New differential operators and discretization methods for eddy-viscosity models for LES

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7^{th} Workshop on Research in Turbulence and Transition

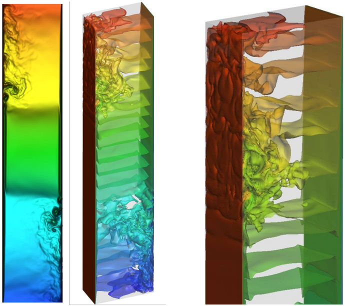
Terrassa, October 14^{th} 2013

Contents

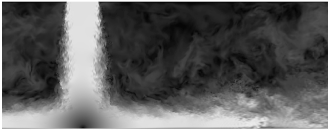
DNS of turbulent incompressible flows

Main features of the DNS code:

- Structured staggered grids
- High-order symmetry-preserving schemes
- Fully-explicit second-order time-integration method
- Poisson solver for 2.5D problems: FFT + PCG
- Hybrid MPI+OpenMP parallelization
- OpenCL-based extension for its use on GPGPU

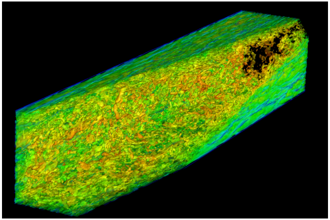


Air-filled differentially heated cavity at $Ra = 10^{11}$ (111M grid points)

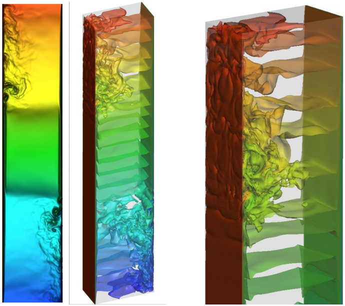


Plane impingement jet at $Re = 20000$ (102M grid points)

DNS of turbulent incompressible flows



Turbulent square duct at $Re_\tau = 1200$ (172M grid points)



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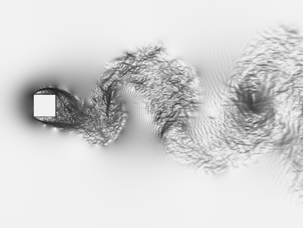
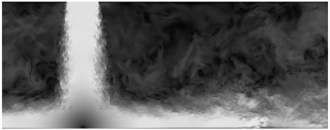


Square cylinder at $Re = 22000$ (300M grid points)



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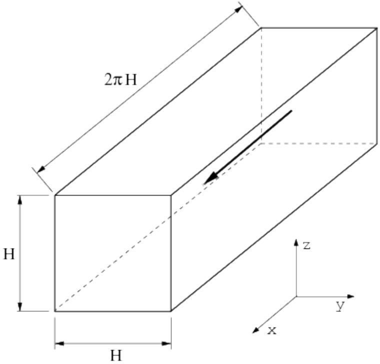
New DNS: turbulent square duct up to at $Re_\tau = 1200$



- Series of DNS simulations
  $Re_\tau = 300, 600, 900, 1200$

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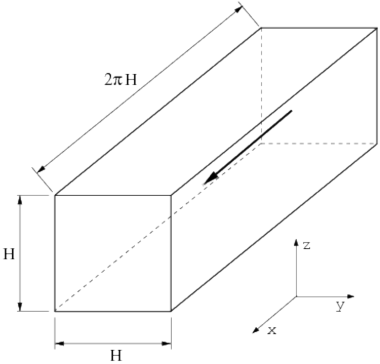

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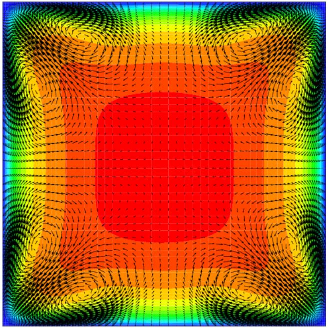

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Some details about **DNS** at $Re_\tau = 1200$

