int_1^{\text{Re}^{3/4}} \hat{\omega}_k^{2n} \text{Re}^{2\tilde{\alpha}} k^{2\beta} dk$$

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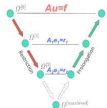
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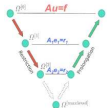
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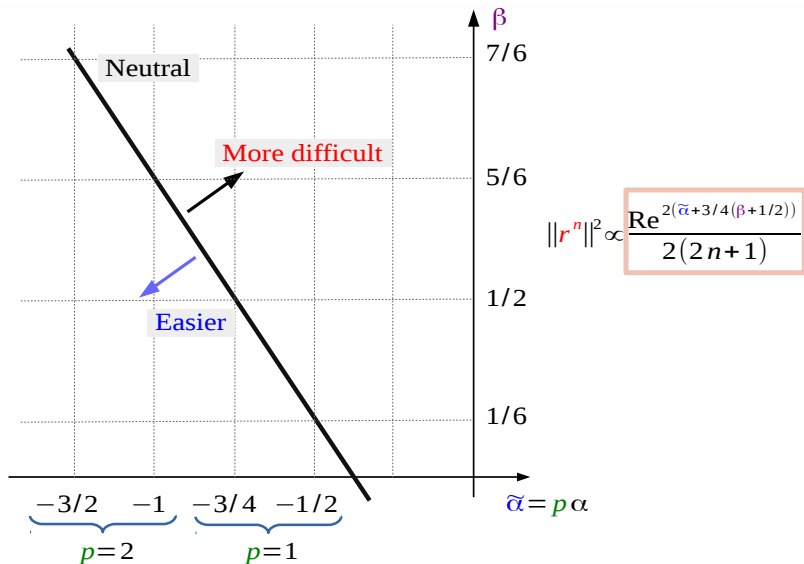
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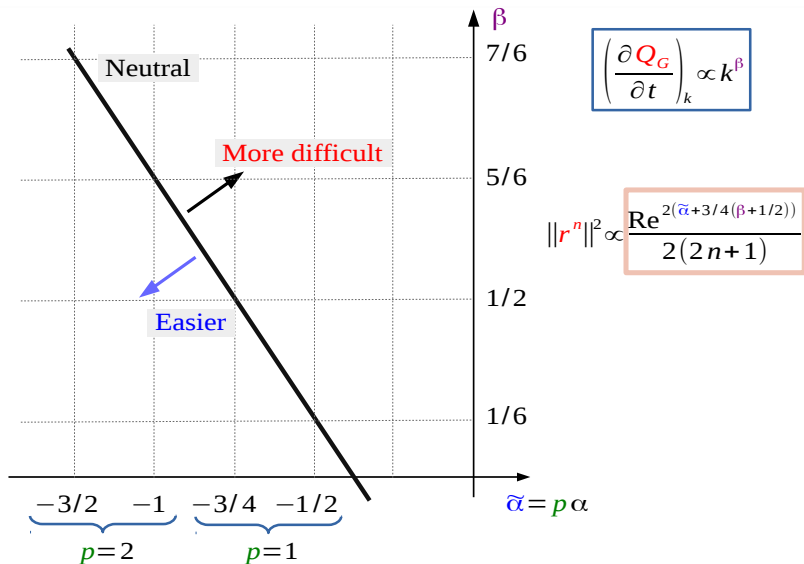
$\{\tilde{\alpha}, \beta\}$  phase space





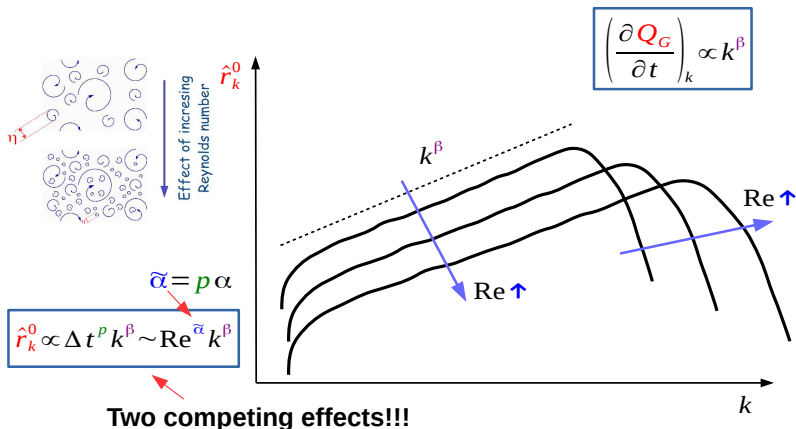
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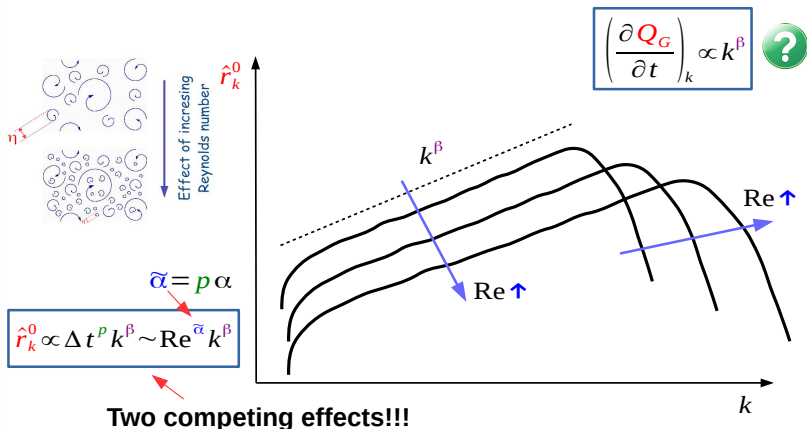