- Mesh size: $640 \times 518 \times 518$
- 392 CPUs on the MareNostrum ($P_x = 2$, $P_y = 14$, $P_z = 14$)
- 4^{th}-order symmetry-preserving scheme
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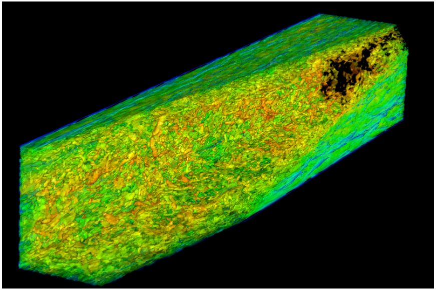
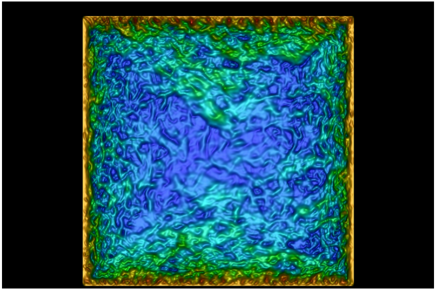
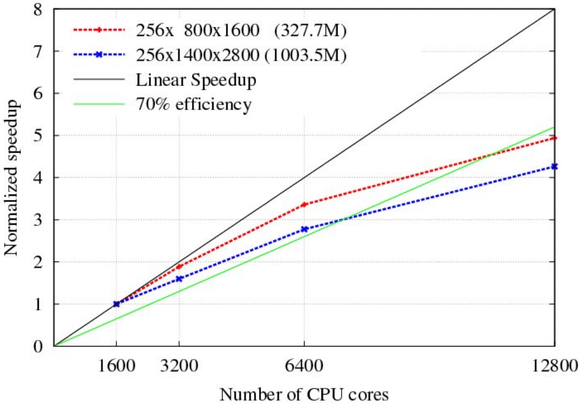




Scaling? Yes[1], we can... but never enough



[1] A. Gorobets *et al.* *"Hybrid MPI+OpenMP parallelization of an FFT-based 3D Poisson solver with one periodic direction"* **Computers&Fluids**, 49:101-109, 2011

Governing equations

Incompressible Navier-Stokes equations:

$$\begin{aligned} \nabla \cdot u &= 0 \\ \partial_t u + \mathcal{C}(u, u) &= \mathcal{D} u - \nabla p \end{aligned}$$

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where the **nonlinear convective** term is given by

$$\mathcal{C}(u, \phi) = (u \cdot \nabla)\phi$$

and the **linear dissipative** term is given by

$$\mathcal{D}\phi = \nu \Delta \phi$$

Stopping the vortex-stretching[2]

Taking the curl of momentum equation the **vorticity transport equation** follows

$$\partial_t \omega + \mathcal{C}(u, \omega) = \mathcal{C}(\omega, u) + \mathcal{D}(\omega)$$

[2] F.X. Trias *et al.* **Computers&Fluids**, 39:1815-1831, 2010

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Let us now consider an arbitrary part of the flow domain, Ω, with **periodic boundary conditions**. Then, taking the L^2 innerproduct with $\omega = \nabla \times u$ leads to the **enstrophy equation**

$$\frac{1}{2}\frac{d}{dt}(\omega, \omega) = (\omega, \mathcal{C}(\omega, u)) - \nu\left(\nabla\omega, \nabla\omega\right)$$

where $(a, b) = \int_\Omega a \cdot b d\Omega$.

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where $(a, b) = \int_\Omega a \cdot b d\Omega$. Unless, the grid is fine enough convection dominates diffusion (in a discrete sense)

$$(\omega, \mathcal{C}(\omega, u)) \; > \; \nu (\nabla \omega, \nabla \omega)$$

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Then, recalling that $\nabla \times \omega = \nabla(\nabla \cdot u) - \Delta u$ and the boundary contribution vanishes *, the diffusive term is given by the $L^2(\Omega)$-norm of Δu

$$(\nabla\omega, \nabla\omega) \stackrel{*}{=} -(\omega, \Delta\omega) = (\omega, \nabla \times \nabla \times \omega)$$
$$\stackrel{*}{=} (\nabla \times \omega, \nabla \times \omega) = (\Delta u, \Delta u) = \|\Delta u\|^2$$

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The **overall damping** introduced by a model should be given by

$$H^{\Omega} = \min\left\{\frac{\nu\|\Delta u\|^2}{4|\widetilde{R}|}, 1\right\}$$

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- Laminar flows ($R \to 0$)
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- In the wall (near-wall behavior is given by $R \propto y^1$ and $\|\Delta u\|^2 \propto y^0$)

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$$\nu_t \approx \frac{4|\widetilde{R}|}{\|\Delta u\|^2}$$

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Taking $\|\Delta u\|^2 \leq -\lambda_\Delta(\omega, \omega) = 4\lambda_\Delta \widetilde{Q}$, it becomes the eddy-viscosity model[3] based on the invariants $R = -1/3tr(S^3) = -det(S)$ and $Q = -1/2tr(S^2)$.

$\lambda_\Delta < 0$ is the largest (smallest in absolute value) non-zero eigenvalue of Laplacian operator Δ on Ω. In a periodic box of size h, $\lambda_\Delta = -(\pi/h)^2$.

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Alternatively, **regularizations** of the non-linear convective term results into a damping of vortex-stretching term, *i.e.* $f^{reg}|\tilde{R}|$ (where $0 < f \leq 1$)

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Or a combination of both?

Towards a simple LES model

Hence, a **new eddy-viscosity** model for LES

$$\partial_t \overline{u} + \mathcal{C}(\overline{u}, \overline{u}) = \mathcal{D}\overline{u} - \nabla \overline{p} - \nabla \cdot \tau(\overline{u}) \;\; ; \qquad \nabla \cdot \overline{u} = 0$$

$$\tau \; (\overline{u}) = -2\nu_t S(\overline{u})$$

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- No problems with $4|\widetilde{R}|$ and $\|\Delta \overline{u}\|^2$.
- But, what about the **discretization** of $\nabla \cdot \tau(\overline{u})$?

Discretization of $2\nabla \cdot (\nu_t S(u))$: a new simple approach[4]

$$\partial_t u + \mathcal{C}(u, u) = \mathcal{D}u - \nabla p \quad , \quad \nabla \cdot u = 0$$

$$\Omega_s \frac{du_s}{dt} + \mathsf{C}\left(u_s\right) u_s = \mathsf{D} u_s + \mathsf{M}^T p_c \quad , \quad \mathsf{M} u_s = 0_c$$

[4] F.X.Trias *et al. A simple approach to discretize the viscous term with spatially varying (eddy-)viscosity* **Journal of Computational Physics**, 253:405-417, 2013

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$$\boxed{\underbrace{-\mathsf{M}^T \Omega_c^{-1} \mathsf{M} \tilde{u}_s}_{\approx \nabla(\nabla \cdot (\nu_t u))} - \underbrace{\mathsf{C}(u_s)(-\Omega_s^{-1} \mathsf{M}^T \nu_{t,c})}_{\approx \mathcal{C}(u, \nabla \nu_t)}}$$

where $[\tilde{u}_s]_f = [\nu_{t,s}]_f [u_s]_f$. **Straightforward implementation!!!**

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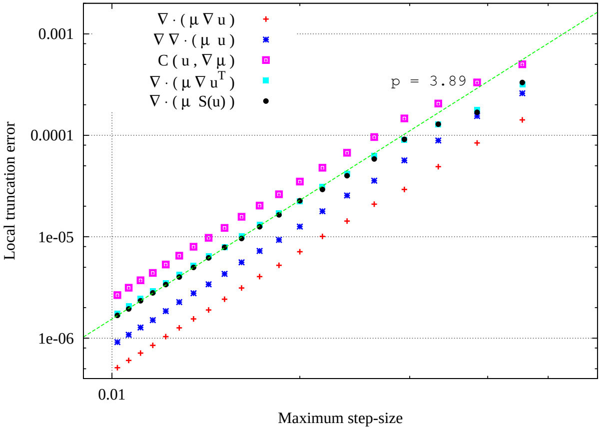
Discretization of $2\nabla \cdot (\nu_t S(u))$: a new simple approach

4^{th}-order FVM on a staggered Cartesian grid



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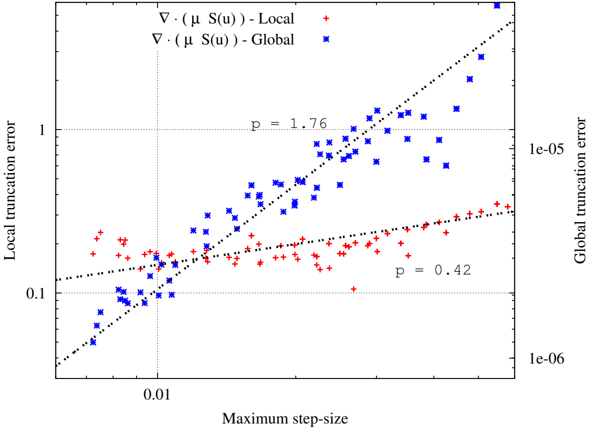
2^{th}-order FVM on a collocated unstructured grid