$$\frac{\Delta t}{t_l} \sim \text{Re}^\alpha$$

- $\alpha = -1/2$  (K41 or diffusion dominated)
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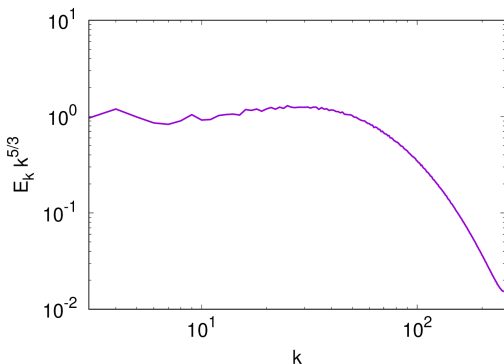
# Homogeneous isotropic turbulence

## Kolmogorov theory predictions

$$E_k = C_k \varepsilon^{2/3} k^{-5/3}$$

$$\left( \frac{\partial Q_G}{\partial t} \right)_k \propto k^\beta$$

$$\hat{r}_k^0 \propto \Delta t^p k^\beta \sim \text{Re}^{\tilde{\alpha}} k^\beta$$



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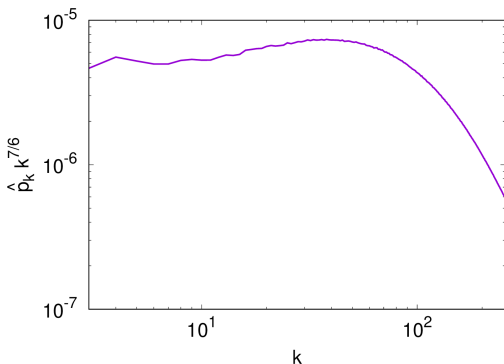
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T.Gotoh, D.Fukayama. *Pressure spectrum in homogeneous turbulence*. *Physical Review Letters*, 86(17), 3775–3778, 2001.

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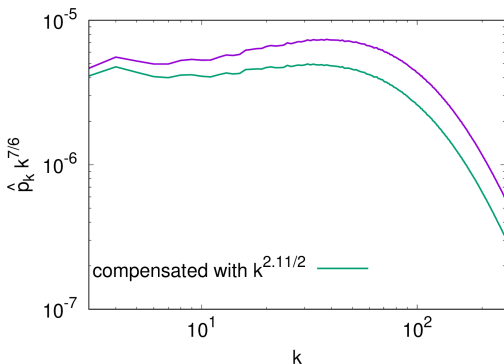
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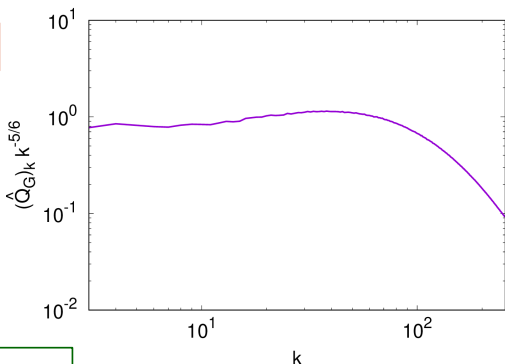
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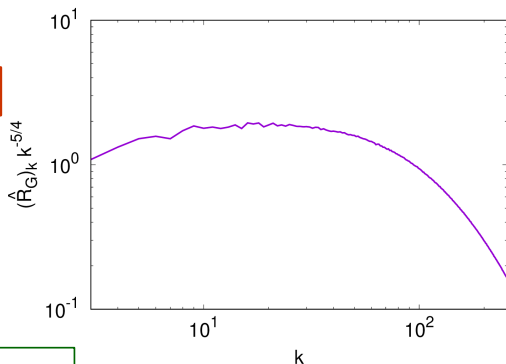
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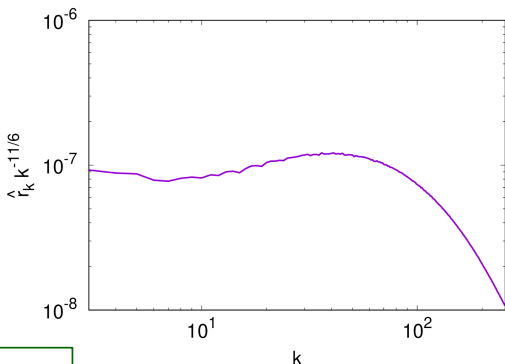
$$(\hat{r})_k \propto k^{5/6+1} = k^{11/6}$$

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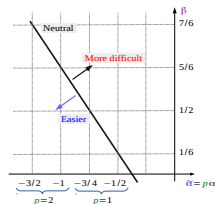
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## Concluding remarks

- **Two competing effects** on the convergence of Poisson's equation have been identified.

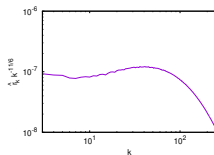
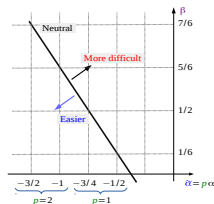
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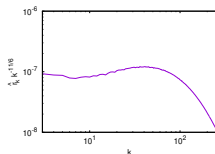
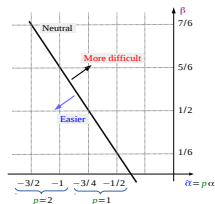
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On-going and near future research:

- Carrying out simulations at higher  $Re_\lambda$
- Extending the analysis to more complex flows

# Thank you for your attendance