Discretization of $2\nabla \cdot (\nu_t S(u))$: a new simple approach

Let's make it even easier...

$$\nabla \cdot (\nu_t \nabla u^T) = \nabla(\nabla \cdot (\nu_t u)) - \mathcal{C}(u, \nabla \nu_t)$$

$$\boxed{\underbrace{-\mathsf{M}^T \Omega_c^{-1} \mathsf{M} \tilde{u}_s}_{\approx \nabla(\nabla \cdot (\nu_t u))} - \underbrace{\mathsf{C}(u_s)(-\Omega_s^{-1} \mathsf{M}^T \nu_{t,c})}_{\approx \mathcal{C}(u, \nabla \nu_t)}}$$

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Since $\nabla(\nabla \cdot (\nu_t u))$ is a gradient of a scalar field, this term can be **absorbed into the pressure**, $\pi = p - \nabla \cdot (\nu_t u)$.

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Since $\nabla(\nabla \cdot (\nu_t u))$ is a gradient of a scalar field, this term can be **absorbed into the pressure**, $\pi = p - \nabla \cdot (\nu_t u)$.

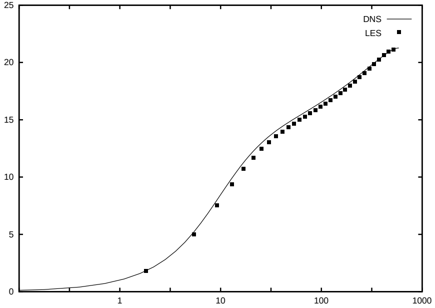
Therefore, the only term that needs to be discretized is

$$-\underbrace{\mathsf{C}(u_s)(-\Omega_s^{-1} \mathsf{M}^T \nu_{t,c})}_{\approx \mathcal{C}(u, \nabla \nu_t)}$$

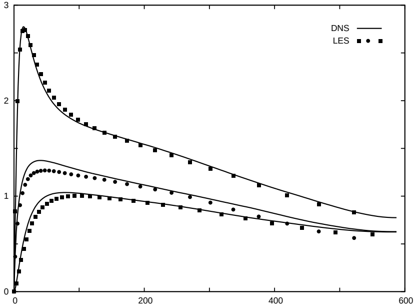
Preliminary results

Turbulent channel flow

$Re_\tau = 590$ DNS Moser et al. LES 64^3



mean velocity

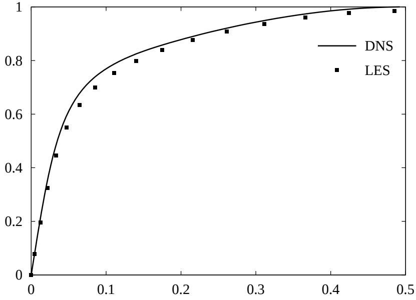


rms fluctuations

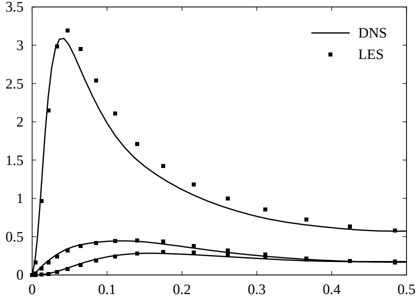
Preliminary results

Turbulent square duct

$Re_\tau = 300$ LES $64 \times 32 \times 32$



mean velocity



rms fluctuations

Conclusions

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- Test the performance of new eddy-viscosity type LES for other configurations.
- Try to properly combine regularization modeling and LES.

Thank you for your attention

Further reading

- Roel Verstappen, *"When does eddy viscosity damp subfilter scales sufficiently?"*, Journal of Scientific Computing, 49 (1): 94-110, 2011
- F.X.Trias, R.W.C.P.Verstappen, A.Gorobets, M.Soria, A.Oliva, *"Parameter-free symmetry-preserving regularization modeling of a turbulent differentially heated cavity"*, Computers & Fluids, 39:1815-1831, 2010.
- F.X.Trias, A.Gorobets, A.Oliva, *A simple approach to discretize the viscous term with spatially varying (eddy-)viscosity*, Journal of Computational Physics, 253:405-417, 2013.